

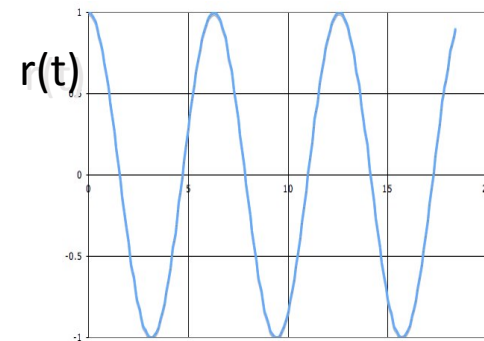
# Oscillations and Waves

What's the difference?

Oscillations involve a discrete set of quantities that vary in time (usually periodically).

Examples: pendula, vibrations of individual molecules, firefly lights, currents in circuits.

Separation between two atoms in a molecule  $r(t)$



Waves involve continuous quantities that vary in both space and time. (variation may be periodic or not)

Examples: light waves, sound waves, elastic waves, surface waves, electro-chemical waves on neurons

# Learning about Oscillations and waves

- Why to learn it
  - How the ear senses sound
  - Sound itself
  - Brain waves
  - Heart contraction waves
  - Molecule oscillations
- What to learn
  - How to describe oscillations mathematically (sin, cos)
  - How to think about waves
  - Resonances

- Heart beat
- [http://www.youtube.com/watch?annotation\\_id=annotation\\_611436&feature=iv&src\\_vid=Pes9O5z8efk&v=uR4t\\_B-Zwg](http://www.youtube.com/watch?annotation_id=annotation_611436&feature=iv&src_vid=Pes9O5z8efk&v=uR4t_B-Zwg)

- Ventricular Fibrillation

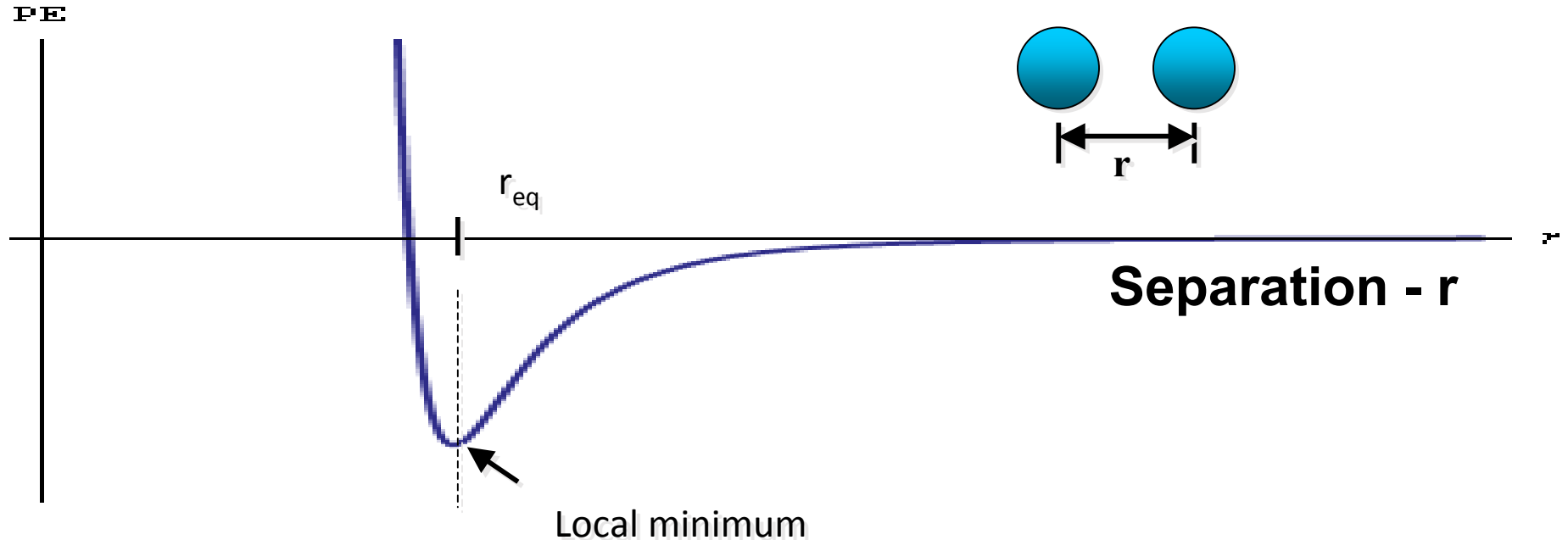
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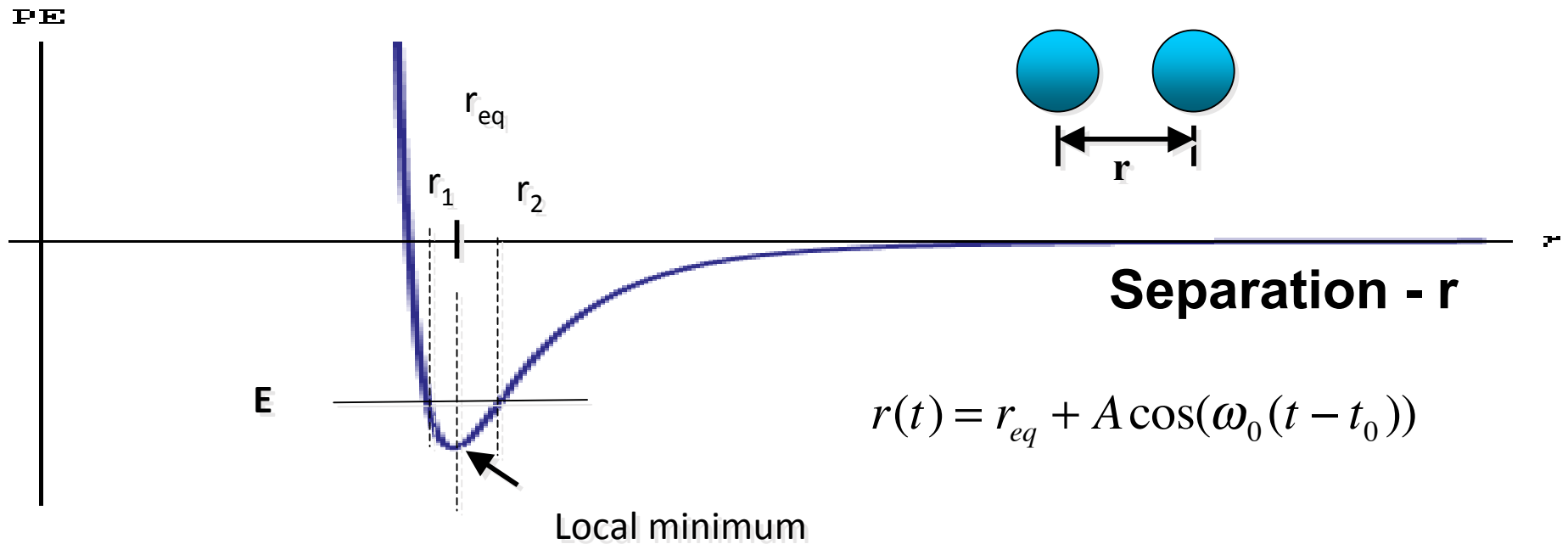
# What are the key ingredients for oscillatory motion?

1. Motion given by Newton's laws  $a = F/m$
2. Force  $F$  is derived from a potential with a local minimum.

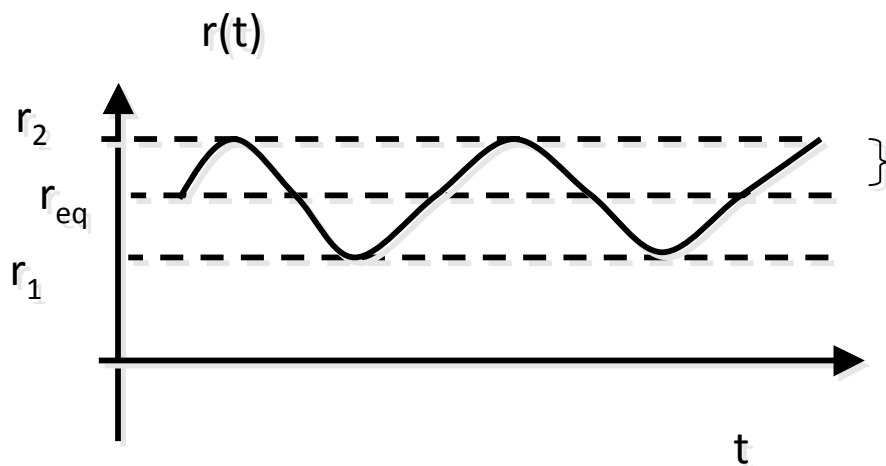
$$F(r) = -\frac{dU(r)}{dr} \simeq -k(r - r_{eq}) \quad U(r) \simeq \frac{k}{2}(r - r_{eq})^2 + U(r_{eq})$$

## Potential Energy - $U(r)$





$$r(t) = r_{eq} + A \cos(\omega_0(t - t_0))$$



$$\omega_0 = \sqrt{\frac{k}{m}}$$

natural frequency of oscillation

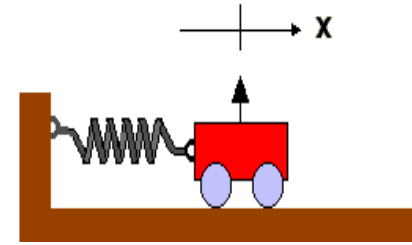
$$A = (r_2 - r_{eq})$$

Amplitude of Oscillation

$$t_0 =$$

Any time  $r(t) = r_2$

# Model system: Mass on a Spring

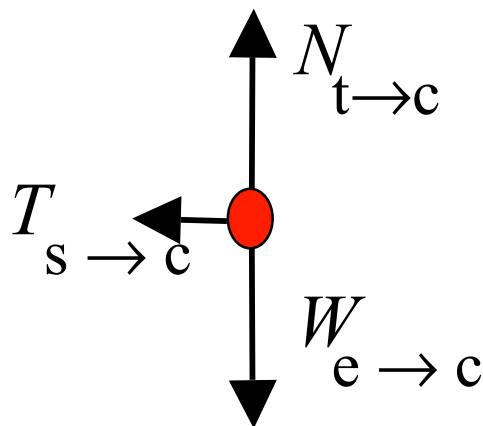
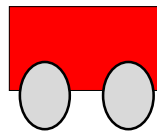
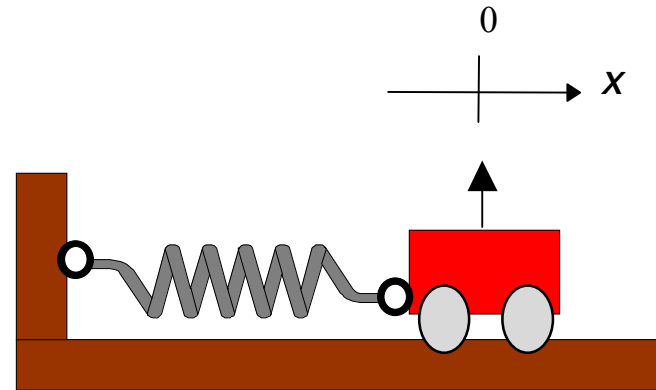


- Consider a cart of mass  $m$  attached to a light (mass of spring  $\ll m$ ) spring.
- Choose the coordinate system so that when the cart is at 0 the spring is at its rest length
- Recall the properties of an ideal spring.
  - When it is pulled or pushed on both ends it changes its length.
  - $T$  is tension,  $T > 0$  means the spring is being stretched

$$T = k\Delta l$$

# Analyzing the forces: cart & spring

- What are the forces acting on the cart?

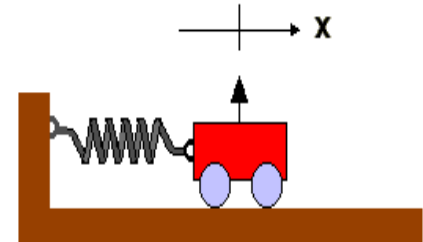


What do the  
subscripts mean  
?

A mass connected to a spring is oscillating back and forth. Consider two possibilities:

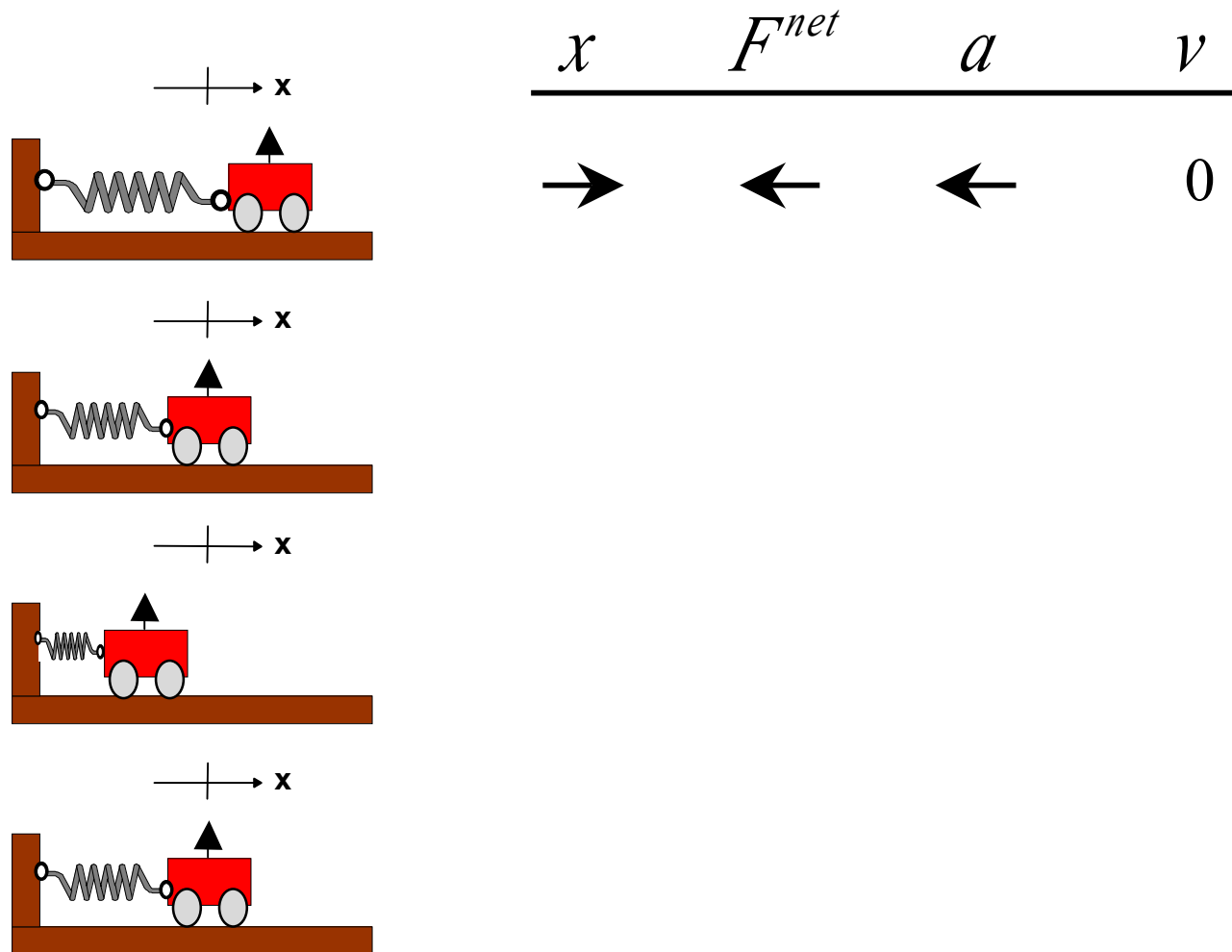
- (i) at some point during the oscillation the mass has  $v = 0$  but  $a \neq 0$
- (ii) at some point during the oscillation the mass has  $v = 0$  and  $a = 0$ .

1. Both occur sometime during the oscillation.
2. Neither occurs during the oscillation.
3. Only (i) occurs.
4. Only (ii) occurs.

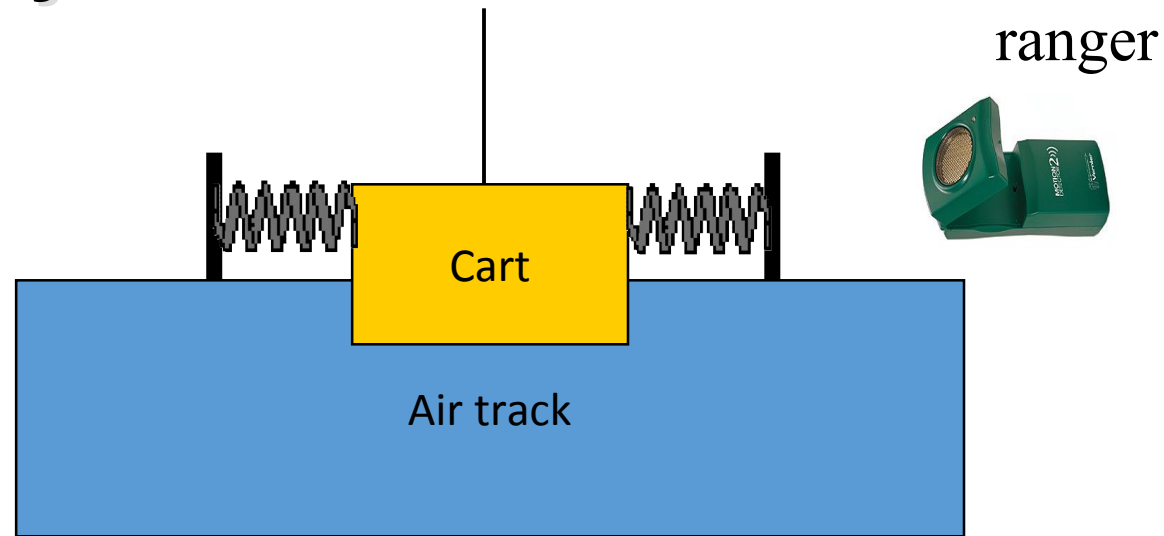




# Tracking the motion

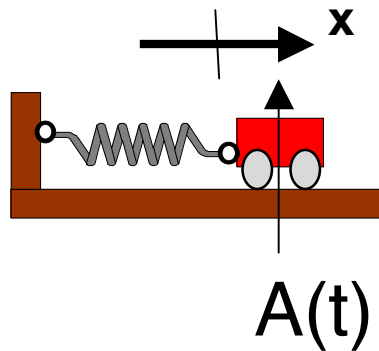


# Let's try it

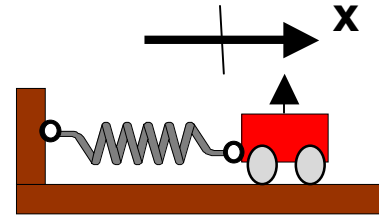


## Why do we have two springs?

- Position of the cart depends on time  $t$
- Lets call the  $x$  position of the cart:  $A(t)$



# Doing the Math: The Equation of Motion



■ Newton's equation for the cart is

$$a = \frac{F_{net}}{m} = \frac{-k x(t)}{m} = -\left(\frac{k}{m}\right)x(t)$$

Acceleration is defined

$$a = \frac{d^2 x(t)}{dt^2}$$

# Solving a differential equation

- Express *acceleration*  $a$  as a derivative of  $x(t)$ .

$$a = \frac{d^2 x(t)}{dt^2} = -\left(\frac{k}{m}\right)x(t)$$

- Verify solution  $x(t) = A \cos(\omega_0(t - t_0))$

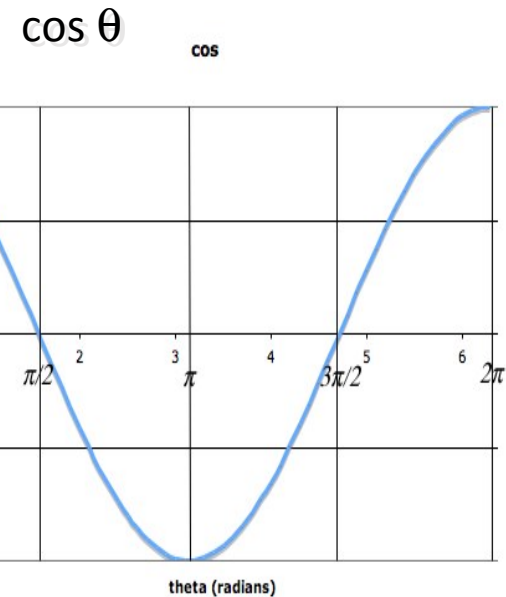
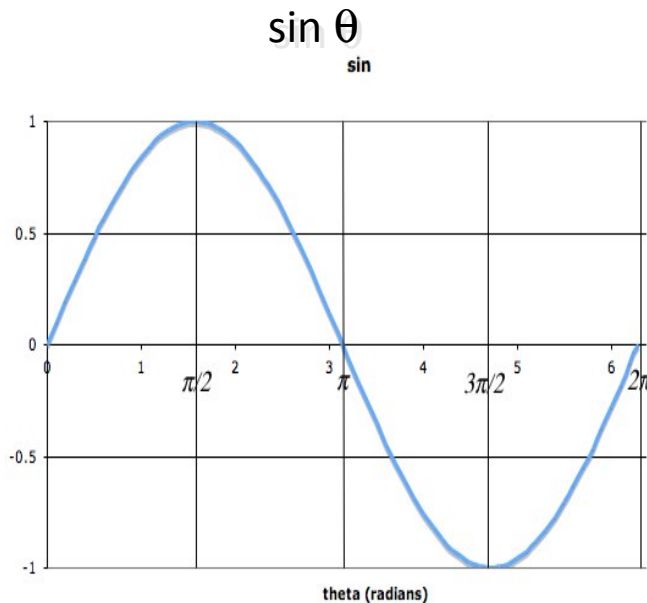
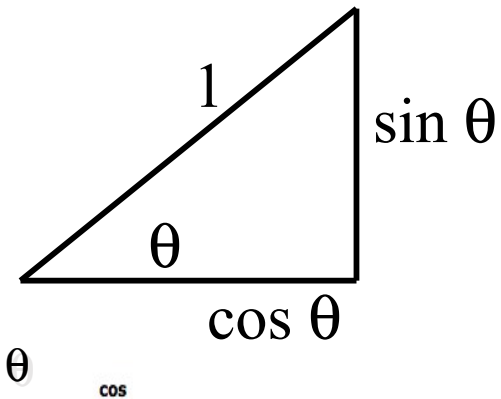
$$\omega_0 = \sqrt{\frac{k}{m}}$$

$$\frac{d^2 \cos \theta}{d\theta^2} = -\cos \theta$$

What determines  $A$  and  $t_0$ ?

# Graphs: $\sin(\theta)$ vs $\cos(\theta)$

- Which is which? How can you tell?
- The two functions sin and cos are derivatives of each other (slopes), but one has a minus sign. Which one? How can you tell?

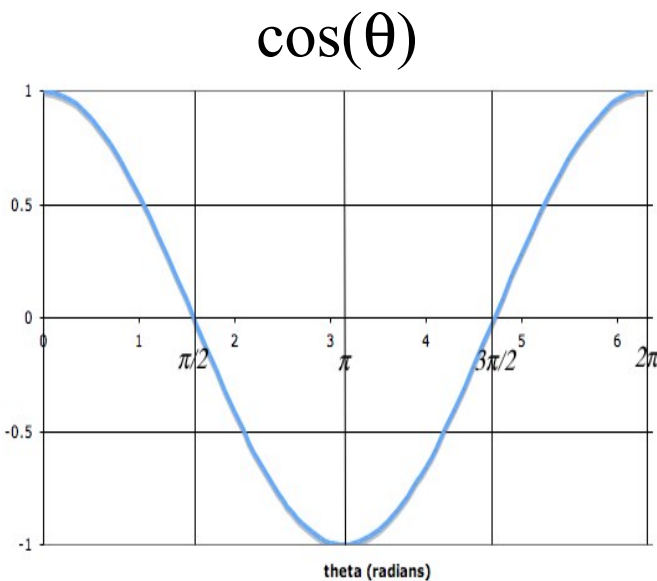


# Graphs: $\sin(\theta)$ vs $\sin(\omega_0 t)$

- For angles,  $\theta = 0$  and  $\theta = 2\pi$  are the same so you only get one cycle.

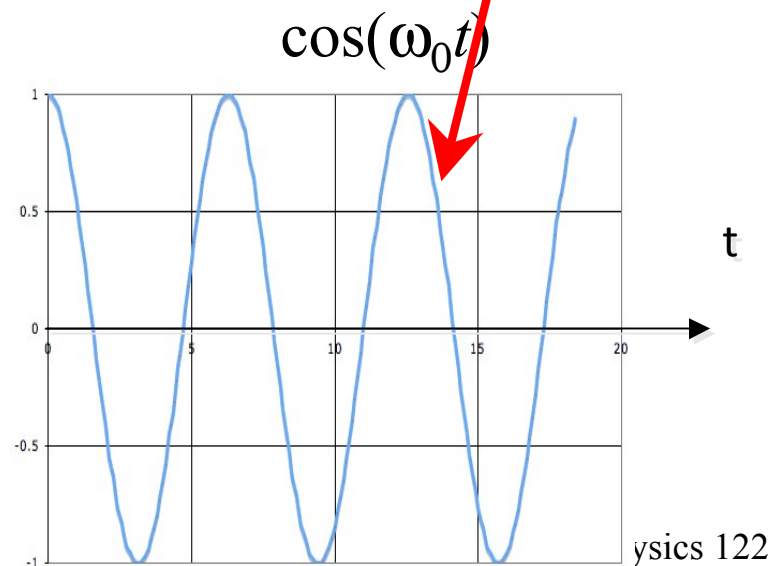
- For time,  $t$  can go on forever so the cycles repeat.

What does changing  $\omega_0$  do to this graph?



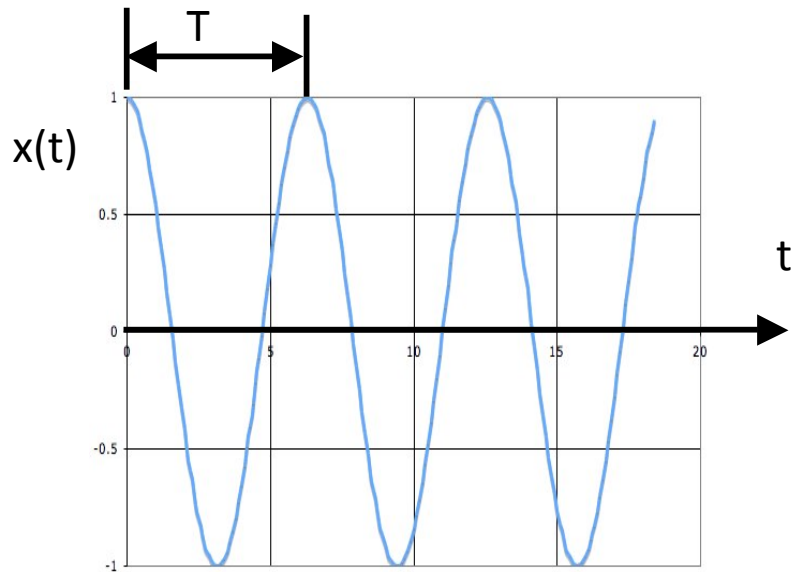
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# Interpreting the Result

- What do the various terms mean?
  - $A$  is the maximum displacement — the *amplitude* of the oscillation.
  - What is  $\omega_0$ ? If  $T$  is the *period* (how long it takes to go through a full oscillation) then

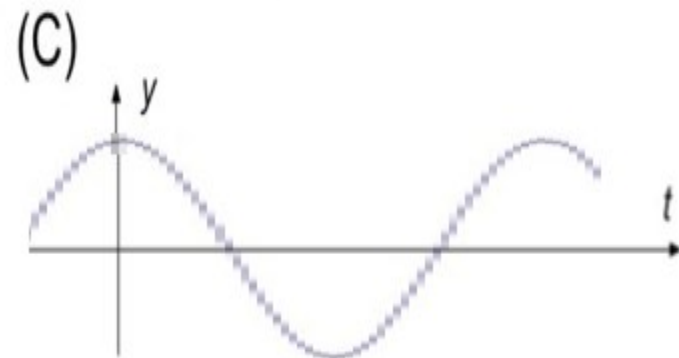
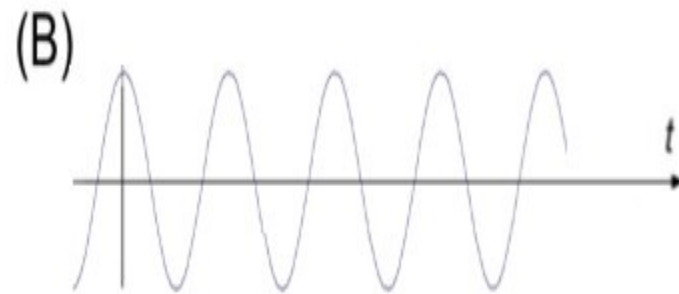
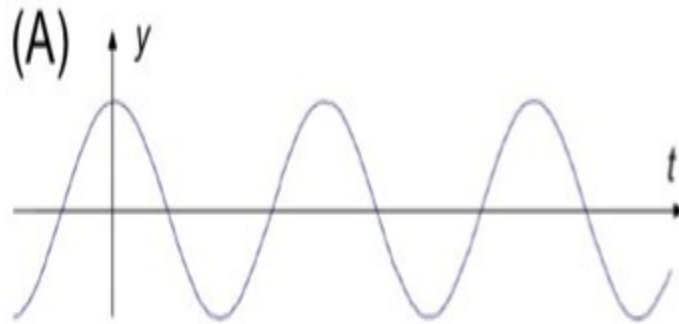


$$\omega_0 t : 0 \rightarrow 2\pi$$

$$t : 0 \rightarrow T$$

$$\omega_0 T = 2\pi \quad \Rightarrow \quad \omega_0 = \frac{2\pi}{T}$$





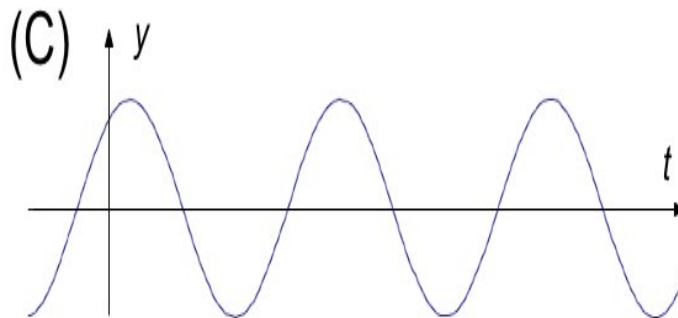
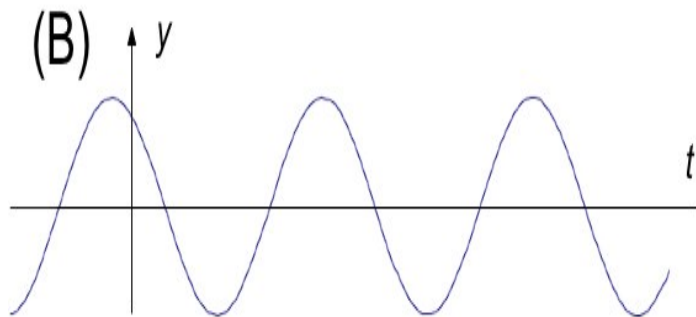
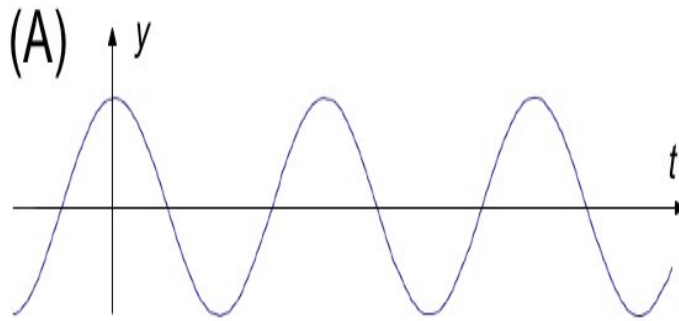
If curve (A) is

$$A \cos(\omega_0 t)$$

which curve is

$$A \cos(2\omega_0 t)?$$

1. (A)
2. (B)
3. (C)
4. None of the above.



Which of these curves is described by

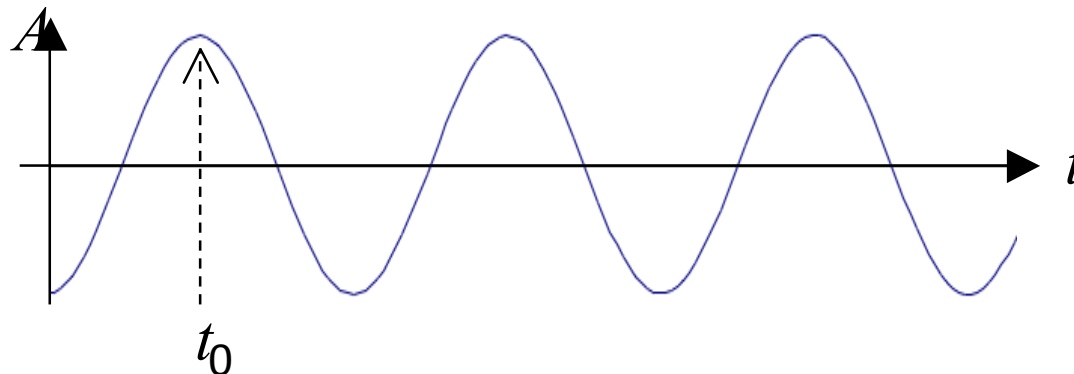
$$A \cos(\omega_0 t + \phi)$$

with  $\phi > 0$  (and  $\phi \ll 2\pi$ )?

1. (A)
2. (B)
3. (C)
4. None of the above.

# Oscillations

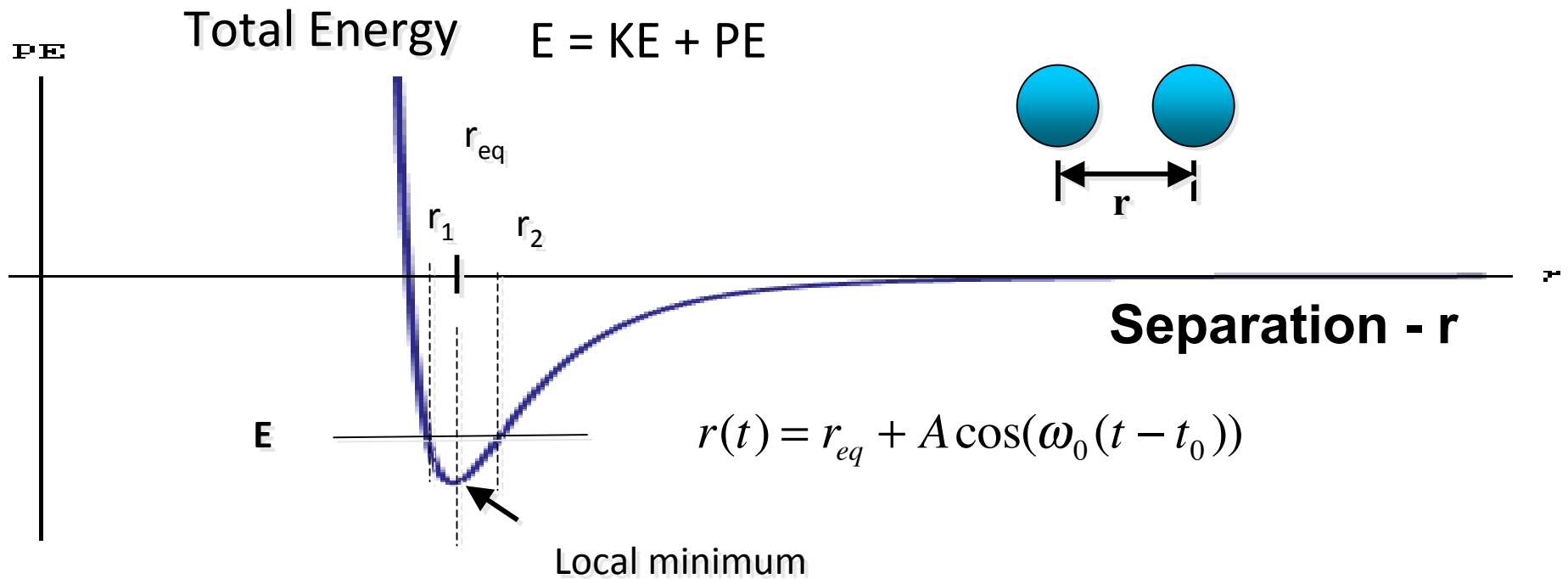
- A typical oscillation



$$x(t) = A \cos(\omega_0(t - t_0))$$

$$A(t) = A_0 \cos(\omega_0(t - t_0))$$

$$= A_0 \cos(\omega_0 t - \omega_0 t_0) = A_0 \cos(\omega_0 t - \phi)$$



Potential  $PE = U$

$$U(r) \simeq \frac{k}{2}(r - r_{eq})^2 + U(r_{eq})$$

$$U = \frac{k}{2}A^2 \cos^2(\omega(t - t_0)) + U(r_{eq})$$

Kinetic

$$KE = \frac{m}{2}(v(t))^2 = \frac{m}{2}\left(\frac{dr(t)}{dt}\right)^2$$

$$KE = \frac{m\omega_0^2}{2}A^2 \sin^2(\omega(t - t_0))$$

add them together

$$\omega_0^2 = \frac{k}{m}$$



$$\begin{aligned} E &= U(r_{eq}) + \frac{k}{2} A^2 \cos^2(\omega(t - t_0)) + \frac{m\omega_0^2}{2} A^2 \cos^2(\omega(t - t_0)) \\ &= U(r_{eq}) + \frac{k}{2} A^2 [\cos^2(\omega(t - t_0)) + \sin^2(\omega(t - t_0))] \\ &= U(r_{eq}) + \frac{k}{2} A^2 \end{aligned}$$

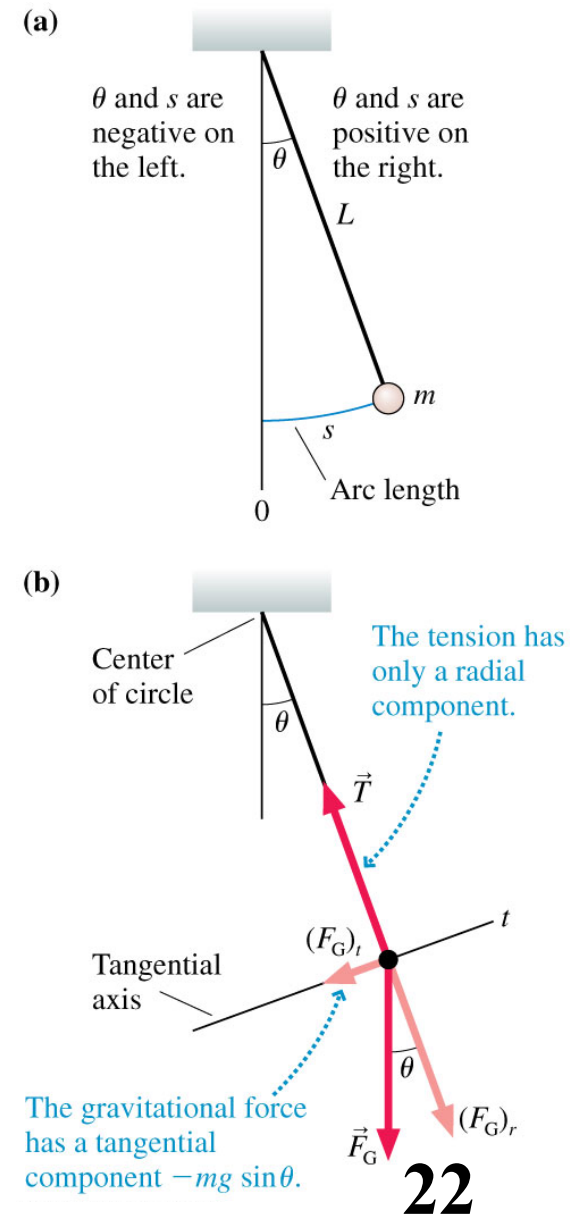
Total energy is constant

# The Simple Pendulum

- Consider a mass  $m$  attached to a string of length  $L$  which is free to swing back and forth.
- If it is displaced from its lowest position by an angle  $\theta$ , Newton's second law for the tangential component of gravity, parallel to the motion, is:

$$(F_{\text{net}})_t = \sum F_t = (F_G)_t = -mg \sin \theta = ma_t$$

$$\frac{d^2 s}{dt^2} = -g \sin \theta$$



# The Simple Pendulum



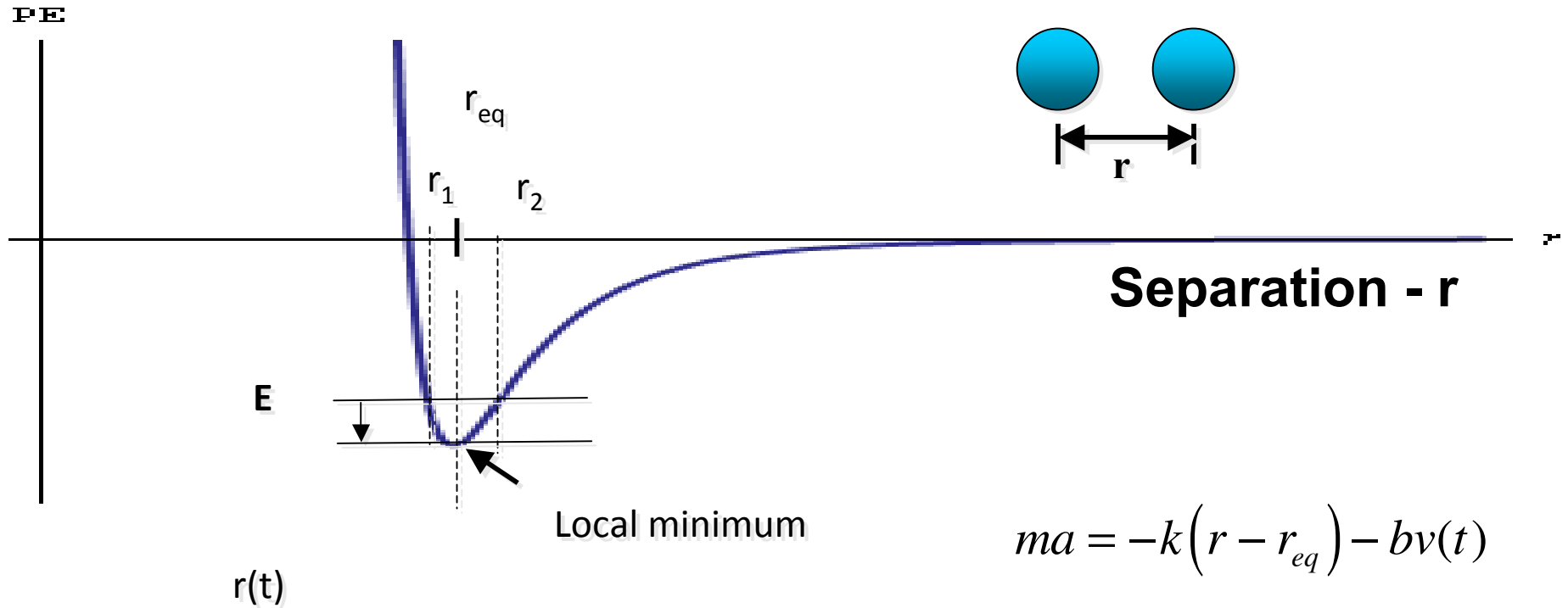
If we restrict the pendulum's oscillations to small angles ( $< 10^\circ$ ), then we may use the **small angle approximation**  $\sin \theta \approx \theta$ , where  $\theta$  is measured in radians.

$$(F_{\text{net}})_t = -mg \sin \theta \approx -mg\theta = -\frac{mg}{L}s$$

and the angular frequency of the motion is found to be:

$$\omega = 2\pi f = \sqrt{\frac{g}{L}}$$

Now add friction



$$ma = -k(r - r_{eq}) - bv(t)$$

$$r(t) = r_{eq} + A \exp\left[-\frac{\gamma}{2}(t - t_0)\right] \cos(\omega_1(t - t_0))$$

$$\frac{d^2}{dt^2} r(t) = -\omega_0^2 r(t) - \gamma \frac{dr}{dt}$$

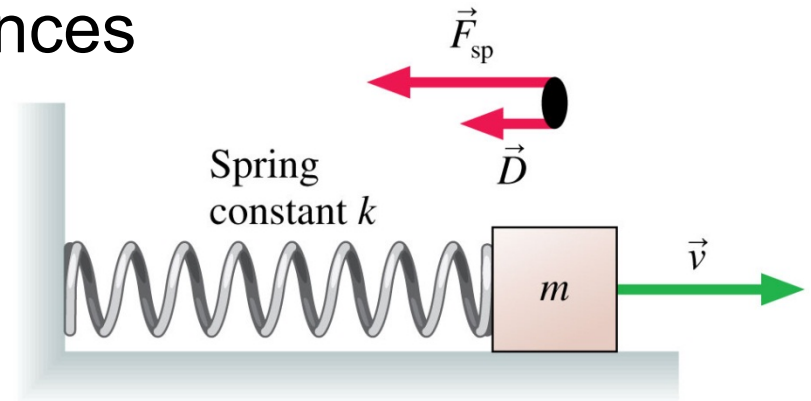
$$\omega_1^2 = \omega_0^2 - \frac{\gamma^2}{4}$$

$$\gamma = \frac{b}{m}$$



# Damped Oscillations

When a mass on a spring experiences the force of the spring as given by Hooke's Law, as well as a linear drag force of magnitude  $|D| = bv$ , the solution is:



$$x(t) = Ae^{-bt/2m} \cos(\omega t + \phi_0) \quad (\text{damped oscillator})$$

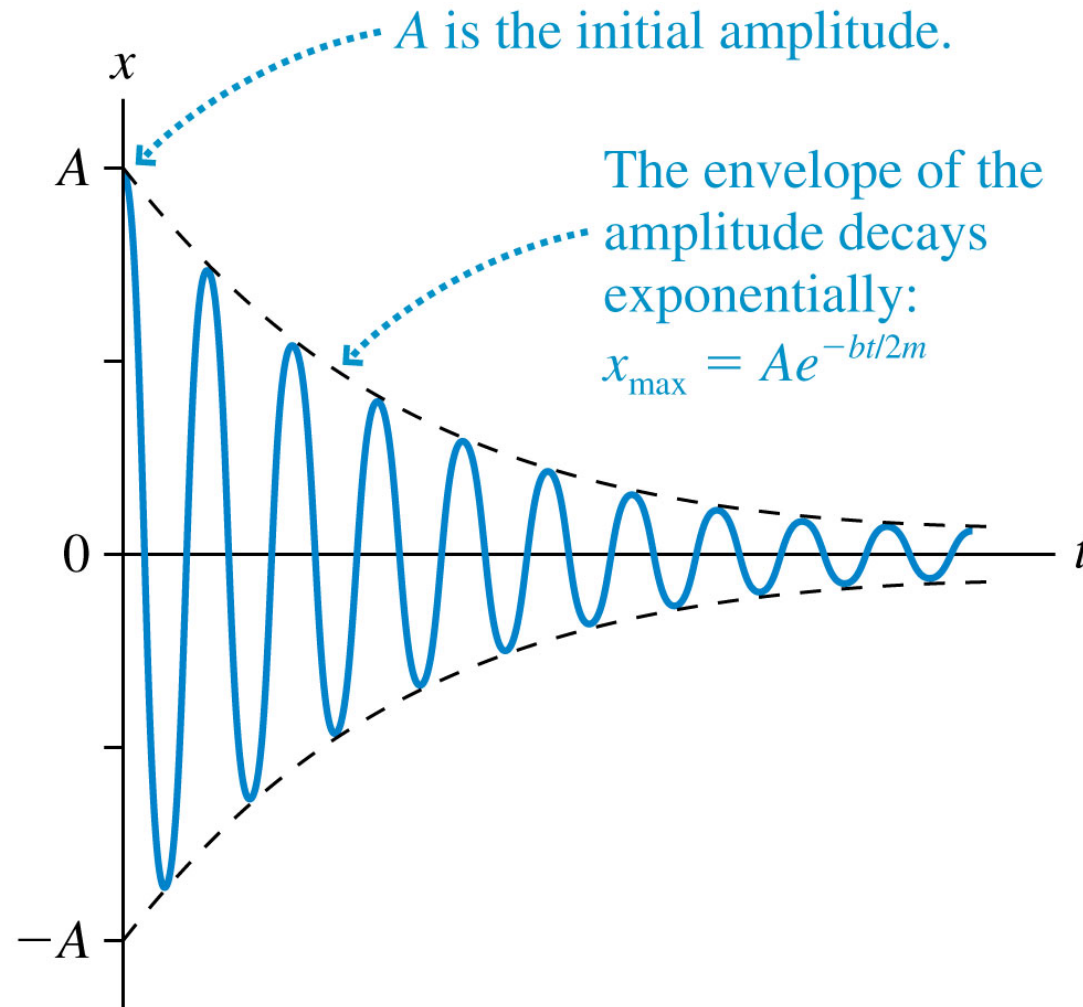
where the angular frequency is given by:

$$\omega = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}} = \sqrt{\omega_0^2 - \frac{b^2}{4m^2}}$$

Here  $\omega_0 = \sqrt{k/m}$  is the angular frequency of the undamped oscillator ( $b = 0$ ).

# Damped Oscillations

Position-versus-time graph for a damped oscillator.

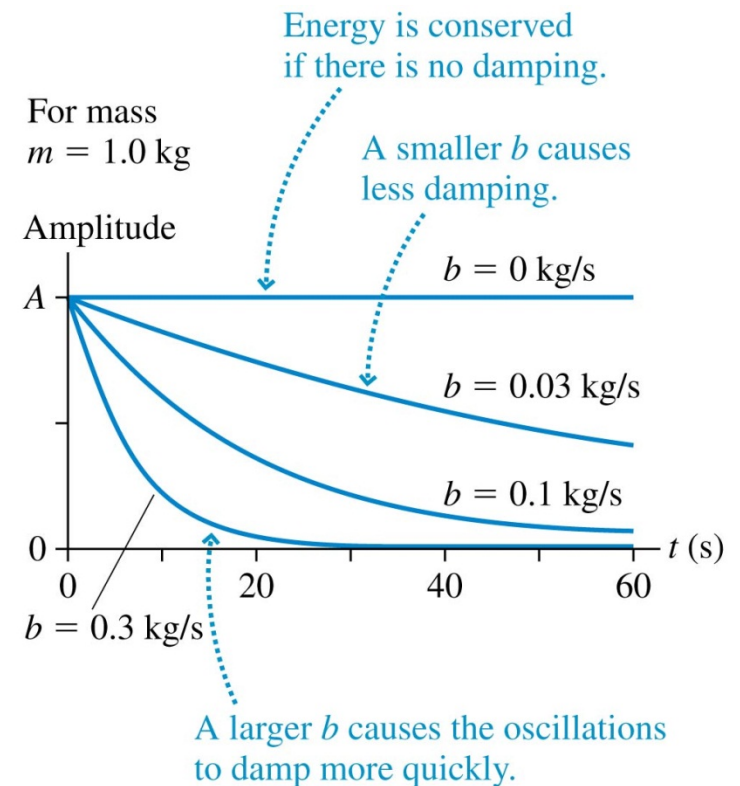


# Damped Oscillations

- A damped oscillator has position  $x = x_{\max} \cos(\omega t + \phi_0)$ , where:

$$x_{\max}(t) = Ae^{-bt/2m}$$

- This slowly changing function  $x_{\max}$  provides a border to the rapid oscillations, and is called the **envelope**.
- The figure shows several oscillation envelopes, corresponding to different values of the damping constant  $b$ .



# Energy in Damped Systems

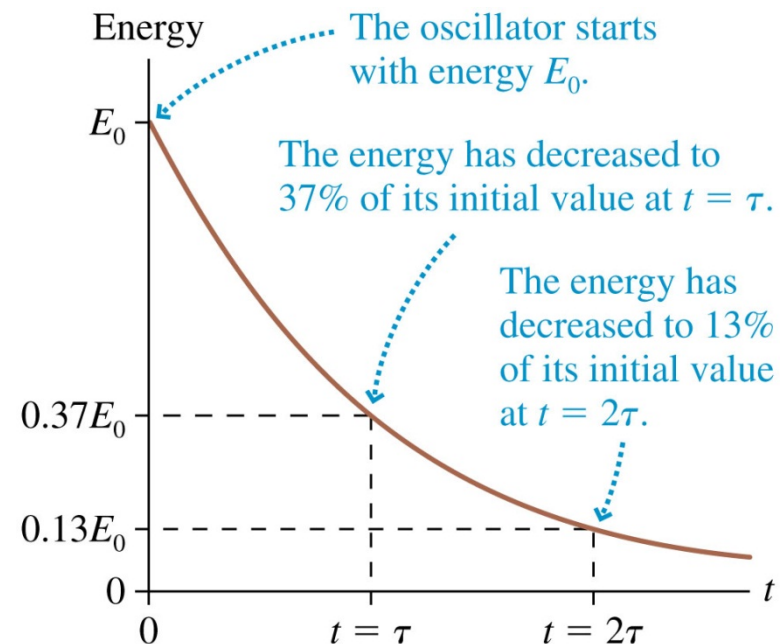
- Because of the drag force, the mechanical energy of a damped system is no longer conserved.
- At any particular time we can compute the mechanical energy from:

$$E(t) = \frac{1}{2}k(x_{\max})^2 = \frac{1}{2}k(Ae^{-t/2\tau})^2 = \left(\frac{1}{2}kA^2\right)e^{-t/\tau} = E_0e^{-t/\tau}$$

- Where the decay constant of this function is called the **time constant**  $\tau$ , defined as:

$$\tau = \frac{m}{b}$$

- The oscillator's mechanical energy decays exponentially with time constant  $\tau$ .



# Driven Oscillations and Resonance

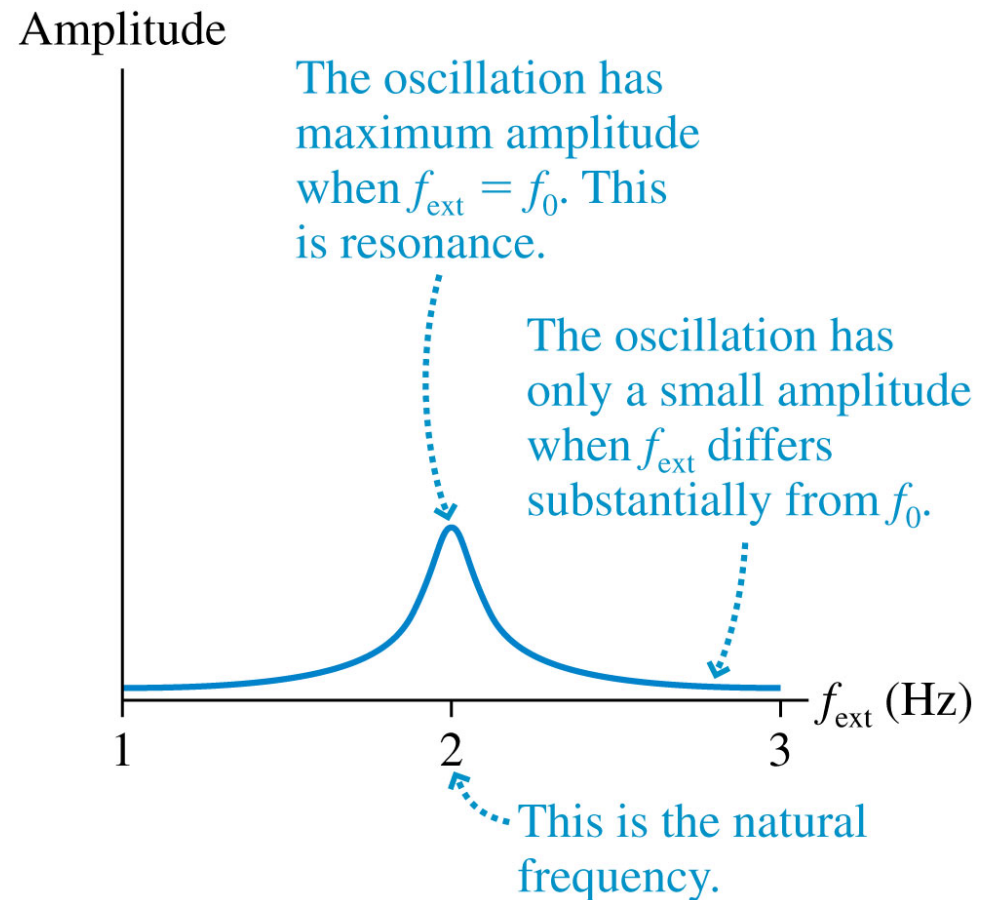
- Consider an oscillating system that, when left to itself, oscillates at a **natural frequency**  $f_0$ .
- Suppose that this system is subjected to a *periodic* external force of **driving frequency**  $f_{\text{ext}}$ .
- The amplitude of oscillations is generally not very high if  $f_{\text{ext}}$  differs much from  $f_0$ .
- As  $f_{\text{ext}}$  gets closer and closer to  $f_0$ , the amplitude of the oscillation rises dramatically.



A singer or musical instrument can shatter a crystal goblet by matching the goblet's natural oscillation frequency.

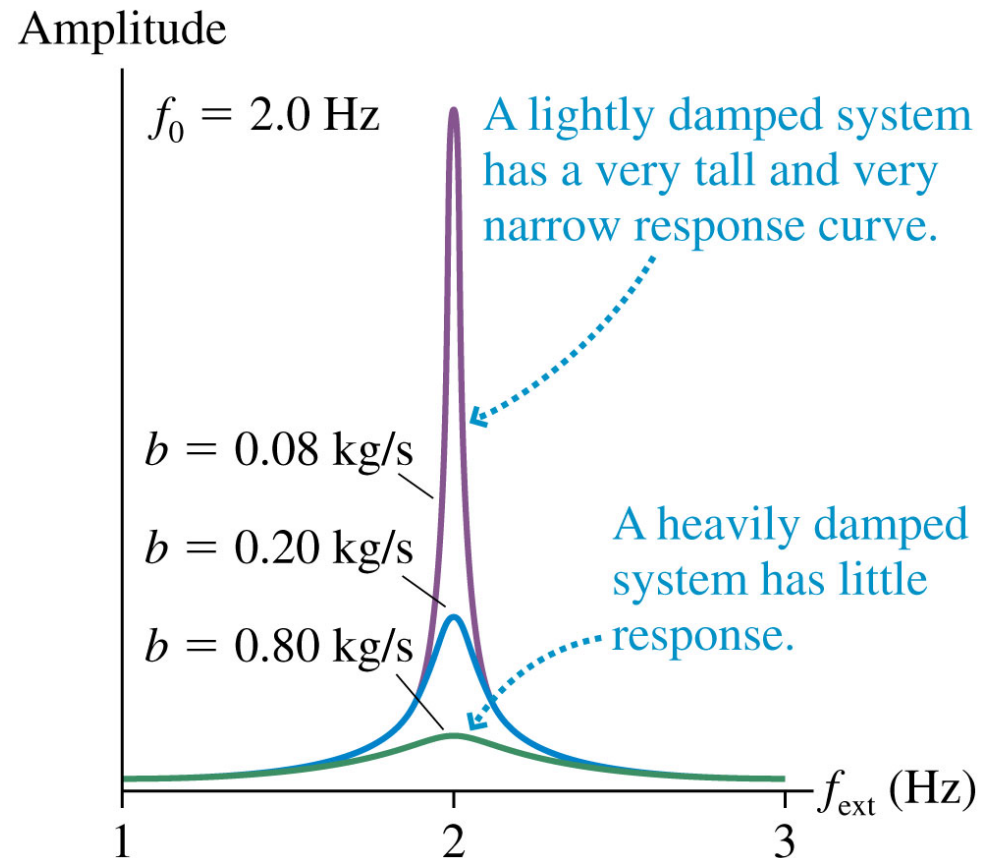
# Driven Oscillations and Resonance

The response curve shows the amplitude of a driven oscillator at frequencies near its natural frequency of 2.0 Hz.



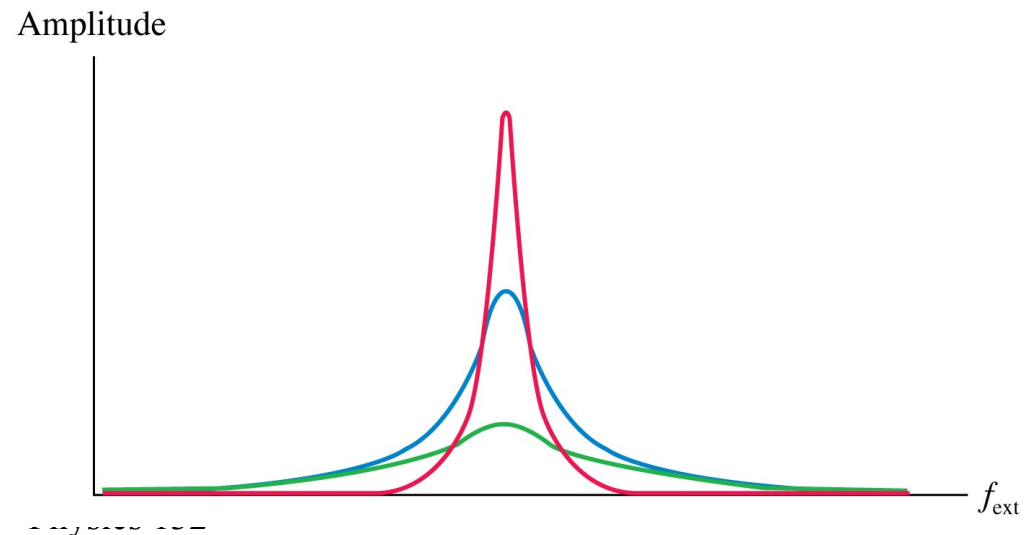
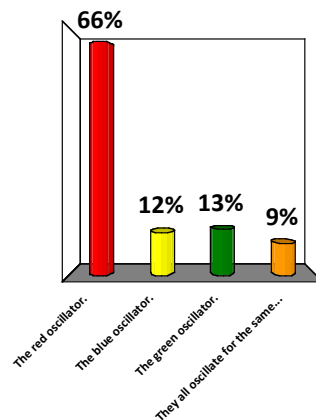
# Driven Oscillations and Resonance

- The figure shows the same oscillator with three different values of the damping constant.
- The resonance amplitude becomes higher and narrower as the damping constant decreases.



The graph shows how three oscillators respond as the frequency of a driving force is varied. If each oscillator is started and then left alone, which will oscillate for the longest time?

- 😊 A. The red oscillator.
- B. The blue oscillator.
- C. The green oscillator.
- D. They all oscillate for the same length of time.







- <https://www.youtube.com/watch?v=x9BVSu7Ok>