

THE VARIATIONAL PRINCIPLE: YET ANOTHER EXAMPLE

The basic idea is very simple: since the ground state minimizes the energy of a system, the expectation value of the Hamiltonian on a generic state is necessarily ^{than} larger (or equal to) the ground state energy:

$$E_{gs} \leq \langle \psi | H | \psi \rangle$$

where ψ is my guess at the ground state wavefunction. More interestingly, we can make ψ depend on some parameters $\lambda_1, \lambda_2, \lambda_3, \dots$. Since the expression above shall be valid for any ψ , we can minimize it as a function of the λ_i parameters, thus obtaining a better estimate of E_{gs} . In practice, we write a normalized wavefunction $\psi(\vec{r}, \lambda_1, \lambda_2, \dots)$ and then require:

$$\frac{\partial}{\partial \lambda_i} \langle \psi(\vec{r}, \lambda_1, \lambda_2, \dots) | H | \psi(\vec{r}, \lambda_1, \lambda_2, \dots) \rangle = 0$$

for all λ_i parameters. Let us note that there is no way to tell how close the minimized $\langle \psi | H | \psi \rangle$ is to E_{gs} , the true ground state energy.

It is useful to make the λ_i parameters correspond to a physically motivated quantity, such as an effective charge, the width of a wavefunction, the distance between two nuclei in a molecule.

Let us study the problem of a particle of mass m subject to a potential $V(x) = -\alpha \delta(x)$ (a Dirac δ function).

Exact solution:

$$\psi_{gs}(x) = \sqrt{\frac{m\alpha}{\hbar}} \exp(-m\alpha|x|/\hbar^2)$$

$$E_{gs} = -m\alpha^2/2\hbar^2$$

Let us guess a solution in the form of a Gaussian wavefunction:

$$\psi(x, b) = \left(\frac{2b}{\pi}\right)^{1/4} e^{-bx^2}$$

b tells us how wide the wavefunction is in the real space (the larger b , the narrower $\psi \implies$ the larger the uncertainty on the particle's momentum, hence the larger the expectation value of the particle's kinetic energy!).

$\psi(x, b)$ must be normalized to unity: $\langle \psi(x, b) | \psi(x, b) \rangle = 1$

The Hamiltonian is simple:

$$H = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} - \alpha \delta(x) = T + V$$

Let us calculate the expectation values:

$$\begin{aligned} \langle T \rangle &= -\frac{\hbar^2}{2m} \int_{-\infty}^{\infty} \left(\frac{2b}{\pi}\right)^{1/2} e^{-bx^2} \frac{\partial^2}{\partial x^2} (e^{-bx^2}) = \\ &= -\frac{\hbar^2}{2m} \left(\frac{2b}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} e^{-bx^2} (-2be^{-bx^2} + 4bx^2 e^{-bx^2}) = \\ &= -\frac{\hbar^2}{2m} \left(\frac{2b}{\pi}\right)^{1/2} \cdot (-2b) \cdot \left(\frac{\pi}{2b}\right)^{1/2} + \left(-\frac{\hbar^2}{2m}\right) \left(\frac{2b}{\pi}\right)^{1/2} \cdot (4b) \cdot \frac{\sqrt{8\pi}}{8} \cdot \frac{1}{b^{3/2}} = \\ &= \frac{\hbar^2 b}{m} - \frac{\hbar^2 b}{2m} = \boxed{\frac{\hbar^2 b}{2m}} = \langle T \rangle \end{aligned}$$

$$\langle V \rangle = -\alpha \int_{-\infty}^{\infty} \left(\frac{2b}{\pi}\right)^{1/2} e^{-2bx^2} \delta(x) = -\left(\frac{2bx^2}{\pi}\right)^{1/2}$$

Then:

$$\langle H \rangle = \frac{\hbar^2 b}{2m} - \left(\frac{2b\alpha^2}{\pi}\right)^{1/2} \implies \frac{\partial \langle H \rangle}{\partial b} = \frac{\hbar^2}{2m} - \frac{1}{2} \left(\frac{2\alpha^2}{\pi}\right)^{1/2} b^{-1/2}$$

Finally:

$$\frac{\partial \langle H \rangle}{\partial b} = 0 \Rightarrow b = \frac{2}{\pi} \frac{m^2 \alpha^2}{\hbar^4}$$

Correspondingly:

$$\langle H \rangle = -\frac{1}{\pi} \frac{m \alpha^2}{\hbar^2} \geq -\frac{1}{2} \frac{m \alpha^2}{\hbar^2} \text{ exact solution!}$$

Crucial notes:

1) the Hamiltonian cannot be modified. One can rearrange its pieces, and make it more explicit how expectation values depend on the λ_i parameters used to model the guessed wavefunction ψ (e.g.: He atoms and Z : we wrote $-\frac{Z}{r} = -\frac{Z}{r} + \frac{Z-2}{r}$!), but the Hamiltonian itself cannot depend on λ_i : we are allowed to modify only our guessed ψ

2) when making ψ depend on λ_i parameters, do not forget to normalize ψ :

$$\langle \psi(\lambda_i) | \psi(\lambda_i) \rangle = 1$$

the normalization constant may depend on λ_i , and that dependency is important!