

Here we will show (following section 49 of chapter 7 from Landau, Lifshitz) that action variable (for 1D system with period  $T$ ), i.e.,

$$I(E; \lambda) = \frac{1}{2\pi} \int_0^T \left( \text{denoted by } \oint \right) p(q; E, \lambda) dq \quad (1)$$

where  $p(q; E, \lambda)$  is obtained by solving

$$H(q, p; \lambda) = E \quad (2)$$

is *adiabatically* invariant, i.e.,

$$\frac{dI(E; \lambda)}{dt} = 0 \text{ at } O(\dot{\lambda} \text{ or } \epsilon) \quad (3)$$

where

$$\begin{aligned} \epsilon &\equiv \frac{\bar{\dot{\lambda}} T}{\bar{\lambda}} \\ &\ll 1 \text{ (i.e., slowly varying) } \lambda \end{aligned} \quad (4)$$

with “dot” standing for time-derivative (as usual) and “bar” denoting average over one period ( $T$ ). Similarly,  $\dot{\lambda}$  itself is assumed to be slowly varying, i.e.,  $\bar{\dot{\lambda}} T / \bar{\lambda} \ll 1$  [for simplicity, we can take this factor to be  $O(\epsilon)$  also].

The above behavior of  $I$  (to be proven below) is to be compared to

$$\dot{E} \propto \dot{\lambda} \text{ (or } \epsilon) \quad (5)$$

i.e., if the Hamiltonian (via its parameter  $\lambda$ ) is time-*dependent*, then the energy of the system is not constant, that too at *leading* order in the rate of change of that parameter (cf.  $I$  above being constant at *this* order). However, given that  $\epsilon \ll 1$ , the *relative* change in  $E$  over  $T$  is *small*. [Of course,  $q$  and  $p$  change by  $O(1)$  over a time of order  $T$ , *even if*  $\lambda$  is constant.]

From Eq. (2), we have (along the actual path)

$$\frac{dE}{dt} = \frac{\partial H}{\partial \lambda} \dot{\lambda} \quad (6)$$

where the 1st factor, i.e.,  $\partial H / \partial \lambda$  is rapidly evolving over  $T$  (since it contains the  $q$  and  $p$ ), while 2nd factor ( $\dot{\lambda}$ ) is slowly varying. So, in taking average over  $T$ , we can (approximately) pull  $\dot{\lambda}$  outside, i.e., neglecting ( $\epsilon^2$ ) effects, we have

$$\frac{\overline{dE}}{dt} \approx \overline{\frac{\partial H}{\partial \lambda}} \dot{\lambda} \quad (7)$$

where “ $\dot{\lambda}$ ” can be thought of as average over  $T$ , but since its *instantaneous* value differs from the average at higher order in  $\epsilon$ , we will keep using it *without* a “bar”. Note that within our approximations, i.e., keep only  $O(\epsilon)$  effects,  $\lambda$  is taken to be *constant* while averaging  $\partial H / \partial \lambda$  above, i.e.,  $\partial H / \partial \lambda$  is varying (only) due to  $q$  and  $p$  (this is the change with time that is being averaged).

Explicitly, we get

$$\overline{\frac{\partial H}{\partial \lambda}} = \frac{1}{T} \oint \frac{\partial H}{\partial \lambda} dt \quad (8)$$

We can use  $\dot{q} = \partial H / \partial p$  or

$$dt = \frac{dq}{\frac{\partial H}{\partial p}} \quad (9)$$

thus

$$\begin{aligned} T &= \oint dt \\ &= \oint \frac{dq}{\frac{\partial H}{\partial p}} \end{aligned} \quad (10)$$

in Eq. (8) so that Eq. (7) then gives

$$\overline{\frac{dE}{dt}} \approx \dot{\lambda} \frac{\oint \frac{\partial H}{\partial \lambda} \frac{dq}{\frac{\partial H}{\partial p}}}{\oint \frac{dq}{\frac{\partial H}{\partial p}}} \quad (11)$$

Again, integral is over actual path for given, *constant*  $\lambda$  so that  $H = E$  (constant), i.e.,

$$\begin{aligned} \frac{dH}{d\lambda} &= \frac{\partial H}{\partial \lambda} + \frac{\partial H}{\partial p} \frac{\partial p}{\partial \lambda} \\ &= 0 \end{aligned} \quad (12)$$

where 2nd term in 1st line above follows from  $p(q; E, \lambda)$  as found using Eq. (2): just to be clear, note that  $q$  is instead taken to be *independent* variable (so there is *no* “ $\partial H / \partial q$ ” in above).

Plugging  $\partial H / \partial \lambda$  from Eq. (12) into Eq. (11) – and massaging a bit – gives

$$\overline{\frac{dE}{dt}} \approx - \frac{d\lambda}{dt} \frac{\oint dq \frac{\partial p}{\partial \lambda}}{\oint dq \frac{\partial p}{\partial E \text{ (or } H)}} \quad (13)$$

which can be re-written as

$$\oint dq \left( \frac{\partial p}{\partial E} \frac{dE}{dt} + \frac{\partial p}{\partial \lambda} \frac{d\lambda}{dt} \right) \approx 0 \quad (14)$$

(where  $\dot{\lambda}$  and  $\overline{dE/dt}$  can be taken inside integral within our approximations) or

$$\frac{d}{dt} \oint [p(q; E, \lambda) dq] \approx 0 \quad (15)$$

i.e., action variable of Eq. (1) is a constant (again, at leading order in  $\epsilon$ ).