

SECOND-ORDER PERTURBATION THEORY (TIME-INDEPENDENT)

We need to at least take a look at how we can continue our approximation of E_m and ψ_m , by adding one more term to the λ expansion. Again, the problem is as follows:

$$H^0 \psi_m^0 = E_m^0 \psi_m^0 \quad ; \quad \psi_m^0 \text{ is (non-degenerate) eigenstate of } \overset{\text{unperturbed}}{H^0} \text{ Hamiltonian } H^0 \text{ with eigenvalue } E_m^0$$

We add a perturbation $\lambda H'$ and now want to know the ψ_m and E_m that satisfy:

$$H \psi_m = E_m \psi_m \quad \text{where } H = H^0 + \lambda H'$$

We write:

$$\psi_m = \psi_m^0 + \lambda \psi_m^1 + \lambda^2 \psi_m^2 + \dots$$

$$E_m = E_m^0 + \lambda E_m^1 + \lambda^2 E_m^2 + \dots$$

and, upon substitution into Schroedinger's equation, we can collect terms with the same power of λ :

$$\lambda^0: H^0 \psi_m^0 = E_m^0 \psi_m^0 \Rightarrow \text{Schroedinger's equation with unperturbed Hamiltonian}$$

$$\lambda^1: \overset{\text{power}}{H'} \psi_m^0 + H^0 \psi_m^1 = E_m^1 \psi_m^0 + E_m^0 \psi_m^1 \Rightarrow E_m^1 = \langle \psi_m^0 | H' | \psi_m^0 \rangle$$

$$\psi_m^1 = \sum_{m \neq m'} \frac{\langle \psi_m^0 | H' | \psi_{m'}^0 \rangle}{E_m^0 - E_{m'}^0} \psi_{m'}^0$$

$$\lambda^2: H' \psi_m^1 + H^0 \psi_m^2 = E_m^2 \psi_m^0 + E_m^1 \psi_m^1 + E_m^0 \psi_m^2$$

Let us write the last equation in Dirac's notation, and take its direct product with $\langle \psi_m^0 |$:

$$\langle \psi_m^0 | H' | \psi_m^1 \rangle + \langle \psi_m^0 | H^0 | \psi_m^2 \rangle = E_m^2 \langle \psi_m^0 | \psi_m^0 \rangle + E_m^1 \langle \psi_m^0 | \psi_m^1 \rangle + E_m^0 \langle \psi_m^0 | \psi_m^2 \rangle$$

$$H^0 \text{ is Hermitian: } \langle \psi_m^0 | H^0 | \psi_m^0 \rangle = \langle H^0 \psi_m^0 | \psi_m^0 \rangle = E_m^0 \langle \psi_m^0 | \psi_m^0 \rangle$$

We therefore have:

$$\langle \psi_m^0 | H^1 | \psi_m^0 \rangle = E_m^2 + E_m^1 \underbrace{\langle \psi_m^0 | \psi_m^0 \rangle}$$

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= 0: remember, we choose

$$\psi_m^1 = \sum_{m' \neq m} (\dots) \psi_{m'}^0$$

$$\sum_{m' \neq m} C_m^{(m')} \langle \psi_{m'}^0 | H^1 | \psi_m^0 \rangle$$

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← used $|\psi_m^1\rangle = \sum_{m' \neq m} C_m^{(m')} \psi_{m'}^0$

$$\sum_{m' \neq m} \frac{\langle \psi_{m'}^0 | H^1 | \psi_m^0 \rangle}{E_m^0 - E_{m'}^0} \cdot \langle \psi_m^0 | H^1 | \psi_m^0 \rangle = \sum_{m' \neq m} \frac{|\langle \psi_{m'}^0 | H^1 | \psi_m^0 \rangle|^2}{E_m^0 - E_{m'}^0}$$

Finally:

$$E_m^2 = \sum_{m' \neq m} \frac{|\langle \psi_{m'}^0 | H^1 | \psi_m^0 \rangle|^2}{E_m^0 - E_{m'}^0}$$

Explicitly:

$$E_m = E_m^0 + \lambda \langle \psi_m^0 | H^1 | \psi_m^0 \rangle + \lambda^2 \sum_{m' \neq m} \frac{|\langle \psi_{m'}^0 | H^1 | \psi_m^0 \rangle|^2}{E_m^0 - E_{m'}^0} + \dots$$

$\lambda = 1$ if I write $H = H^0 + H^1$

third-order correction