

Comments on pion (π) - nucleon (N) system

(1). $SU(2)$ acting on fermion doublet $\Psi = \begin{pmatrix} \psi_p \\ \psi_n \end{pmatrix}$:
 capital Ψ (not ψ !) \leftarrow Pauli matrices

$$\boxed{\Psi \rightarrow \exp(-i\beta_a \sigma_a^a/2) \Psi} \quad \text{scalar triplet}$$

is isomorphic to $SO(3)$ acting on $\Phi = \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$:

$$\boxed{\Phi \rightarrow \exp(-i\beta_a T_a^{\text{adj.}}) \Phi}$$

where $(T_a^{\text{adj.}})_{bc} = -i\epsilon_{abc}$ ($\epsilon_{123} = -\epsilon_{132} = 1 \dots$)

since Lie algebra/commutations relations are the same:

$$\boxed{[T^a, T^b] = i\epsilon^{abc} T_c} \quad \text{for } T_a = \frac{\sigma_a}{2} \quad \text{or} \quad \begin{matrix} T_a^{\text{adj.}} \\ \downarrow \\ 2 \times 2 \text{ matrices} \end{matrix} \quad \begin{matrix} \text{adj.} \\ \downarrow \\ 3 \times 3 \text{ matrices} \end{matrix}$$

\Rightarrow elements of group generated by $T_a^{\text{adj.}}$ (again 3×3 matrices) are in one-to-one correspondence with those generated by $\sigma_a/2$ (2×2 matrices). Namely, (if) we have

$$\underbrace{\exp(-i\beta_a^{(1)} \sigma_a/2) \times \exp(-i\beta_b^{(2)} \sigma_b/2)}_{\substack{\text{product of 2 group elements:} \\ (2 \times 2) \text{ unitary matrices}}} = \underbrace{\exp(-i\beta_c^{(3)} \sigma_c/2)}_{\substack{\text{another group element}}},$$

then (necessarily),

$$\underbrace{\exp[-i\beta_a^{(1)} T_a^{\text{adj.}}]}_{\substack{(3 \times 3) \text{ orthogonal} \\ \text{matrix}}} \times \underbrace{\exp[-i\beta_b^{(2)} T_b^{\text{adj.}}]}_{\substack{(3 \times 3) \text{ orthogonal} \\ \text{matrix}}} = \exp[-i\beta_c^{(3)} T_c^{\text{adj.}}]$$

(2). Summary of symmetries : we have

[field $\rightarrow \exp[-i\beta_a T^a - i\alpha B] \times$ field], i.e.,

$SU(2) \times U(1)_B$, where $B = \mathbb{1}_{2 \times 2}$ for fermion doublet; $B = 0$ for scalar triplet;

$T_{ij}^a = (\sigma^a/2)_{ij}$ ($a=1,2,3$; $i,j=1,2$) for fermion doublet and $(T^a)_{bc} = (T^a_{adj.})_{bc} = -i\epsilon_{abc}$ for scalar triplet $\xrightarrow{\text{adjoint representation}}$

$-[U(1)_{EM}]$ [local gauge symmetry, after adding photon above is a subgroup of $SU(2) \times U(1)_B$ generated by $\Theta_{EM} = T_3 + \frac{1}{2}B$], i.e., $\xrightarrow{\text{3rd component of } SU(2) \text{ [isospin]}}$

[field $\rightarrow \exp[-i\beta_3 T^3 - i\beta_{3/2} B] \times$ field]

$\beta_{3/2}$ for fermion or $-i\epsilon_{3bc}$ for scalar

[set $\beta_1 = \beta_2 = 0$ and $\beta_3 = \alpha$ in the full symmetry transformation]