

PERTURBATION THEORY (TIME-INDEPENDENT) (2)

Problem: we have an Hamiltonian H^0 of which we know all eigenstates ψ_m^0 and eigenvalues E_m^0 that solve Schrodinger's equation:

$$H^0 \psi_m^0 = E_m^0 \psi_m^0$$

We perturb the Hamiltonian by adding a term $\lambda H'$ to it, and ask ourselves what ^{we} can say about the eigenstates ψ_m and eigenvalues E_m of

$$H = H^0 + \lambda H'$$

Let us write ψ_m and E_m as a power series in λ :

$$\psi_m = \psi_m^0 + \lambda \psi_m^1 + \lambda^2 \psi_m^2 + \dots$$

$$E_m = E_m^0 + \lambda E_m^1 + \lambda^2 E_m^2 + \dots$$

We can now use those expansions in the Schrodinger's equation

$$H \psi_m = E_m \psi_m$$

and collect terms with the same power of λ . The terms multiplied by

λ form the equation:

$$H^0 \psi_m^1 + H' \psi_m^0 = E_m^0 \psi_m^1 + E_m^1 \psi_m^0$$

From the above equation we derive:

$$E_m^1 = \langle \psi_m^0 | H' | \psi_m^0 \rangle$$

First order correction to energy is expectation value of perturbation on unperturbed states

What can we say about the first-order correction on the wavefunction ψ_m^1 ?

Let us exploit the fact that the ψ_m^0 form a complete set: this means that any state can be written as a linear combination of ψ_m^0 . In particular:

$$\psi_m^1 = \sum_n c_{nm}^{(1)} \psi_n^0$$

and orthonormality

Completeness, also mean that :

$$\sum_m |\psi_m^0\rangle \langle \psi_m^0| = \mathbb{1} : \text{the identity operator}$$

Easy to prove :

completeness: any generic $|\psi\rangle$ can be written as $\sum_m C_m |\psi_m^0\rangle$

orthonormality: $\langle \psi_m^0 | \psi_n^0 \rangle = \delta_{mn}$ (Kronecker delta: $= 1$ if $m=n$, $= 0$ if $m \neq n$)

Then :

I insert this expression

$$\begin{aligned} \underline{|\psi\rangle} &= \sum_m C_m |\psi_m^0\rangle = \sum_m C_m \sum_n |\psi_m^0\rangle \langle \psi_n^0 | \psi_m^0 \rangle = \\ &= \sum_m C_m \sum_n \delta_{nm} |\psi_m^0\rangle = \sum_m |\psi_m^0\rangle \sum_n \delta_{nm} C_n = \\ &= \sum_m C_m |\psi_m^0\rangle \quad \text{same as above (m is just the name of} \\ &= \underline{|\psi\rangle} \quad \text{the index in the sum)} \end{aligned}$$

Let us now go back to our expression:

$$H^0 \psi_m^1 + H^1 \psi_m^0 = E_m^0 \psi_m^1 + E_m^1 \psi_m^0$$

$$H^0 |\psi_m^1\rangle + H^1 |\psi_m^0\rangle = E_m^0 |\psi_m^1\rangle + E_m^1 |\psi_m^0\rangle$$

\Downarrow

$$\sum_m C_m^{(m)} H^0 |\psi_m^0\rangle + H^1 |\psi_m^0\rangle = E_m^0 \sum_m C_m^{(m)} |\psi_m^0\rangle + E_m^1 |\psi_m^0\rangle$$

$$\text{use } |\psi_m^1\rangle = \sum_m C_m^{(m)} |\psi_m^0\rangle$$

$$\text{use } |\psi_m^1\rangle = \sum_m C_m^{(m)} |\psi_m^0\rangle$$

But $H^0 |\psi_m^0\rangle = E_m^0 |\psi_m^0\rangle$: $|\psi_m^0\rangle$ are solution of unperturbed H^0 .

Therefore :

$$\sum_m C_m^{(m)} (\overset{\text{numbers}}{E_m^0} - E_m^0) |\psi_m^0\rangle = - \overset{\text{operator}}{H^1} |\psi_m^0\rangle + \overset{\text{numbers}}{E_m^1} |\psi_m^0\rangle$$

Let us also point out that we can choose $C_m^{(m)}$ to be equal to ϕ .
 $C_m^{(m)}$ tells us about the component along $|\psi_m^0\rangle$ of $|\psi_m^1\rangle$.

The equation that ψ_m^1 needs to solve is:

$$H^0 \psi_m^1 + H^1 \psi_m^0 = E_m^0 \psi_m^1 + E_m^1 \psi_m^0$$

Let us rewrite it:

$$(H^0 - E_m^0) \psi_m^1 = - (H^1 - E_m^1) \psi_m^0$$

Since $(H^0 - E_m^0) \psi_m^0 = 0$ [ψ_m^0 solves the Schrodinger's equation of the unperturbed system], if ψ_m^1 solves the equation above, also $\psi_m^1 + \alpha \psi_m^0$ does, for any value of α : I can use this freedom to choose $C_m^{(m)} = 0$.

Then, I can write:

$$\sum_{m \neq m} C_m^{(m)} (E_m^0 - E_m^0) |\psi_m^0\rangle = - H^1 |\psi_m^0\rangle + E_m^1 |\psi_m^0\rangle$$

Let us take the inner product with $\langle \psi_e^0 |$:

$$\sum_{m \neq m} C_m^{(m)} (E_m^0 - E_m^0) \langle \psi_e^0 | \psi_m^0 \rangle = - \langle \psi_e^0 | H^1 | \psi_m^0 \rangle + E_m^1 \langle \psi_e^0 | \psi_m^0 \rangle$$

$$e = m: \sum_{m \neq m} C_m^{(m)} (E_m^0 - E_m^0) \delta_{mm} = - \langle \psi_m^0 | H^1 | \psi_m^0 \rangle + E_m^1$$

$$\text{oh, but } E_m^1 = \langle \psi_m^0 | H^1 | \psi_m^0 \rangle: \phi = \phi!$$

$$e \neq m: \sum_{m \neq m} C_m^{(m)} (E_m^0 - E_m^0) \delta_{em} = - \langle \psi_e^0 | H^1 | \psi_m^0 \rangle$$

$$C_e^{(m)} (E_e^0 - E_m^0) = - \langle \psi_e^0 | H^1 | \psi_m^0 \rangle$$

$$\Rightarrow C_e^{(m)} = \frac{\langle \psi_e^0 | H^1 | \psi_m^0 \rangle}{E_m^0 - E_e^0} \quad \text{and}$$

$$\psi_m^1 = \sum_{m \neq m} \frac{\langle \psi_m^0 | H^1 | \psi_m^0 \rangle}{E_m^0 - E_m^0} \psi_m^0$$