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Assignment 8

2-1/3

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PHYS 402

1. Particle of mass m in an infinite square well between $x=0$ and $x=L$ with a perturbation $\hat{H}' = a\delta(x-L/4)$

a) Find: $E_n^{(1)}$. The energy eigenvalues and eigenstates for the unperturbed system are: $E_n^{(0)} = \left(\frac{\pi^2 \hbar^2}{2mL^2}\right)n^2$, $\phi_n^{(0)}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$

$$E_n^{(1)} = \langle n | H' | n \rangle = \frac{2a}{L} \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) \delta(x-L/4) dx = \boxed{\frac{2a}{L} \sin^2\left(\frac{n\pi}{4}\right)}$$

i.e. $E_n^{(1)} = \begin{cases} a/L & \text{if } n \text{ is odd} \\ 0 & \text{if } n \text{ is a multiple of 4} \\ \frac{2a}{L} & \text{otherwise} \end{cases}$

b) Second order shift in energy: $E_n^{(2)} = \sum_{k \neq n} \frac{|\langle k | H' | n \rangle|^2}{E_n^{(0)} - E_k^{(0)}}$

$$= \frac{\left(\frac{2a}{L}\right)^2 \sum_{k \neq n} \left| \int_0^L \sin\left(\frac{k\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) \delta(x-L/4) dx \right|^2}{\left(\frac{\pi^2 \hbar^2}{2mL^2}\right)(k^2 - n^2)} = \boxed{\frac{8ma^2}{\pi^2 \hbar^2} \sum_{k \neq n} \frac{\left(\sin\left(\frac{k\pi}{4}\right) \sin\left(\frac{n\pi}{4}\right)\right)^2}{k^2 - n^2}}$$

For the ground state, $n=1$.

$$E_1^{(2)} = \frac{4ma^2}{\pi^2 \hbar^2} \sum_{k \neq 1} \frac{\sin^2(k\pi/4)}{k^2 - 1} = \frac{4ma^2}{\pi^2 \hbar^2} \left[\frac{1}{2^2 - 1} + \frac{1}{2 \cdot (3^2 - 1)} + \frac{1}{2 \cdot (5^2 - 1)} + \frac{1}{6^2 - 1} + \dots \right]$$

$$E_1^{(2)} = \frac{4ma^2}{\pi^2 \hbar^2} \left[\frac{1}{3} + \frac{1}{16} + \frac{1}{48} + \frac{1}{35} + \frac{1}{96} + \frac{1}{160} + \frac{1}{99} + \dots \right]$$

0.472

$$\boxed{E_1^{(2)} = \frac{1.888 ma^2}{\pi^2 \hbar^2}}$$

c) What is the condition to be considered small, in terms of a , m , L and $\hbar$?

$$E_1 = E_1^{(0)} + E_1^{(1)} + E_1^{(2)} = \frac{\pi^2 \hbar^2}{2mL^2} + \frac{a}{L} + \frac{1.89 ma^2}{\pi^2 \hbar^2}$$

This is valid when $\frac{a}{L} \ll \frac{\pi^2 \hbar^2}{2mL^2}$ i.e. when

$$\boxed{a \ll \frac{\pi^2 \hbar^2}{2mL}}$$

2.

$$a) \hat{H} = \frac{A \hat{S}_1 \cdot \hat{S}_2}{\hbar^2}$$

$$\left. \begin{array}{l} |1,1\rangle \\ |1,0\rangle \\ |1,-1\rangle \\ |0,0\rangle \end{array} \right\} \text{possible states}$$

$$\begin{aligned} \hat{H}|1,1\rangle &= \frac{A}{\hbar^2} \frac{(\hat{S}^2 - \hat{S}_1^2 - \hat{S}_2^2)}{2} |1,1\rangle \\ &= \frac{A}{2\hbar^2} (\hbar^2(1+1)(1) - 3/2\hbar^2) |1,1\rangle \\ &= \frac{A}{4} |1,1\rangle \end{aligned}$$

$$\hat{H}|1,0\rangle = \frac{A}{4} |1,0\rangle$$

$$\hat{H}|1,-1\rangle = \frac{A}{4} |1,-1\rangle$$

$$\begin{aligned} \hat{H}|0,0\rangle &= \frac{A}{\hbar^2} \frac{(\hat{S}^2 - \hat{S}_1^2 - \hat{S}_2^2)}{2} |0,0\rangle \\ &= \frac{A}{2\hbar^2} (0\hbar^2 - 3/2\hbar^2) |0,0\rangle \\ &= -\frac{3A}{4} |0,0\rangle \end{aligned}$$

unperturbed state vectors

$$\begin{array}{l} |1,1\rangle \\ |1,0\rangle \\ |1,-1\rangle \end{array} \left. \begin{array}{l} \text{have energy} \\ \frac{A}{4} \\ \uparrow \\ E_1^{(0)} = E_{1,1}^{(0)} = E_{1,-1}^{(0)} = E_{1,0}^{(0)} \end{array} \right\} \quad \begin{array}{l} |0,0\rangle \text{ has energy} \\ -\frac{3A}{4} \\ \uparrow \\ E_0^{(0)} \end{array}$$

$$b) E_0^{(1)} = \langle 0,0 | \frac{B}{\hbar} \hat{S}_{1z} | 0,0 \rangle$$

$$|0,0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$\begin{aligned} E_0^{(1)} &= \frac{1}{2} (\langle \uparrow\downarrow | - \langle \downarrow\uparrow |) \frac{B}{\hbar} \hat{S}_{1z} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \\ &= \frac{B}{2\hbar} (\langle \uparrow\downarrow | \hat{S}_{1z} | \uparrow\downarrow \rangle + \langle \downarrow\uparrow | \hat{S}_{1z} | \downarrow\uparrow \rangle) \\ &= \frac{B}{2\hbar} \left(\frac{\hbar}{2} - \frac{\hbar}{2} \right) = 0 \end{aligned}$$

$$\begin{pmatrix} \textcircled{1} \langle 1,1 | \hat{H}' | 1,1 \rangle & \langle 1,0 | \hat{H}' | 1,1 \rangle & \langle 1,-1 | \hat{H}' | 1,1 \rangle \\ \langle 1,1 | \hat{H}' | 1,0 \rangle & \textcircled{2} \langle 1,0 | \hat{H}' | 1,0 \rangle & \langle 1,-1 | \hat{H}' | 1,0 \rangle \\ \langle 1,1 | \hat{H}' | 1,-1 \rangle & \langle 1,0 | \hat{H}' | 1,-1 \rangle & \textcircled{3} \langle 1,-1 | \hat{H}' | 1,-1 \rangle \end{pmatrix}$$

$$|1,1\rangle \longleftrightarrow |\uparrow\uparrow\rangle$$

$$|1,-1\rangle \longleftrightarrow |\downarrow\downarrow\rangle$$

$$|1,0\rangle \longleftrightarrow \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

due to orthogonality, only the diagonal matrix elements are nonzero

$$\textcircled{1} \frac{B}{\hbar} \langle \uparrow\uparrow | \hat{S}_{1z} | \uparrow\uparrow \rangle = \frac{B}{2}$$

$$\begin{aligned} \textcircled{2} \frac{B}{2\hbar} (\langle \uparrow\downarrow | + \langle \downarrow\uparrow |) \hat{S}_{1z} (|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle) &= \\ \frac{B}{2\hbar} (\langle \uparrow\downarrow | \hat{S}_{1z} | \downarrow\uparrow \rangle + \langle \downarrow\uparrow | \hat{S}_{1z} | \uparrow\downarrow \rangle) &= 0 \end{aligned}$$

$$\textcircled{3} \quad \frac{B}{\hbar} \langle \downarrow\downarrow | \hat{S}_x | \downarrow\downarrow \rangle = -\frac{B}{2}$$

$$\begin{pmatrix} B/2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -B/2 \end{pmatrix}$$

$$E_{11}^{(1)} = \frac{B}{2}$$

$$E_{110}^{(1)} = 0$$

$$E_{1,-1}^{(1)} = -\frac{B}{2}$$

in sum:

$$E_{11} = \frac{A}{4} + \frac{B}{2} + O(B^2)$$

$$E_{110} = \frac{A}{4} + O(B^2)$$

$$E_{1,-1} = \frac{A}{4} - \frac{B}{2} + O(B^2)$$

$$E_0 = -\frac{3A}{4} + O(B^2) \quad \checkmark$$

$$\textcircled{c) \quad} E_0^{(2)} = \sum_{(n,j) \neq (0,0)} \frac{|\langle n,j | \hat{H}' | 0,0 \rangle|^2}{E_0 - E_{n,j}} = \frac{|\langle 1,0 | \hat{H}' | 0,0 \rangle|^2}{E_0 - E_{10}}$$

↑ only non vanishing term

~~of the ground state~~

$$\begin{aligned} E_0^{(2)} &= \left| \frac{B}{2\hbar} (\langle \uparrow\downarrow | + \langle \downarrow\uparrow |) \hat{S}_x (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right|^2 \\ &\quad \left(-\frac{3A}{4} - \frac{A}{4} \right) \\ &= \frac{-\left(\frac{B}{2}\right)^2}{A} = -\frac{B^2}{4A} \end{aligned}$$

$$E_{1,1}^{(2)} = \sum_{(n,j) \neq (1,1)} \frac{| \langle n,j | \hat{H}' | 1,1 \rangle |^2}{E_{1,1} - E_{n,j}} = 0$$

likewise

$$E_{1,-1}^{(2)} = 0$$

$$\begin{aligned} E_{1,0}^{(2)} &= \sum_{(n,j) \neq (1,0)} \frac{| \langle n,j | \hat{H}' | 1,0 \rangle |^2}{E_{1,0} - E_{n,j}} = \frac{| \langle 0,0 | \hat{H}' | 1,0 \rangle |^2}{E_{1,0} - E_0} \\ &= \frac{\left(\frac{B}{2}\right)^2}{\left(\frac{A}{4} + \frac{3A}{4}\right)} \\ &= \frac{B^2}{4A} \end{aligned}$$

$$E_0 = -\frac{3A}{4} - \frac{B^2}{4A} + O(B^3)$$

$$E_{1,1} = \frac{A}{4} + \frac{B}{2} + O(B^3)$$

$$E_{1,0} = \frac{A}{4} + \frac{B^2}{4A} + O(B^3)$$

$$E_{1,-1} = \frac{A}{4} - \frac{B}{2} + O(B^3)$$

3. Particle in a 2-D box: $V = \begin{cases} 0 & \text{for } 0 \leq x < L \text{ and } 0 \leq y < L \\ \infty & \text{otherwise} \end{cases}$

a) The eigenenergies are $E_{nk}^{(0)} = \frac{\pi^2 \hbar^2}{2mL^2} (n^2 + k^2)$, since the x & y directions are independent.

and the eigenvalues are: $\psi_{nk}^{(0)}(x, y) = \frac{2}{L} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{k\pi y}{L}\right)$, $n = 1, 2, 3, \dots$
 $k = 1, 2, 3, \dots$

The first excited state is doubly degenerate since it can be formed with $n=1$ and $k=2$ OR $n=2$ and $k=1$. ✓

b) Perturbation: $\hat{H}' = a \cos\left(\frac{\pi x}{L}\right) \cos\left(\frac{\pi y}{L}\right)$

Let $|\alpha\rangle = |n=1, k=2\rangle = \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{2\pi y}{L}\right)$

$|\beta\rangle = |n=2, k=1\rangle = \frac{2}{L} \sin\left(\frac{2\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right)$

Find:

$$H' = \begin{pmatrix} \langle \alpha | H' | \alpha \rangle & \langle \alpha | H' | \beta \rangle \\ \langle \beta | H' | \alpha \rangle & \langle \beta | H' | \beta \rangle \end{pmatrix}$$

$$\langle \alpha | H' | \alpha \rangle = \frac{4a}{L^2} \int_0^L \int_0^L \sin^2\left(\frac{\pi x}{L}\right) \sin^2\left(\frac{2\pi y}{L}\right) \cos\left(\frac{\pi x}{L}\right) \cos\left(\frac{\pi y}{L}\right) dy dx$$

Let $\frac{\pi x}{L} = u$, $\frac{\pi y}{L} = v \Rightarrow \langle \alpha | H' | \alpha \rangle = \frac{4a}{\pi^2} \int_0^\pi \sin^2 u \cos u du \int_0^\pi \sin^2 2v \cos v dv$

$\int_0^\pi \sin^2 u \cos u du = \int_0^0 t^2 dt = 0 \Rightarrow \langle \alpha | H' | \alpha \rangle = 0$

Similarly, $\langle \beta | H' | \beta \rangle = 0$

$$\langle \alpha | H' | \beta \rangle = \frac{4a}{\pi^2} \int_0^\pi du \sin u \sin 2u \cos u \int_0^\pi \sin v \sin 2v \cos v dv = \frac{4a}{\pi^2} \left[\frac{1}{2} \int_0^\pi \sin^2 2u du \right]$$

$$= \frac{a}{4\pi^2} \left(\int_0^{2\pi} \sin^2 t dt \right)^2 = \frac{a}{4\pi^2} \cdot \pi^2 = \frac{a}{4} \Rightarrow H' = \frac{a}{4} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} W_{\alpha\alpha} & W_{\alpha\beta} \\ W_{\beta\alpha} & W_{\beta\beta} \end{pmatrix}$$

This is the Pauli sigma matrix, so its eigenstates and eigenvalues are:

$$|\psi_+\rangle = \frac{1}{\sqrt{2}} (|\alpha\rangle + |\beta\rangle) \text{ and } |\psi_-\rangle = \frac{1}{\sqrt{2}} (|\alpha\rangle - |\beta\rangle), \quad E_\pm = \pm \frac{a}{4}$$

The first order correction is:

$$E_\pm^{(1)} = \frac{1}{2} \left[W_{\alpha\alpha} + W_{\beta\beta} \pm \sqrt{(W_{\alpha\alpha} - W_{\beta\beta})^2 + 4|W_{\alpha\beta}|^2} \right]$$

$$E_\pm^{(1)} = \pm \frac{a}{8} \cdot 2 \cdot 1 = \pm \frac{a}{4}$$