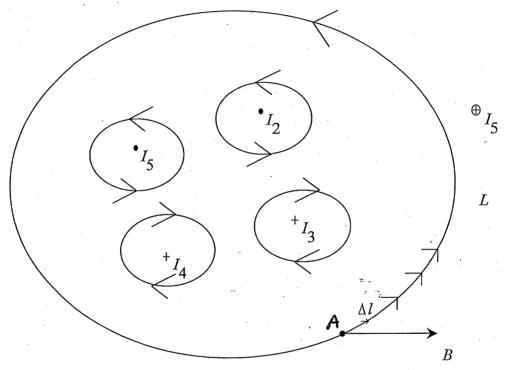
Ampere's Law - Applications

 $\frac{B \text{ along } \Delta l \text{ multiplied by } \Delta l}{\beta \cdot \Delta l} = B \Delta l \cos (\beta, \Delta l).$

If Blue, Brae=0.



Repeat this calculation at every step as shown. Both $B \rightarrow A l$

Write out the sum

$$\Sigma_c \xrightarrow{B \sim \Lambda l};$$

c: closed loop.

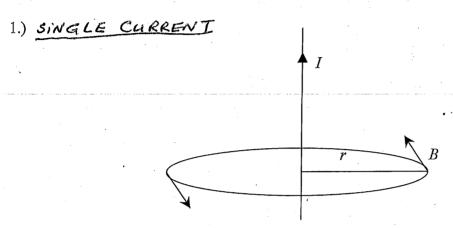
This sum is called circulation of B around closed loop and Ampere says that it is determined solely by currents threading through the surface on which the loop is drawn & only currents within the loop contribute, i.e. exclude I_5 . The mathematical Equation is:

$$\sum_{c} \underbrace{B \Delta l}_{i} = \mu_{0} \sum_{i} I_{i}, \mu_{0} = 4\pi \times 10^{-7} \frac{T - m}{A}$$

In words, circulation of B around closed loop is proportional to the algebraic sum of the currents threading the best suchace on which we hope is drawn.

Note: As in Case of Gauss' Law, Amperels law gives Circulation but not B. To get B you need high symmetry!

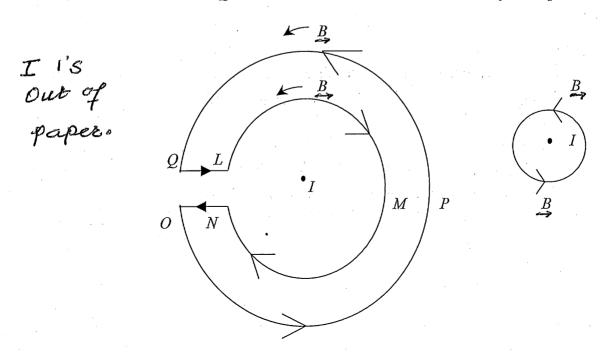
Applications



Single wire with current I, there is cylindrical symmetry so B can be a function of r only & must encircle I. E and A are parallel to one another I. Appropriate loop is circle of radius r centered on the wire

$$B \cdot 2\pi r = \mu_0 I$$
so, $B = \frac{\mu_0 I}{2\pi r} \hat{\varnothing}$
as claimed previously

2.) Next, we begin by showing that if current is outside the loop it contributes nothing to the circulation. Choose LMNOPQ with I at the center of the circles of radii r_1 , and r_2



$$B(r_1) = \frac{\mu_0 I}{2\pi r_1} \hat{\varnothing}$$

$$B(r_2) = \frac{\mu_0 I}{2\pi r_2} \hat{\varnothing}$$

$$\sum_{c \to \infty} B \cdot \Delta I = \frac{\mu_0 I}{2\pi r_1} \cdot 2\pi r_1 + O + \frac{\mu_0 I}{2\pi r_2} \cdot 2\pi r_2 + O$$

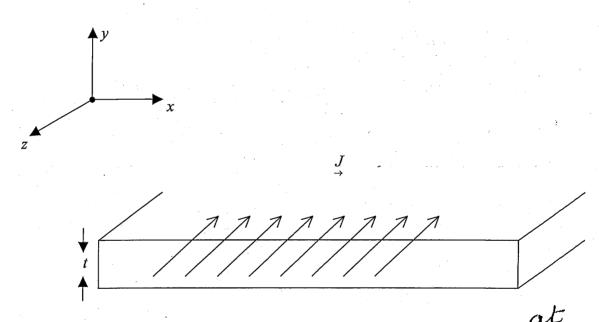
$$= L \to M \to N + N \to 0 + O \to P \to Q + Q \to L$$

$$= O$$

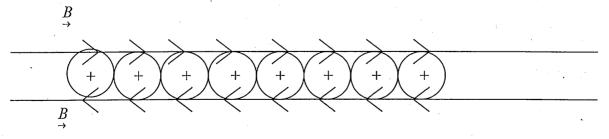
[The first lerm is -ive]

B and Al are
opposite to one
another

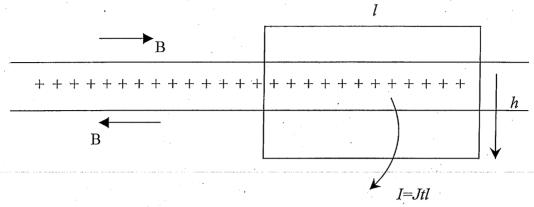
3.)



At y=0 there is a Current sheet of thickness t carrying current density $J=-J\hat{z}$. Looking it endon we see sources as



and we see that y-components of B cancel out. $B || \hat{x}$ survives. Let us take loop of width l and height h.



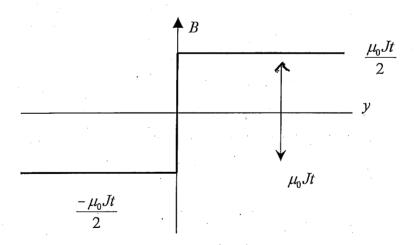
$$\begin{split} \Sigma_c & \underset{\rightarrow}{B} \text{,} \Delta l = Bl + 0 + Bl + 0 \\ & = 2Bl \\ & = \mu_0 Jtl \end{split}$$

so,

$$B = \frac{\mu_0 Jt}{2} \& B = \frac{\mu_0 Jt}{2} \hat{x} \qquad y > 0$$

$$= \frac{-\mu_0 Jt}{2} \hat{x} \qquad y > 0$$

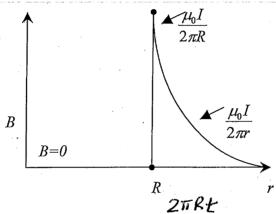
That is, B-field will jump by $\mu_0 Jt$ on crossing the current sheet from y<0 to y>0.

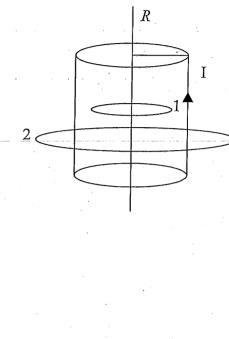


4.) <u>Hollow Cylindrical Conductor</u>- Radius R, carries uniform current. We want B at a distance r from its axis. Since there is a cylindrical symmetry we should use circles centered on the axis for our closed loop. For r < R. use loop 1.

 $B \cdot 2\pi r = 0$. No Current threads through loop 1. $F = \mathcal{F} \times \mathcal{F}$ $B \cdot 2\pi r = \mu_0 I$, the entire current threads loop 2.

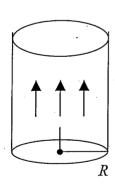
so,
$$B = \frac{\mu_0 I}{2\pi r} \hat{\varnothing}$$





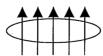
Note: if cylinder has wall thickness t: $I = J \cdot 2$ and field at surface would be $\mu_0 Jt$. Again field would jump by $\mu_0 Jt$ on crossing a current sheet.

5.) SOLID CYLINDRICAL CONDUCTOR - with uniform current



Define
$$J = \frac{I}{\pi R^2}$$

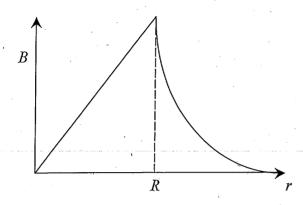
Now for r < R $I = J\pi r^2$



$$B \cdot 2\pi r = \mu_0 J \pi r^2$$

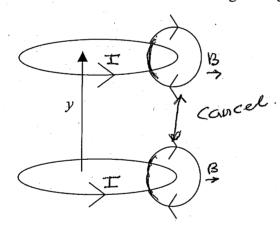
$$B = \frac{\mu_0 J r}{2} \hat{\varnothing} \qquad r < R$$

For >R, entire I contributes $B = \frac{\mu_0 I}{2\pi r} \hat{\varnothing}$



6.) Solenoid: Tightly wound, small radius, length much larger than radius:

N turns, L long, $n = \frac{N}{L}$ =# of turns per meter. Look at two neighboring turns



r-component cancels B_{ν} inside survives.

Long-narrow solenoid. B field lives must $B \approx 0$ just outside. Come out of top, loop around and $Bl = \mu_0 n I l$ enter at $B = \mu_0 n I l$ bottom with mo breaks or bench allowed.

