

Spontaneous breaking of non-abelian (global) symmetries: $SU(2)$ examples

- main point: similar to $U(1)$, but richer, i.e., multiple possibilities for pattern of breaking, since start with > 1 symmetry generators

SSB of $SU(2)$ by (real) scalar triplet VEV (more in HW 5.3)

- Consider (pure) pion part of π -N Lagrangian of before (obviously, fermions - nucleons here - cannot get VEVs!).

$$\mathcal{L}_{\Phi} = \frac{1}{2} (\partial^{\mu} \underline{\Phi}) (\partial_{\mu} \underline{\Phi}) - \frac{1}{2} \mu^2 \underline{\Phi}^T \underline{\Phi} - \frac{1}{4} \lambda (\underline{\Phi}^T \underline{\Phi})^2$$

where $\underline{\Phi} = (\phi_1 \phi_2 \phi_3)^T$ is **real** triplet [adjoint representation of $SU(2)$] scalar field

- Choose $\mu^2 < 0$ such that $\langle 0 | \underbrace{\underline{\Phi}^T \underline{\Phi}}_{\phi_1^2 + \phi_2^2 + \phi_3^2} | 0 \rangle = v^2 = -\frac{\mu^2}{\lambda}$

- Just like $U(1)$ case, where $\langle 0 | (\phi_1 + i\phi_2) / \sqrt{2} | 0 \rangle = \frac{v}{\sqrt{2}}$,

vacuum in $SU(2)$ model is not unique: let us choose (without loss of generality), only $\langle 0 | \Phi_3 | 0 \rangle \neq 0$, i.e., $\langle 0 | \underline{\Phi} | 0 \rangle = \begin{pmatrix} 0 \\ 0 \\ v \end{pmatrix}$
column-vector

– Plug above VEV into $\mathcal{L}_{\underline{\Phi}}$ (see HW 5.3.1) to find $\phi_{1,2}$ are massless, while (fluctuation around VEV of) ϕ_3 is massive ...

... we sort of expected to find massless scalars as NGB's due to SSB, i.e., $\phi_{1,2}$ are (strong) "candidates" for being NGB's

– In fact, Goldstone's theorem says there is a massless scalar (NGB) corresponding to each broken symmetry generator \Rightarrow in this model, at most 2 generators are broken ... out of 3 ($= 2^2 - 1$) to begin with for $SU(2) \Rightarrow$ at least one symmetry is unbroken in vacuum ...

... which was "expected" as follows: $su(2)$ here is really $so(3)$ acting on $\underline{\Phi}$, i.e.,

"rotations" among ϕ_a ($a=1,2,3$) \Rightarrow
 $\langle 0 | \phi_3 | 0 \rangle$ breaks this rotational
symmetry by "picking a direction in 3d"...
but **2** d rotations, i.e., among $\phi_{1,2}$, are
still a symmetry: this is $SO(2)$
rotating 2 real scalar fields into each other
or $U(1)$ on 1 complex scalar field

— So, we have
 $SU(2)$ [3 generators] \rightarrow $U(1)$ [1 generator], i.e.,
(only) **2** generators are broken
(out of original **3**) \Rightarrow **2** NGB's...
... which then "must be" $\phi_{1,2}$

— Check explicitly unbroken $U(1)$
in HW 5.3.1; another way to
"see" unbroken symmetry in
HW 5.3.2; "identity" of NGB
matches Goldstone's theorem HW 5.3.3

- Rank of group is number of diagonal generators in fundamental representation, i.e., 1 for $SU(2)$ & $U(1)$

\Rightarrow In general, rank of group is not changed/reduced if SSB by adjoint VEV [here, triplet for $SU(2)$]

SSB of $SU(2)$ by VEV of (complex)

doublet scalar field [cannot have

real doublet, since (complex) unitary rotations involved... triplet can be real, since real, orthogonal rotations]

- Subtlety (see HW 5.4): 3 massless scalars (NGB's), hence 3 broken generators, i.e., all original $SU(2)$...

... but starting symmetry (number of generators) is actually larger than 3 (again, 3 broken) \Rightarrow

there's remnant symmetry!

... in general, rank reduced by VEV of fundamental scalar