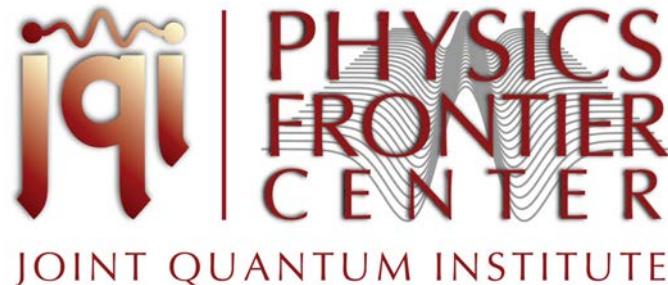


Optical Nanofibers; some experiments in optomechanics.

LENS, Florence November 2019

Luis A. Orozco

www.jqi.umd.edu



Work supported by: National Science Foundation of the USA, the Physics Frontier Center at the Joint Quantum Institute, The Joint Quantum Institute, and The University of Shanxi, China.

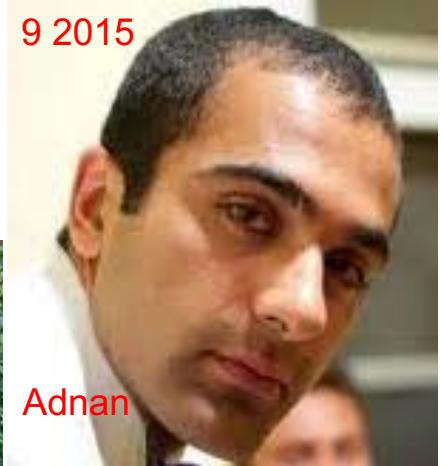
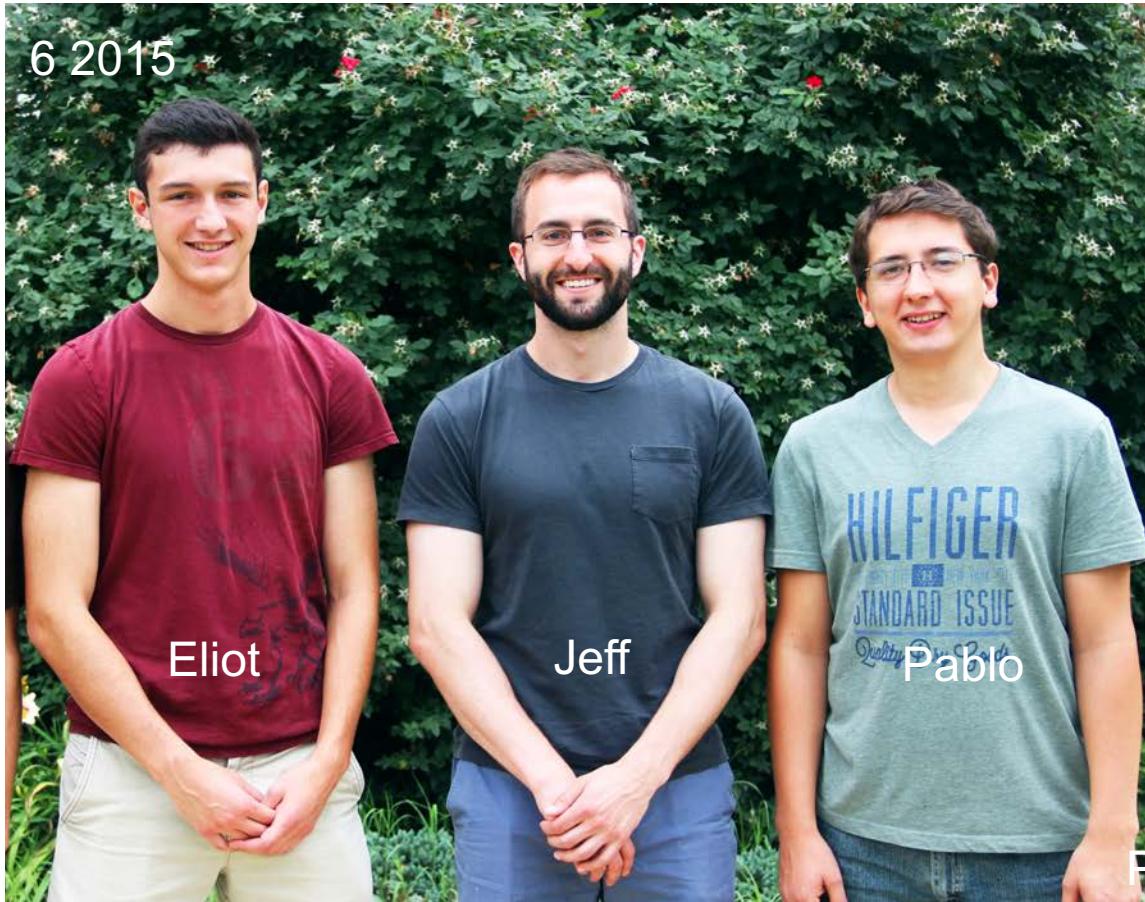
Graduate students: Jeff A. Grover, Pablo Solano,
Dianquiang Su.

Undergraduate students: Eliot Fenton, Adnan Kahn.

Professors and Researchers: Pablo Barberis Blostein,
Fredrik K. Fatemi, Luis A. Orozco, Yanting Zhao.

University of Maryland, Army Research Laboratory,
Universidad Nacional Autónoma de México, Shanxi
University.

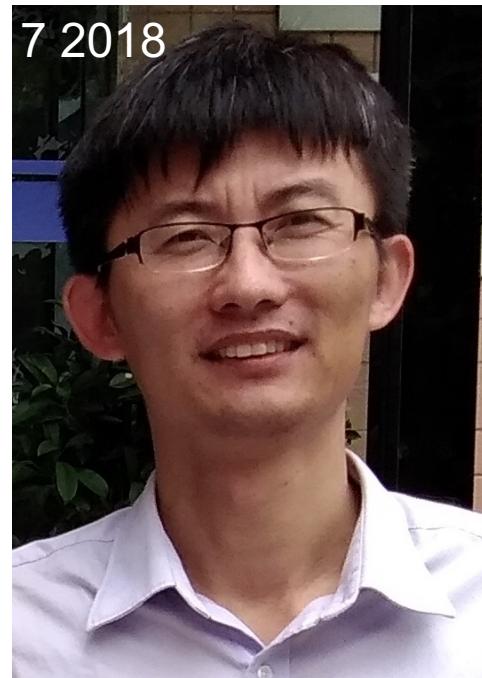
JQI students



Shanxi University, Taiyuan, China



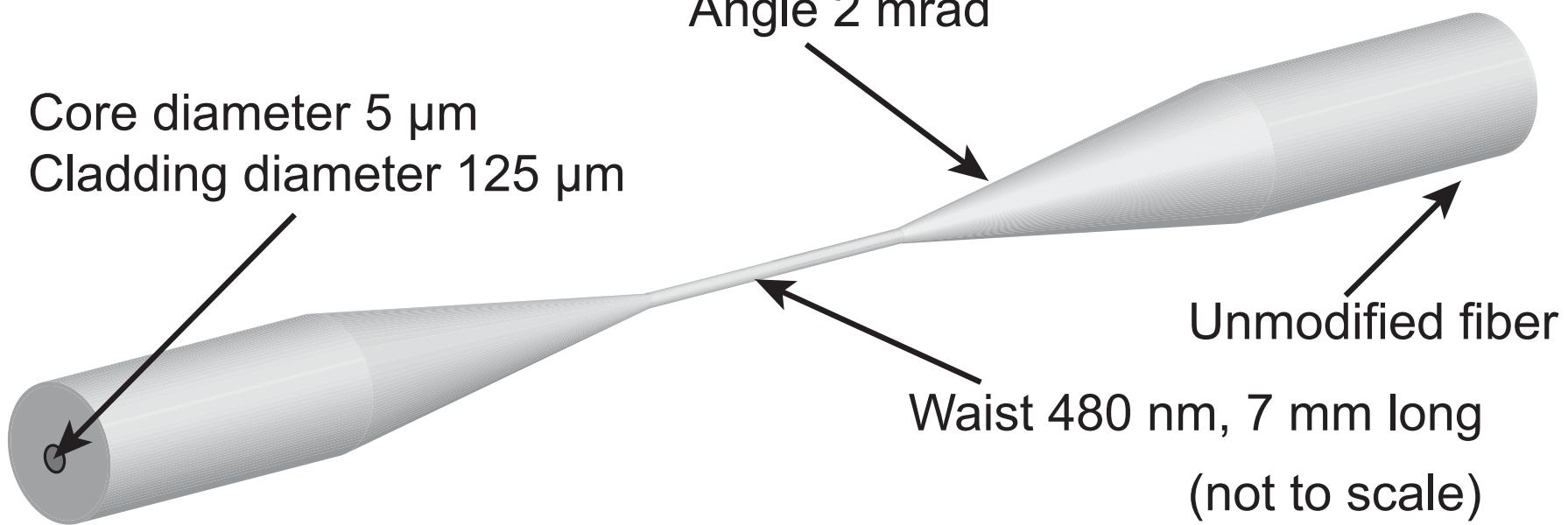
Su Dianquiang



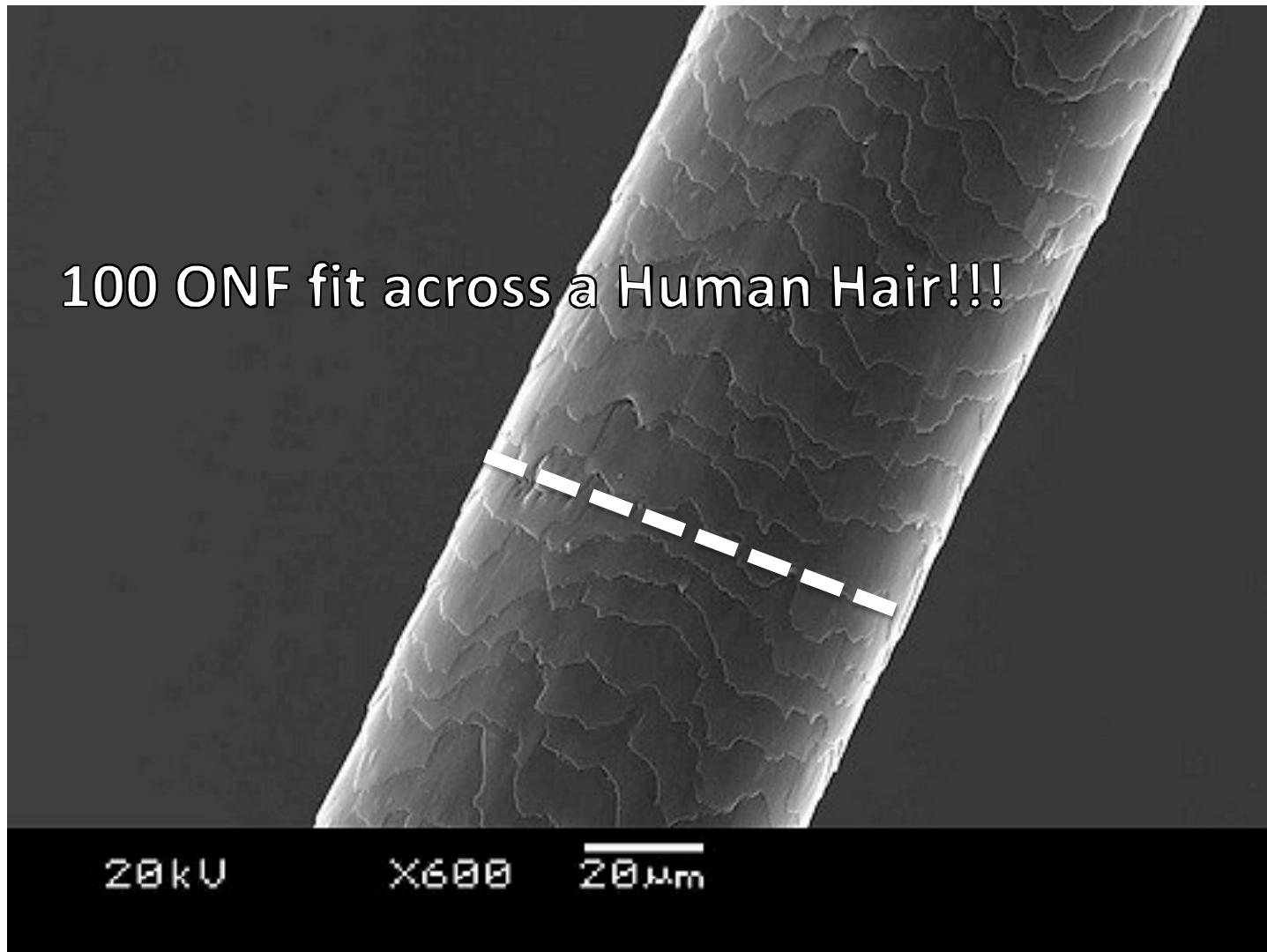
Zhao Yanting

Optical Nanofibers

Optical Nanofibers

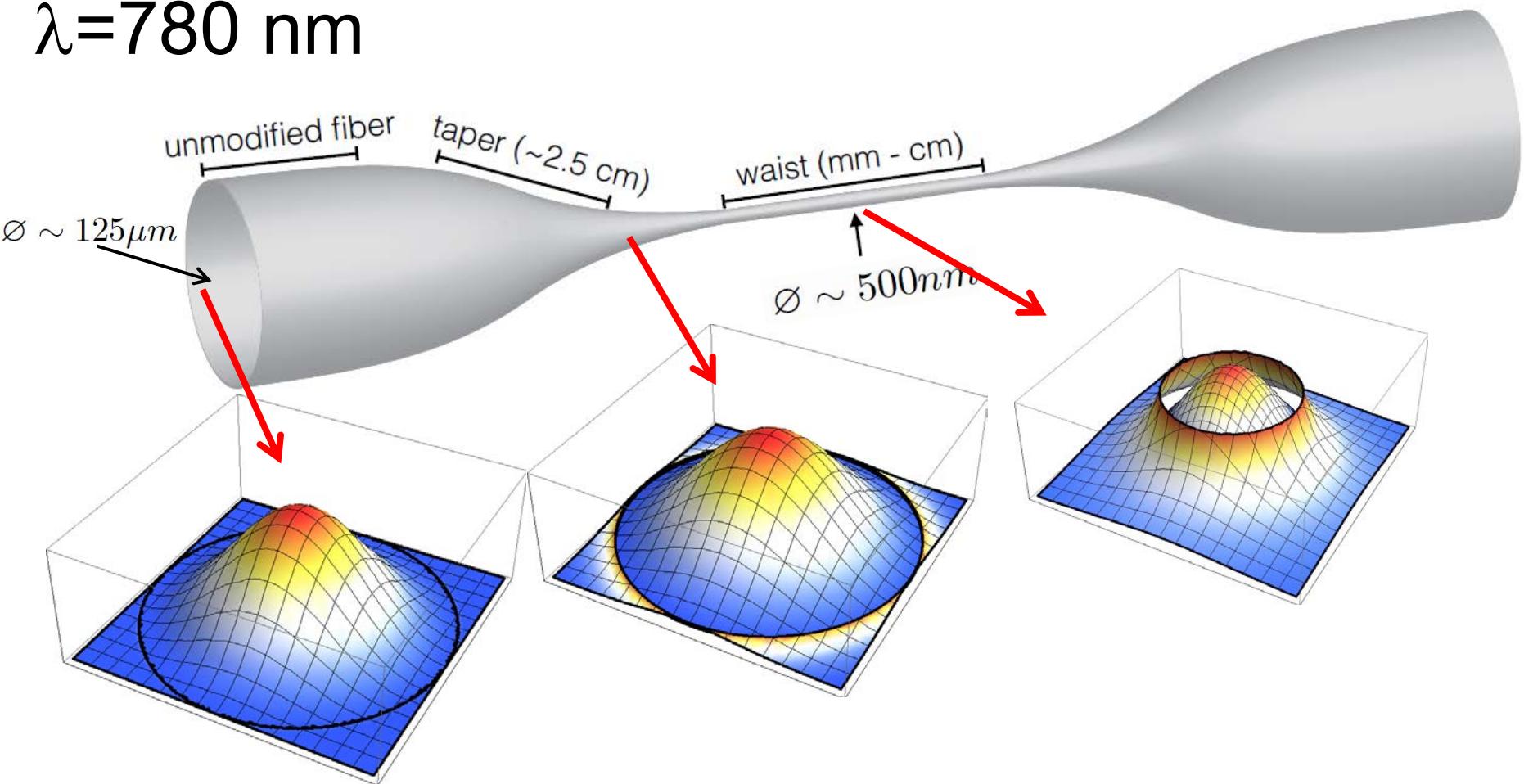


The scale

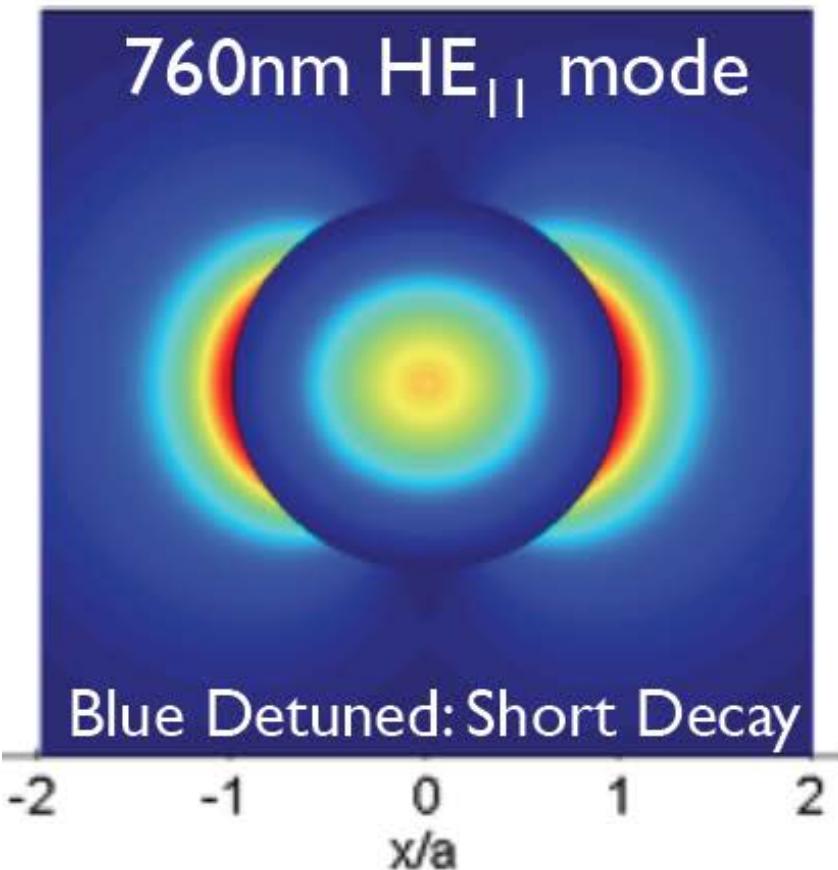


Optical Nanofibers

$\lambda=780 \text{ nm}$



Fundamental mode on an optical nanofiber with vertical linear polarization.

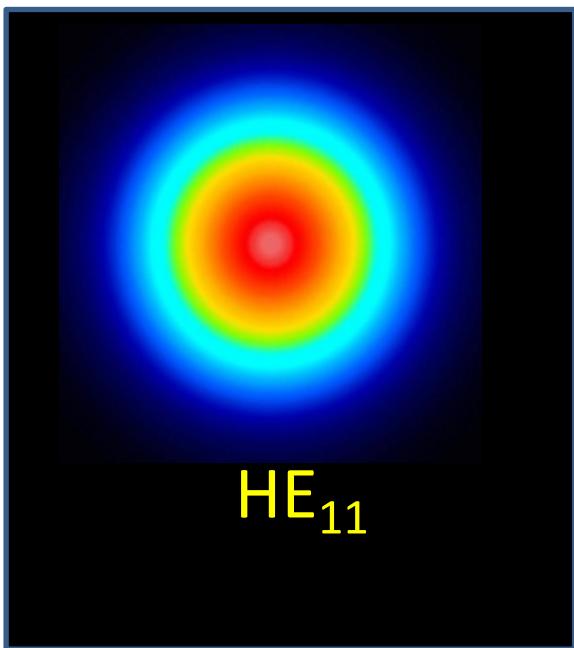


- Radius ~ 250 nm
- Decay length: ~ 100 nm
- Intensity 1 mW in the evanescent field 5×10^8 mW/cm² = $10^8 I_{\text{sat}}$ on the D2 line of Rb.

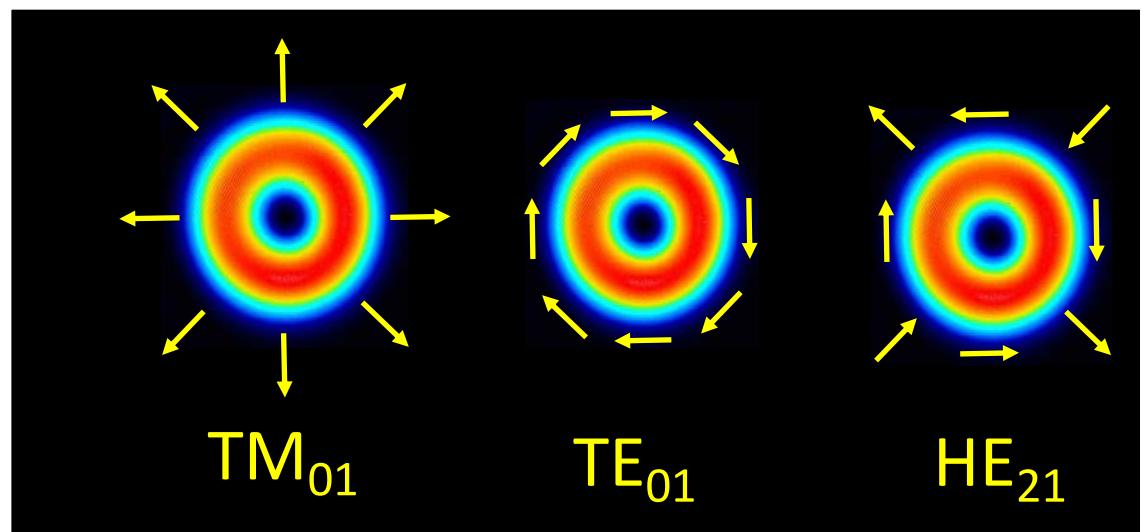
Modes and polarization properties

Lowest order fiber modes

Intensities and polarizations



HE_{11}

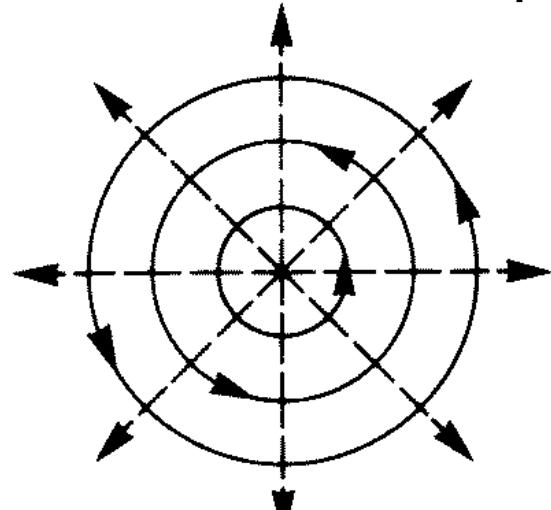


TM_{01}

TE_{01}

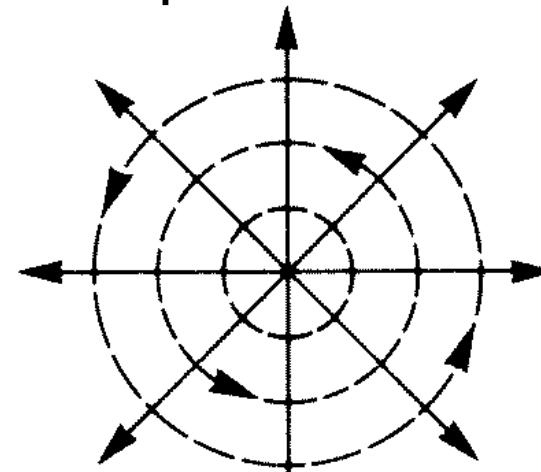
HE_{21}

Transversal component of the polarizations



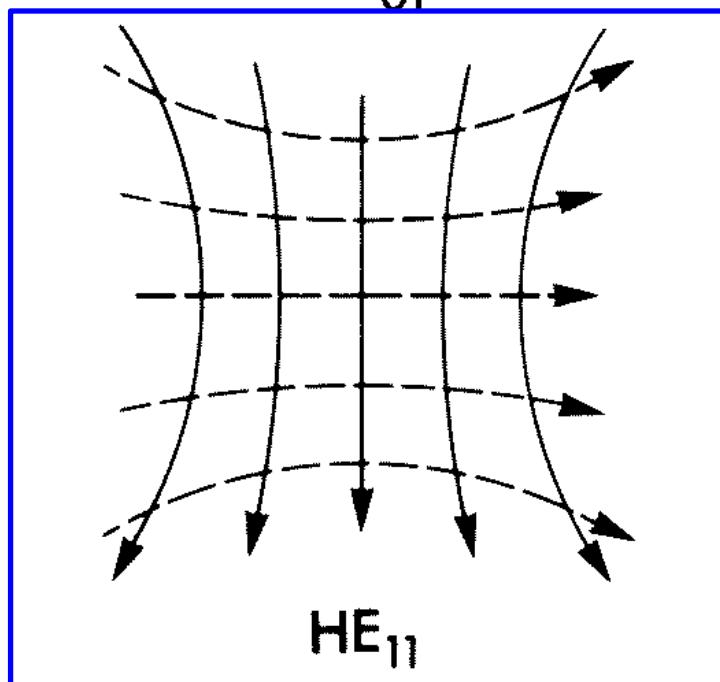
E continuous line

TE₀₁

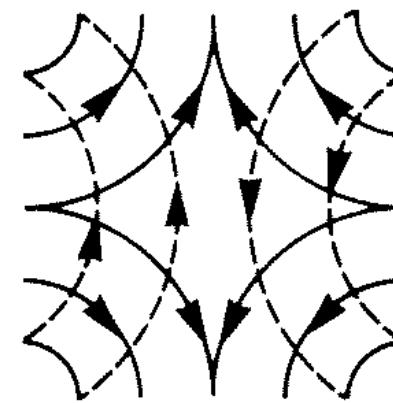


B dashed line

TM₀₁



HE₁₁



HE₂₁

The gradients of E in the radial direction
are large.

- $\text{Div } \mathbf{E}=0$ (Ley de Gauss)
implies large longitudinal
components.

$$\frac{\partial E_r}{\partial r} + \frac{\partial E_z}{\partial z} = 0$$

- The evanescent field has to
have a longitudinal component
to compensate the radial
gradient.

Polarization at the waist of the nanofiber

$$\nabla \cdot \vec{E} = 0$$

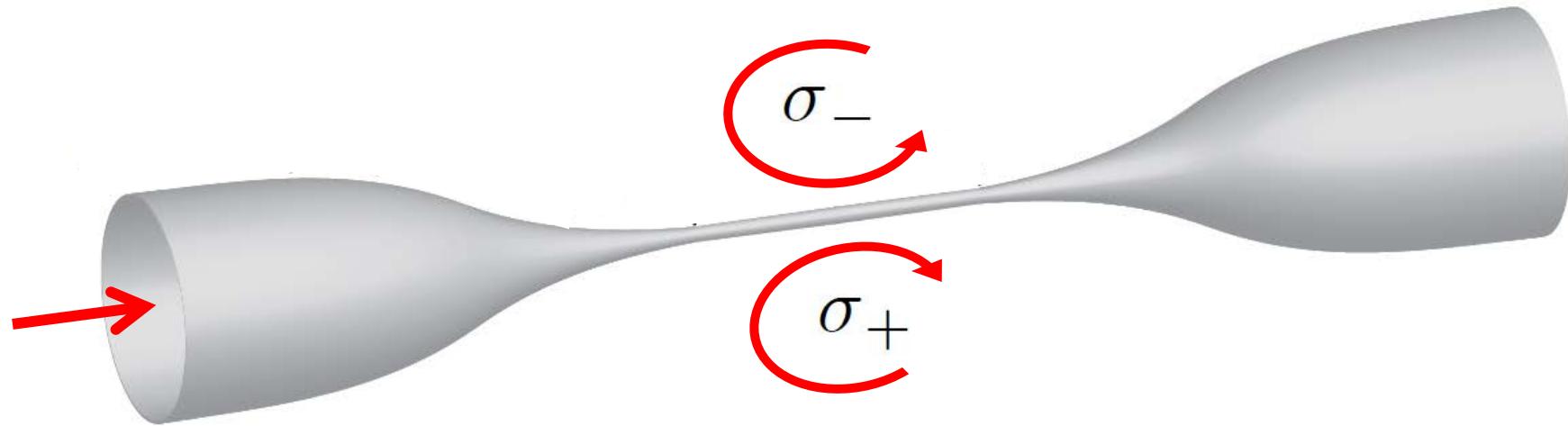
$$\nabla_T \cdot \vec{E} + \frac{2\pi}{\lambda} i E_Z = 0$$



Polarization at the waist of the nanofiber

$$\nabla \cdot \vec{E} = 0$$

$$\nabla_T \cdot \vec{E} + \frac{2\pi}{\lambda} i E_Z = 0$$

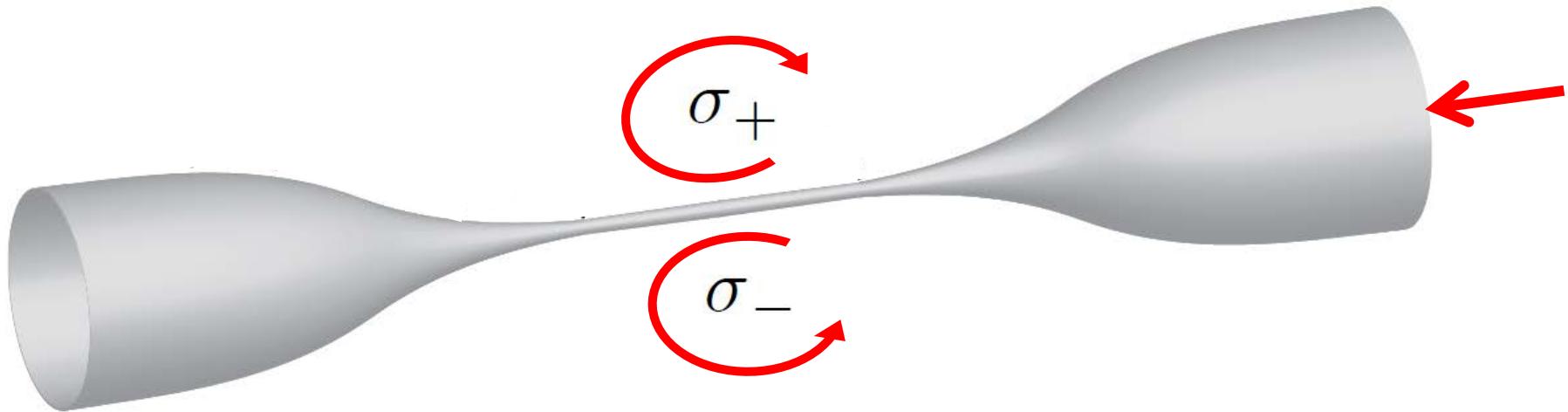


Bicycle circular polarization, not propeller

Polarization at the waist of the nanofiber

$$\nabla \cdot \vec{E} = 0$$

$$\nabla_T \cdot \vec{E} + \frac{2\pi}{\lambda} i E_Z = 0$$



Bicycle circular polarization, not propeller

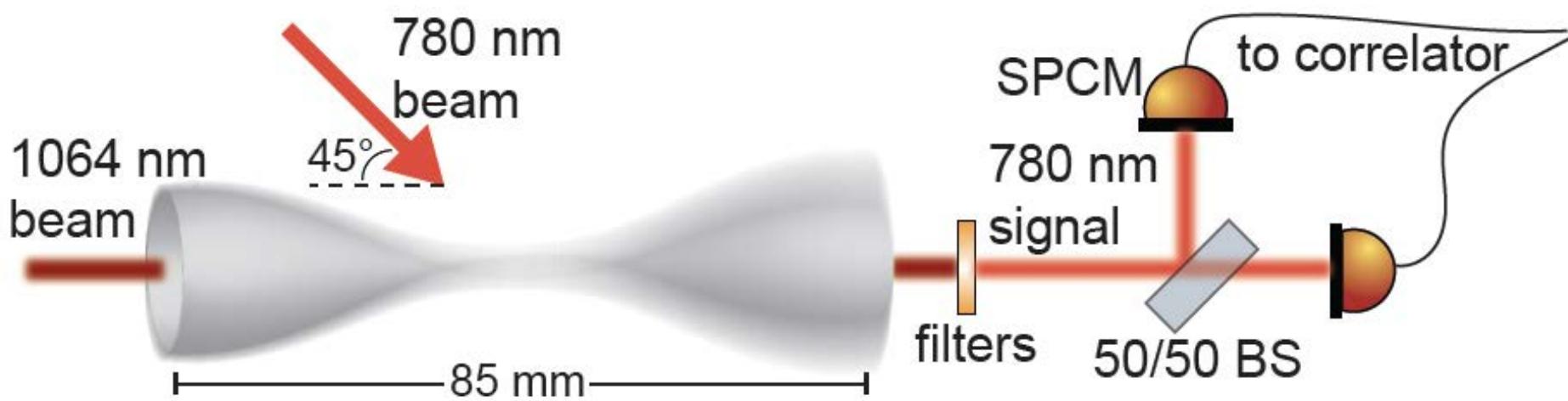
Mechanical modes

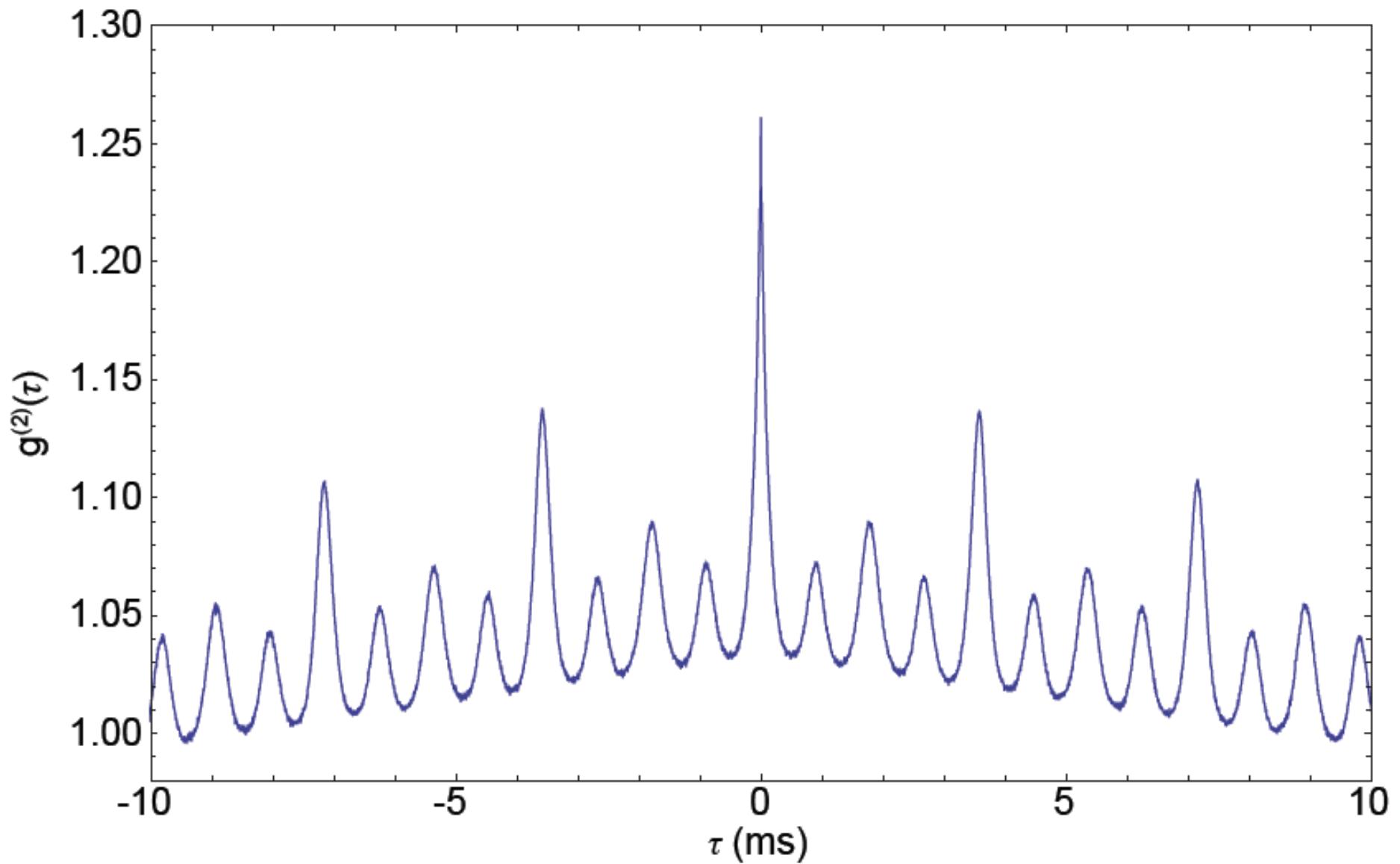
- Vibration (Violin)
- Torsion
- Compression (not discussed here)

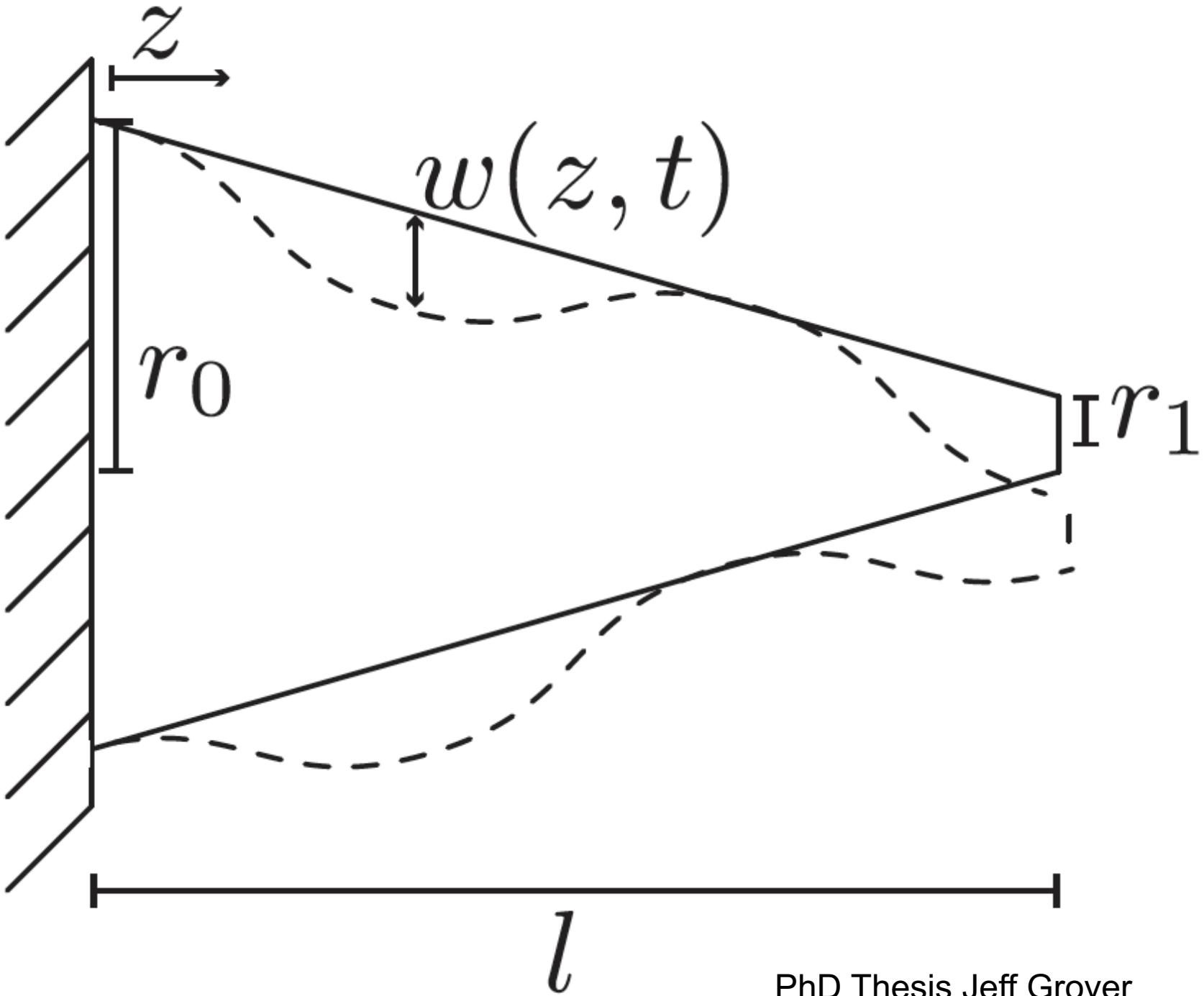


Unmodified Fiber: 125 μm diameter
Tensioning: NONE

Violin modes







The normal modes of a catiliber:

$$w(z, t) = h(z) \sin(\omega t)$$

$$\partial_u^4 h - 8 \frac{1-\alpha}{1-(1-\alpha)u} \partial_u^3 h + 12 \left(\frac{1-\alpha}{1-(1-\alpha)u} \right)^2 \partial_u^2 h = \frac{(lk)^4}{(1-(1-\alpha)u)^2} h,$$

$$\alpha = r_1/r_0, \quad k = 4\rho\omega^2/E r_0^2 \quad (\rho = 2.203 \text{ g}\cdot\text{cm}^3, E = 71.7 \text{ GPa})$$

$$r_0 = 62.5 \mu\text{m} \quad r_1 = 250 \text{ nm} \quad l = 39 \text{ mm}$$

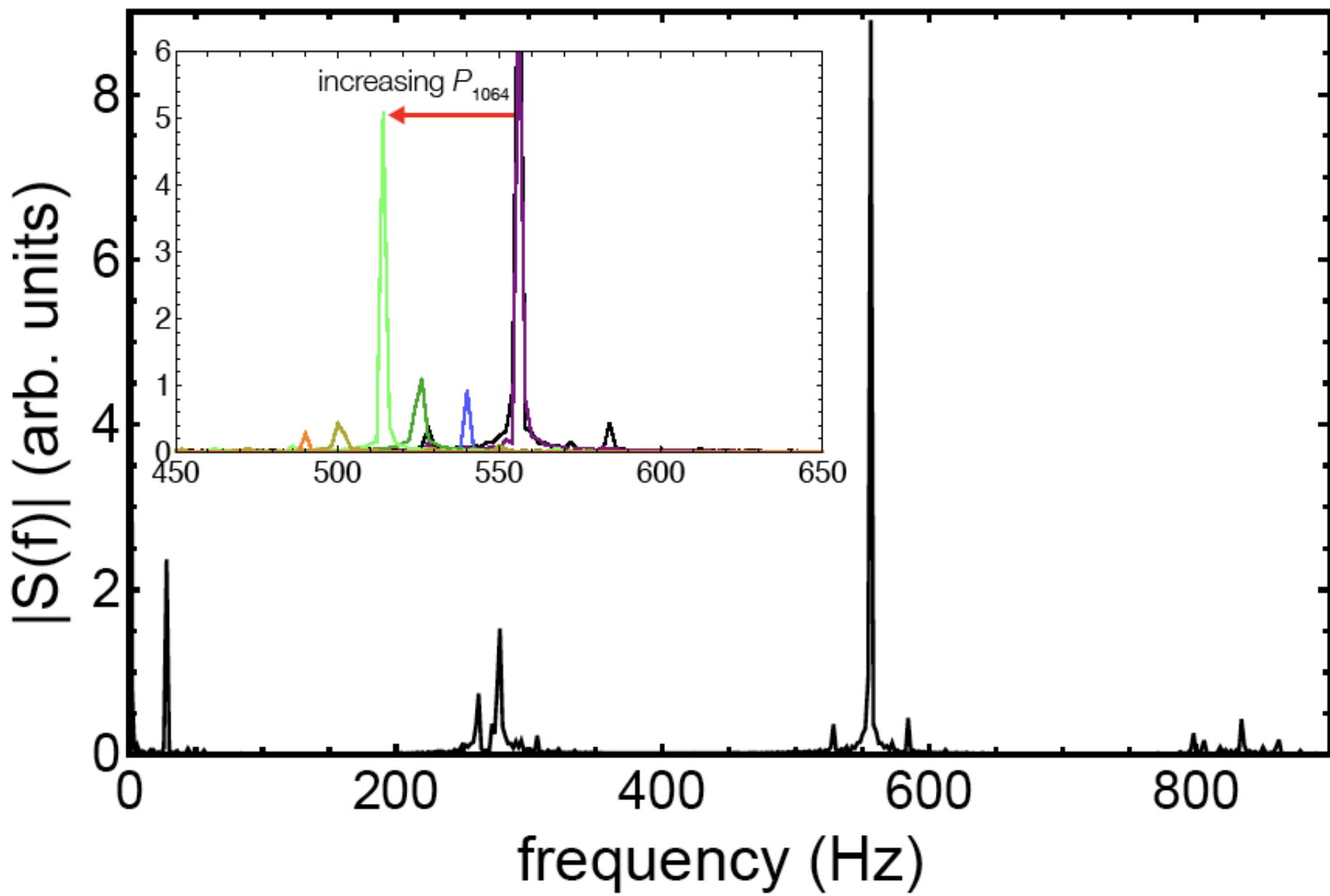
$$\omega/2\pi = 161.5, 392.3$$

With more tension

$$\omega'_n = \omega_n \sqrt{1 + U_n}$$

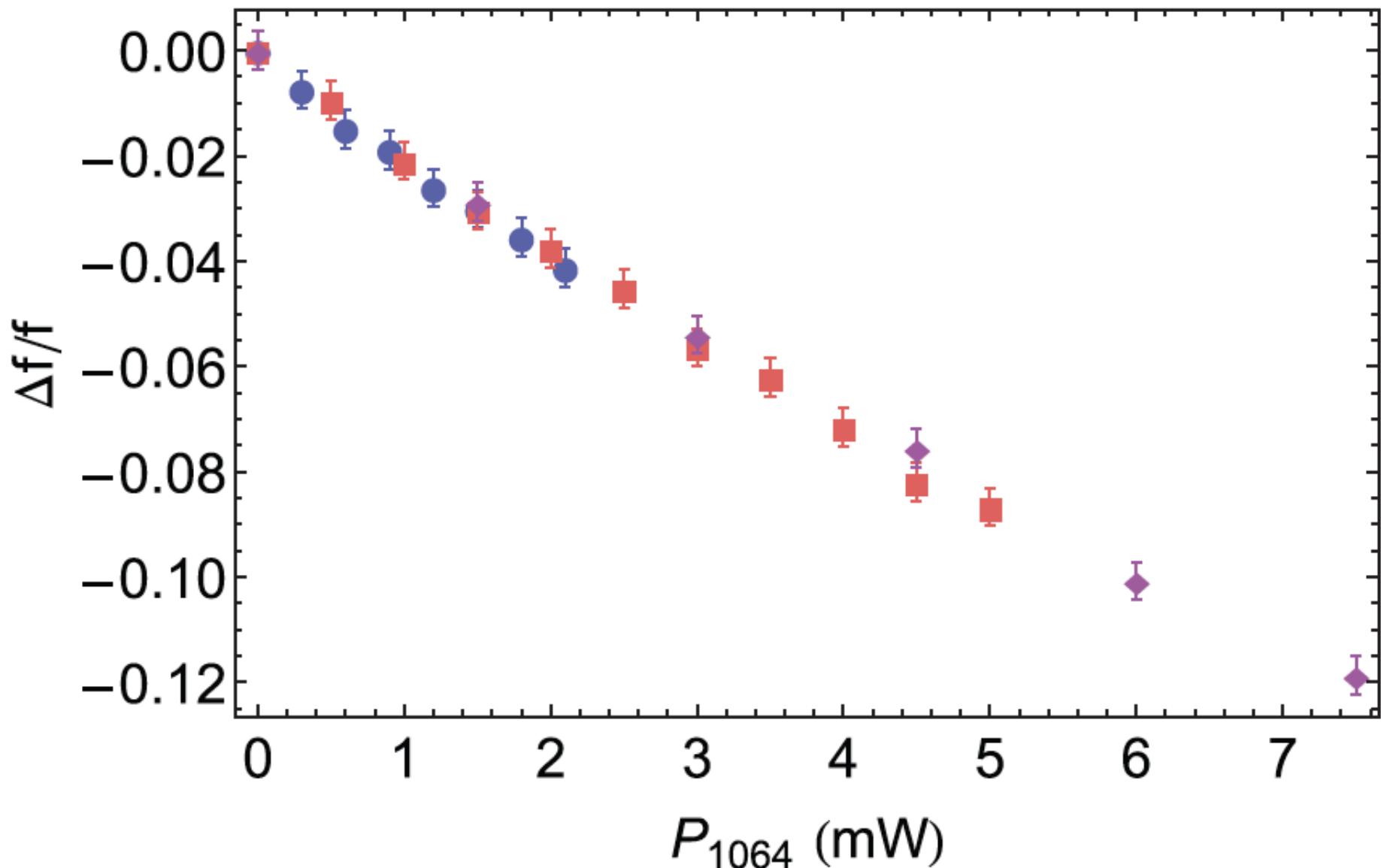
$$U_n = \frac{4}{(2n-1)^2\pi^2} \frac{F_{\text{axial}}l^2}{EI}$$

To obtain the first calculated frequency, the elongation of the fiber should 69 μm



The frequency decreases with T

PhD Thesis Jeff Grover



Torsional Modes

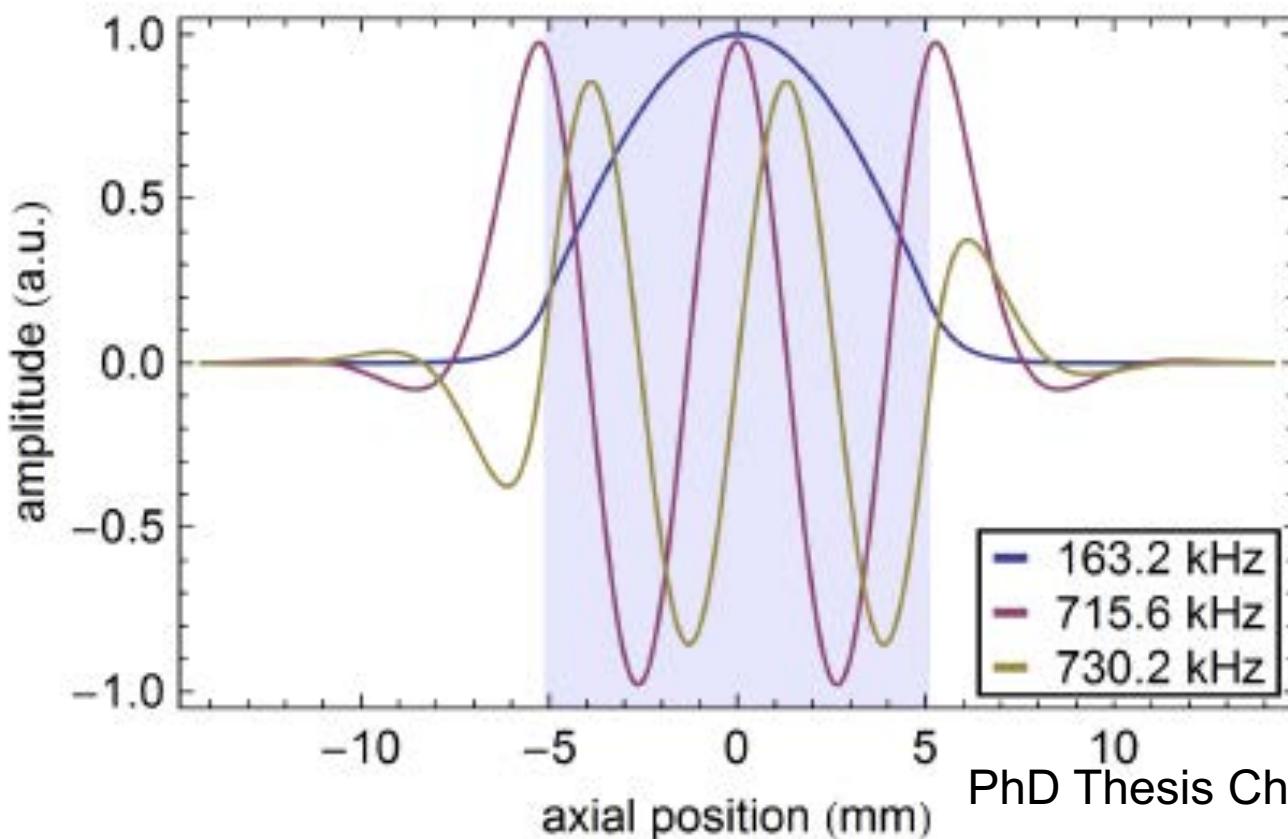
Torsion of a plane of a cylinder by an angle ϕ

$$c_t^{-2} \partial_t^2 \phi(t, z) - \partial_z^2 \phi(t, z) = 0$$

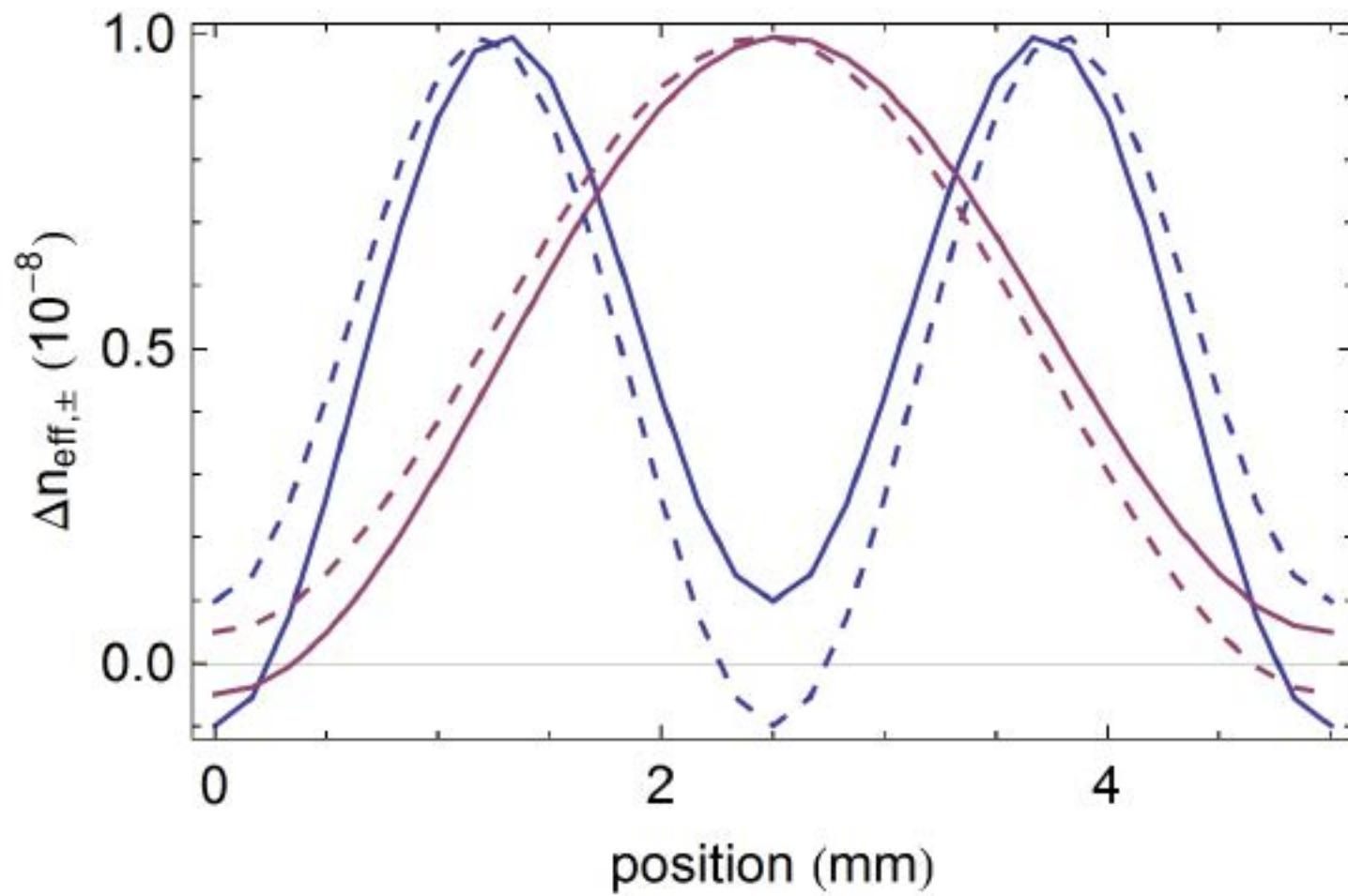
$$\phi(t, z) = \phi(z) \cos(\omega t)$$

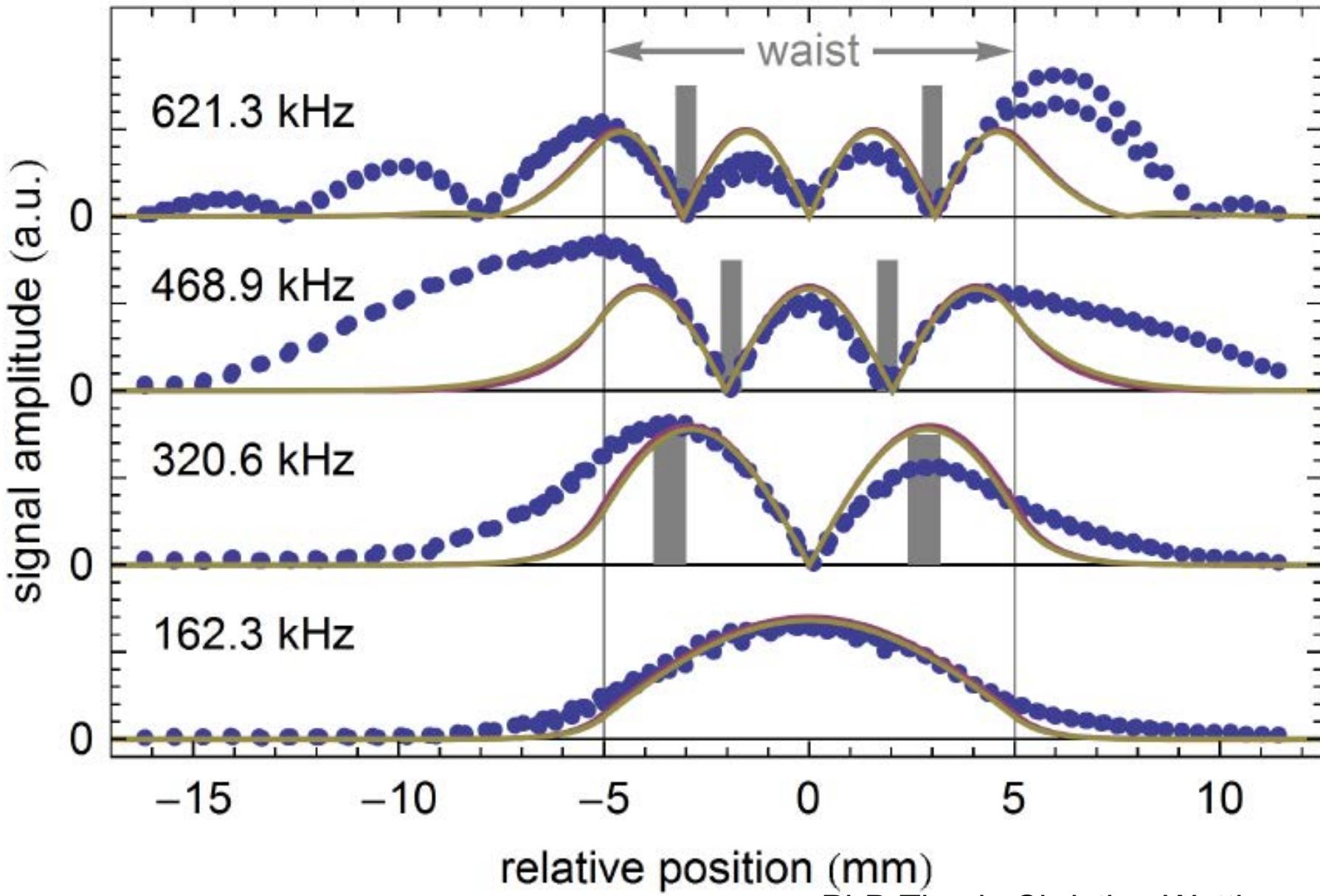
$$\partial_z^2 \phi(t, z) + k_0^2 \phi(t, z) = 0$$

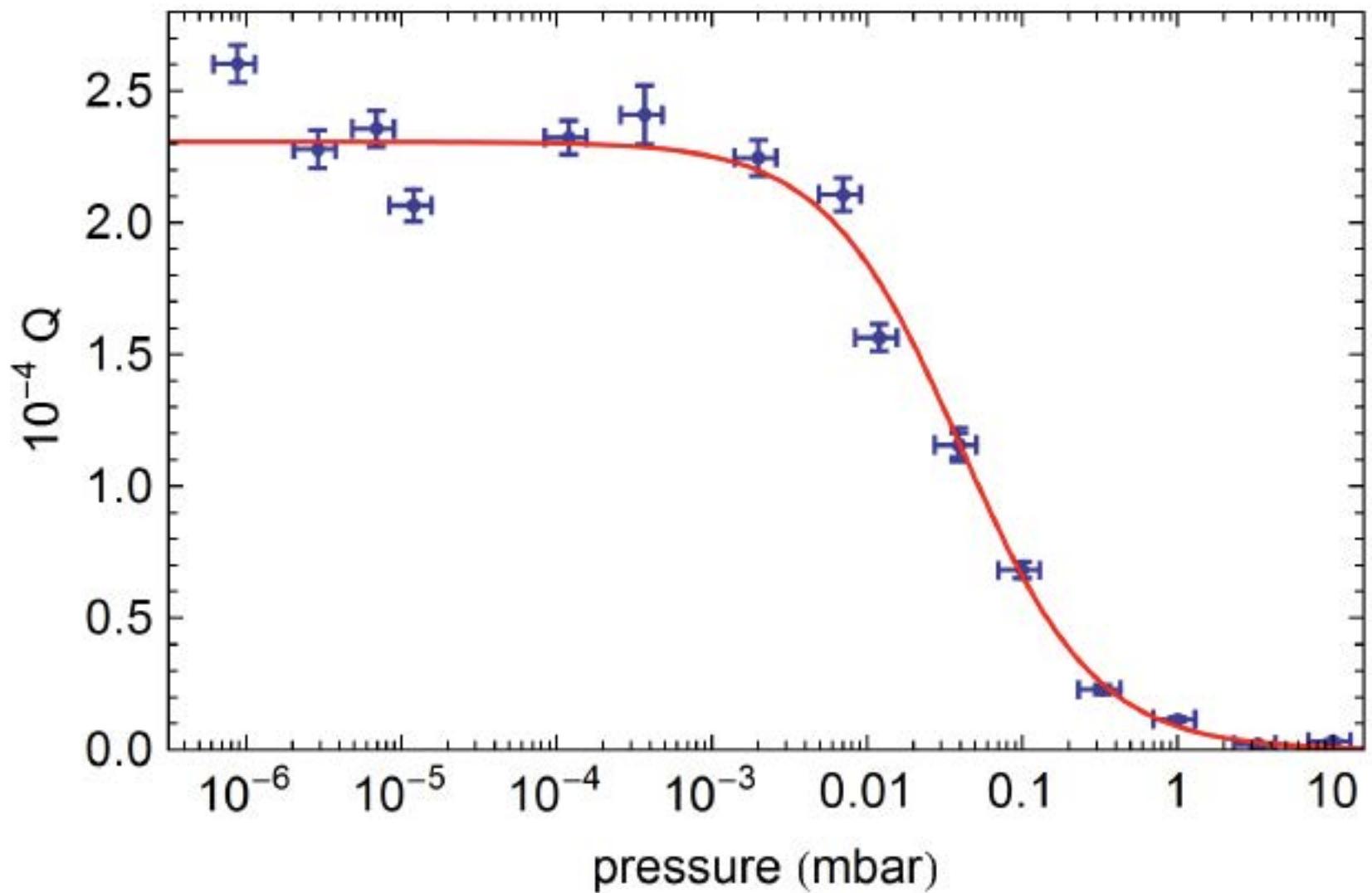
Amplitude of the modes of the nanofiber considering the tapers



Change in the index of refraction



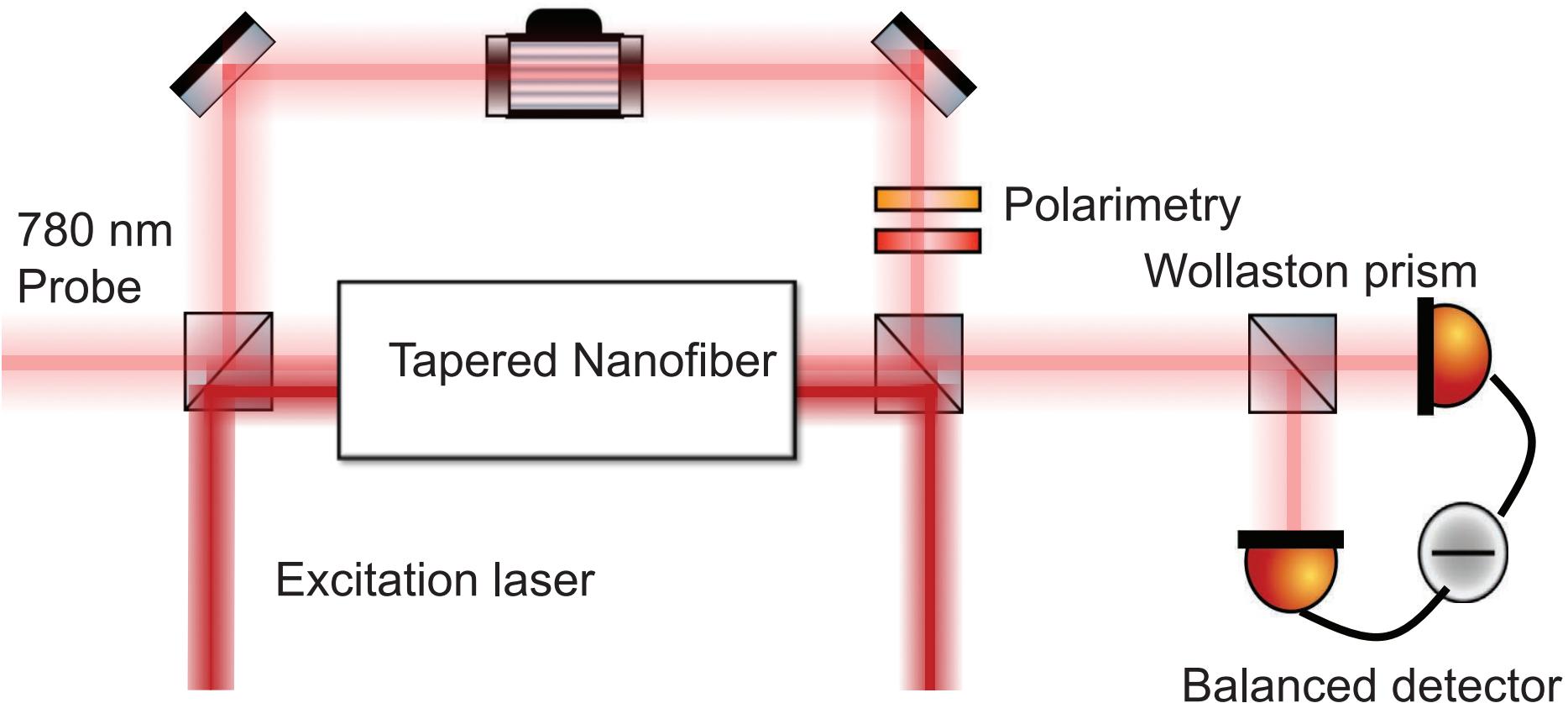




Transfer the intrinsic angular momentum (circular polarization) from the light to the torsional modes of the ONF.

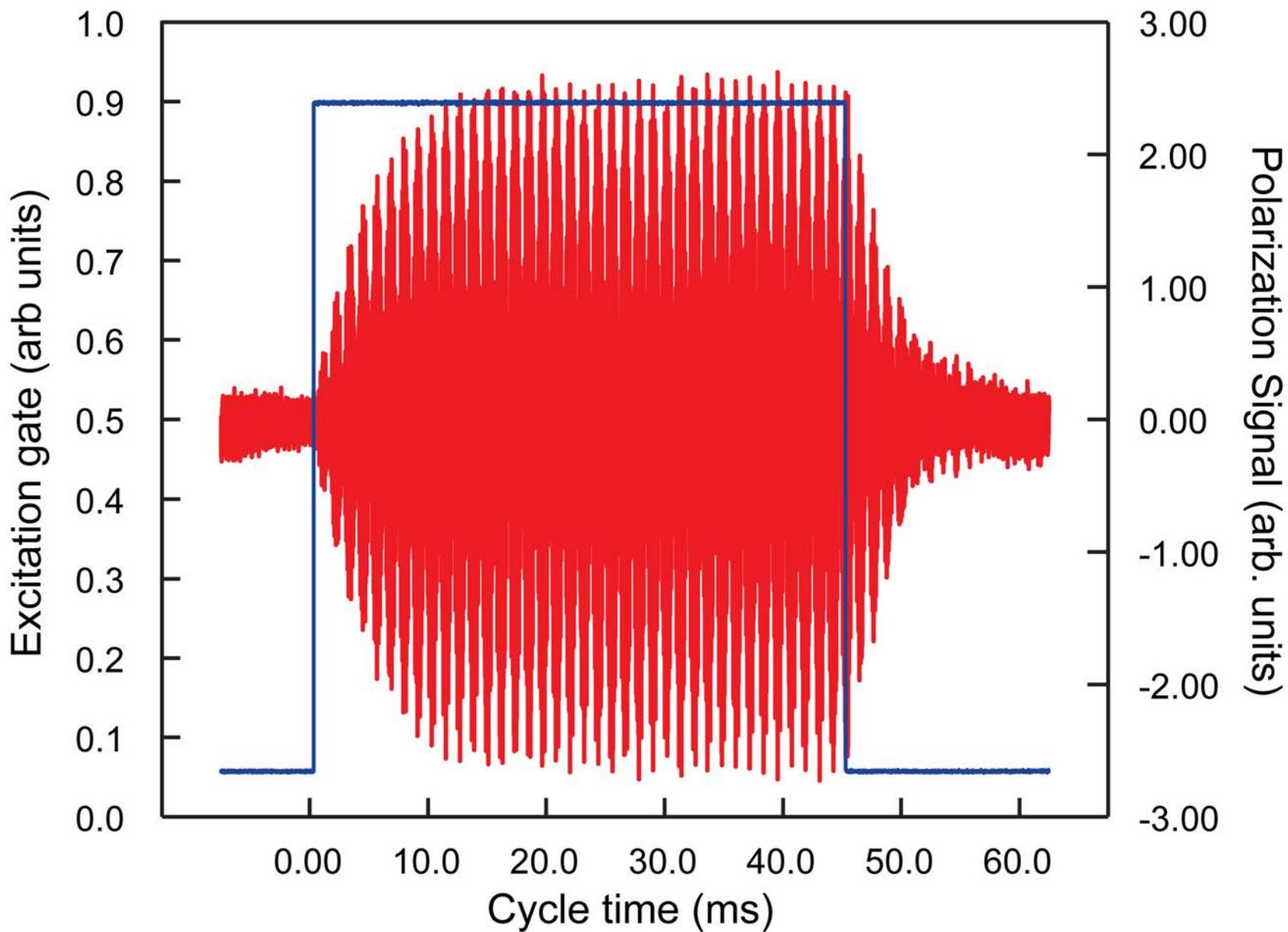
- The stress produced by the torsion on the nanofiber affects its index of refraction, creating a birefringence in the medium.

Frequency shift ~ 1 MHz

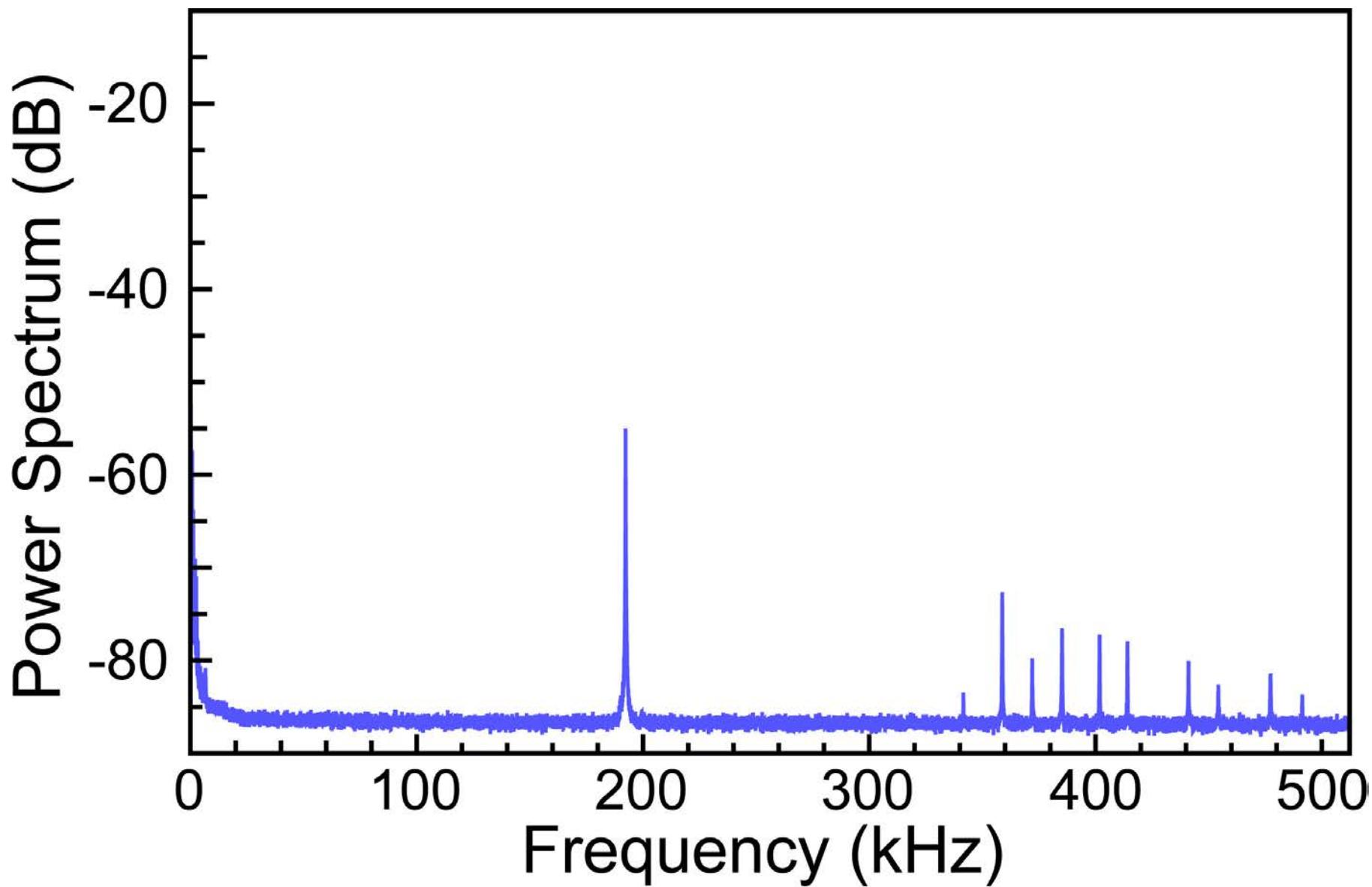


Excitation 1060 nm: circular polarized light (modulated at the mode frequency). Excitation with linear or circular polarization. Probe linear and weak.

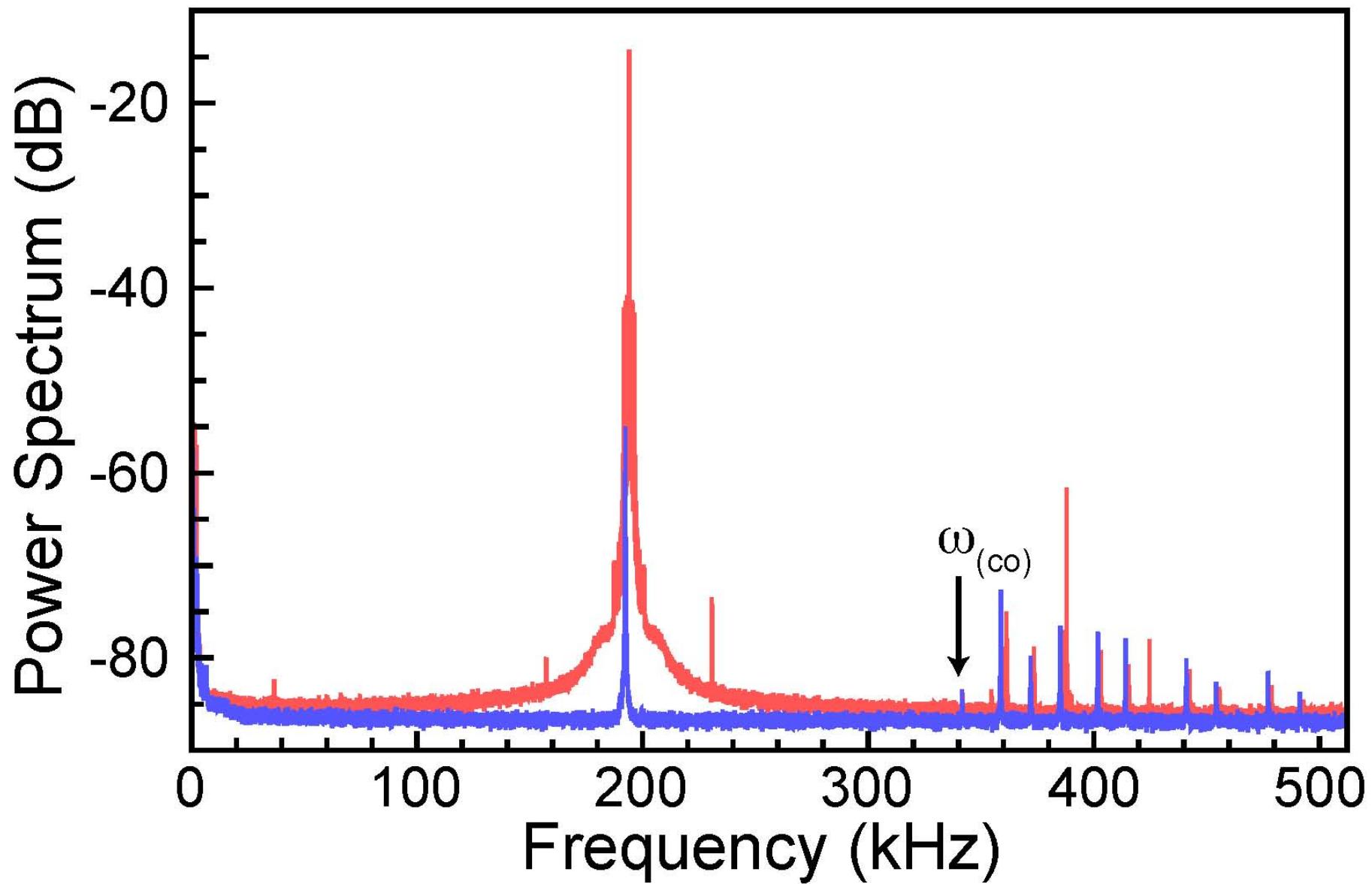
Excitation and response.



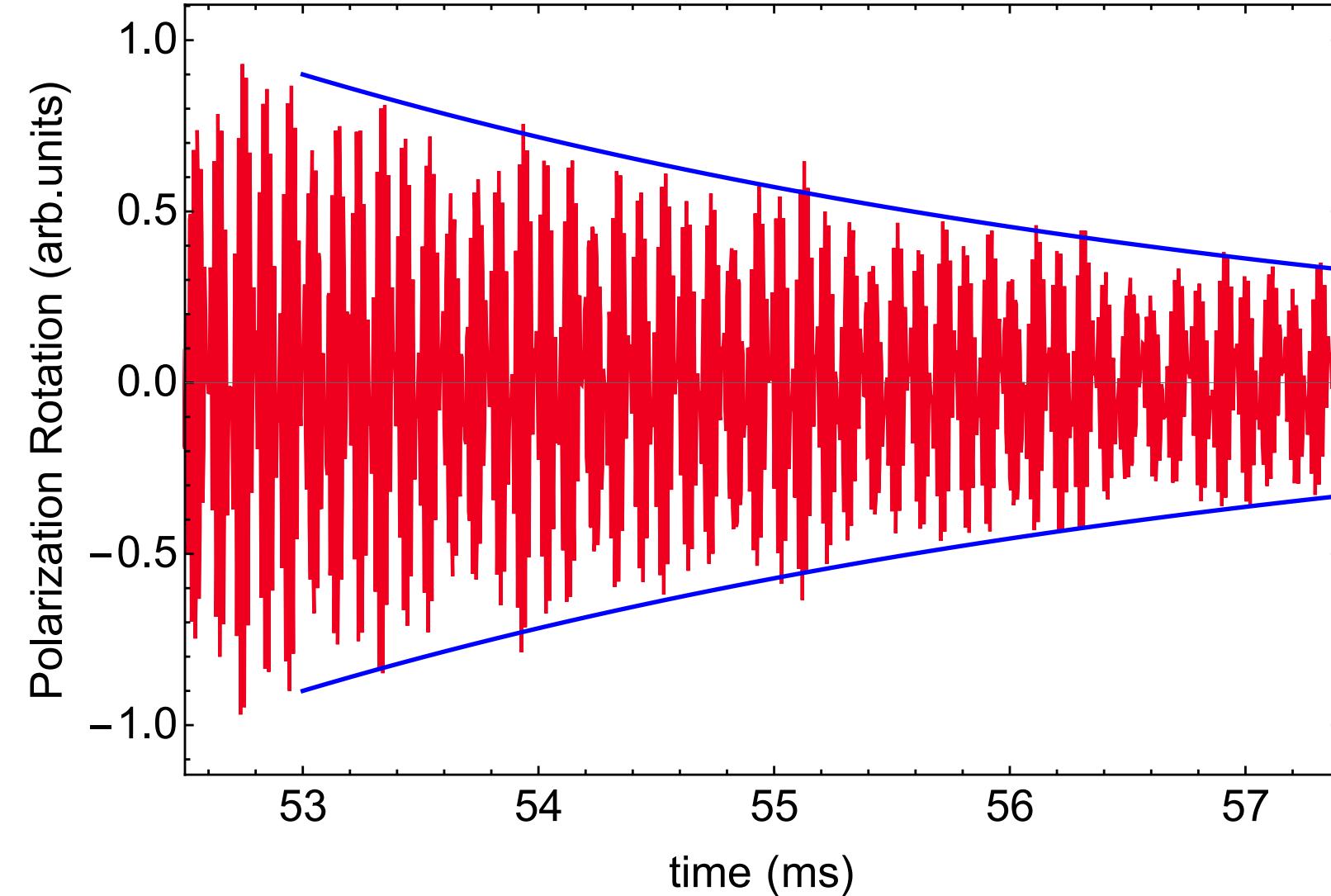
Thermal excitation spectrum

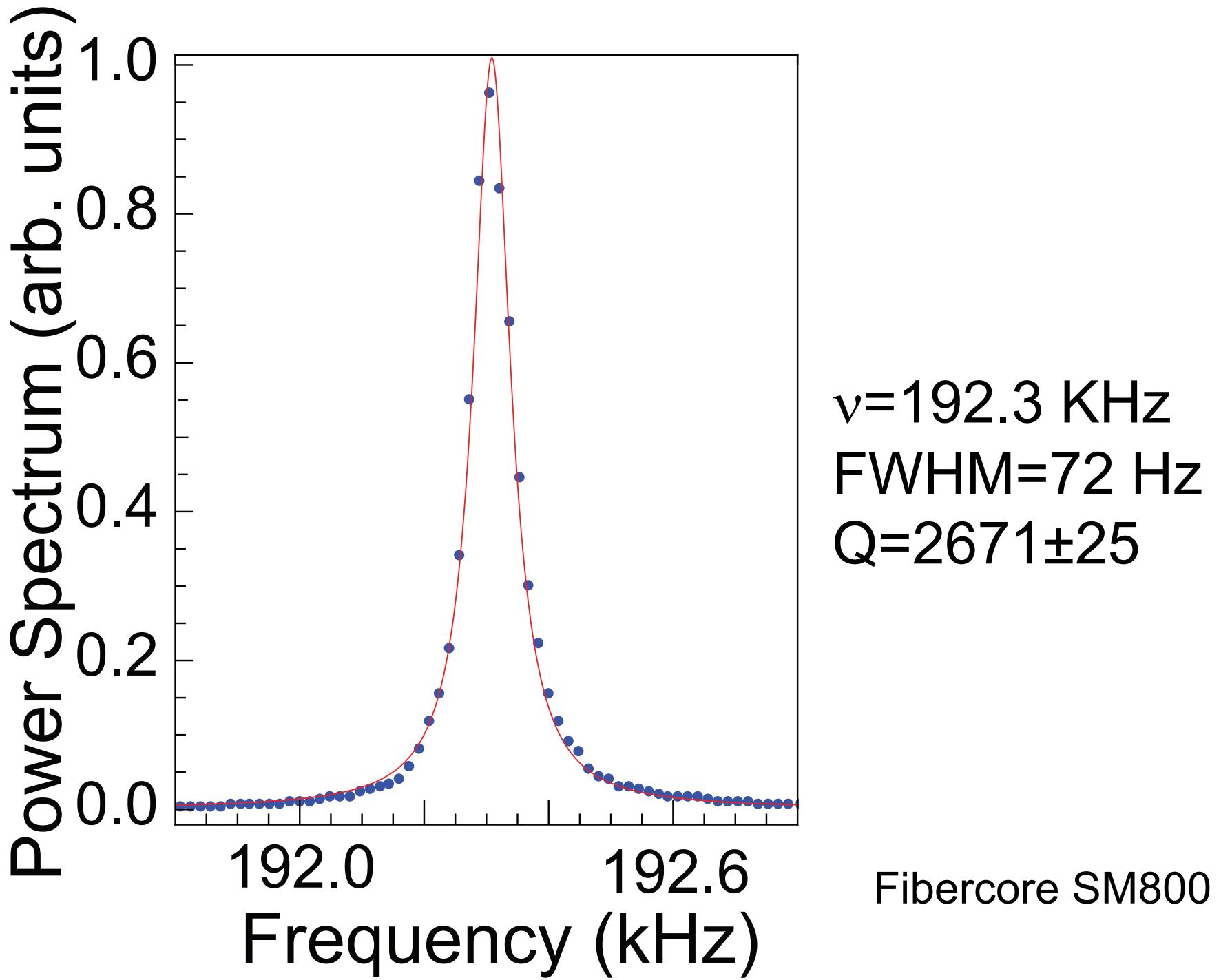


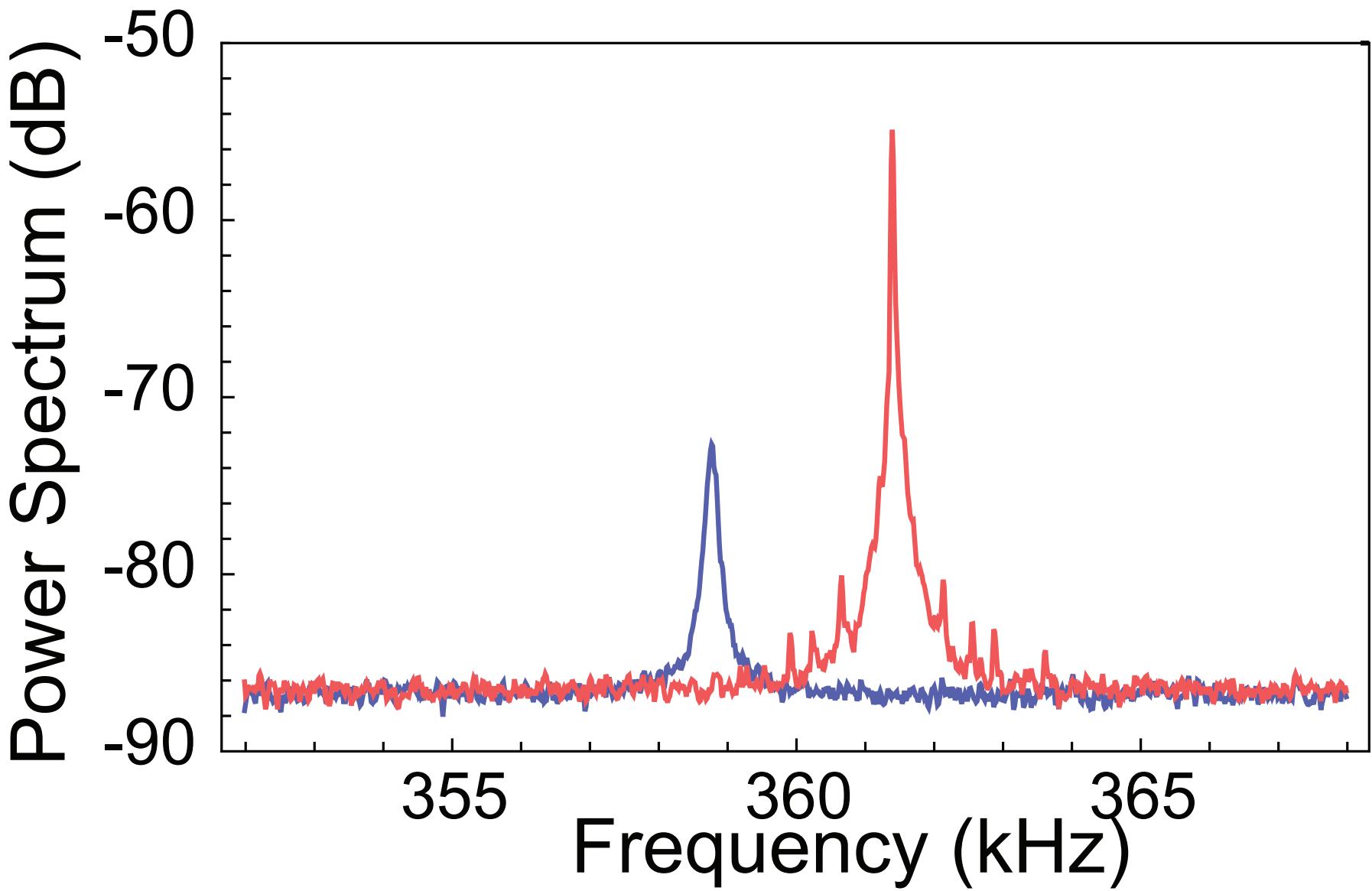
Thermal excitation (blue), Resonant excitation with circular polarization of the first mode (red).



Torsional mode decay with modulation from violin modes

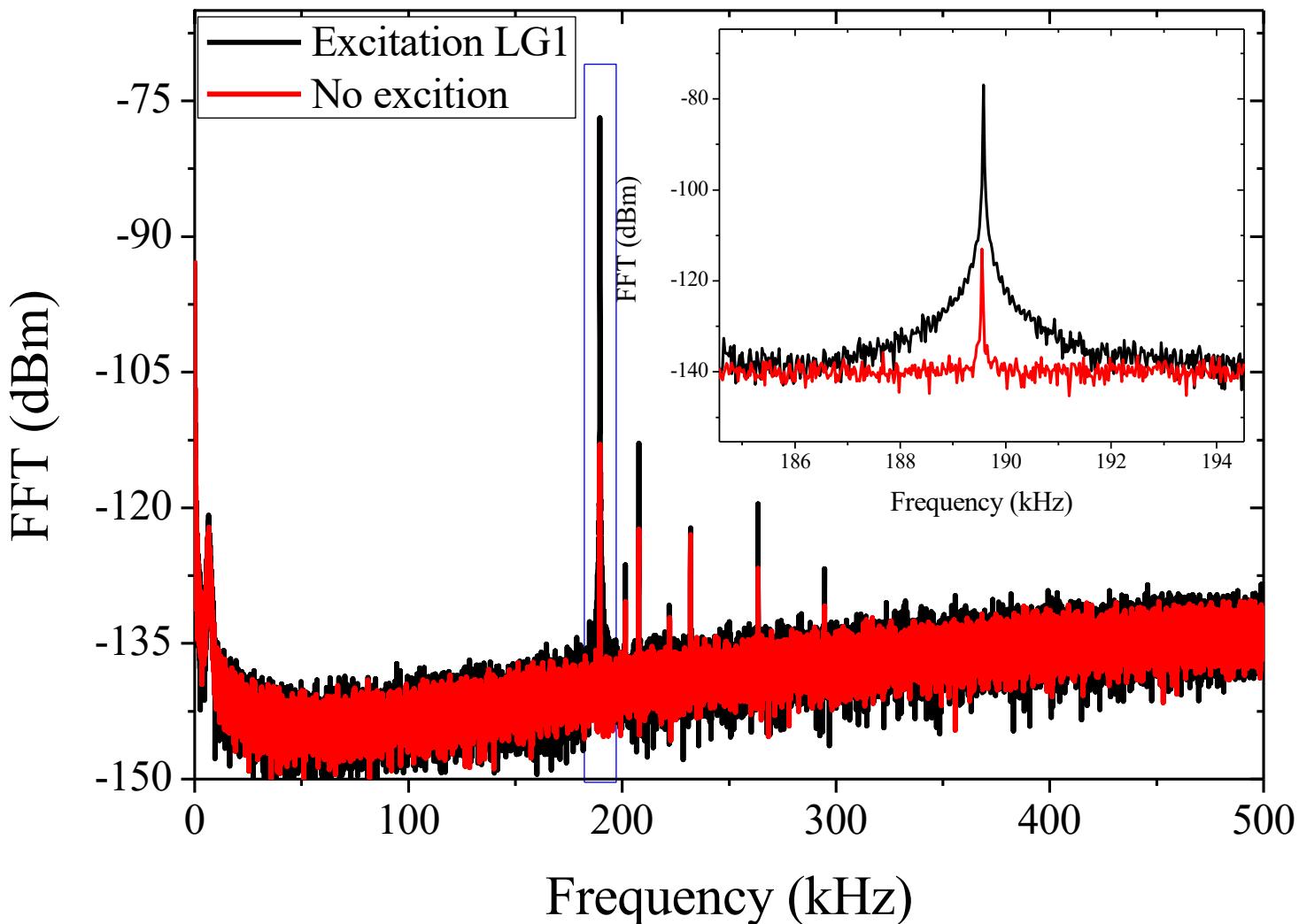






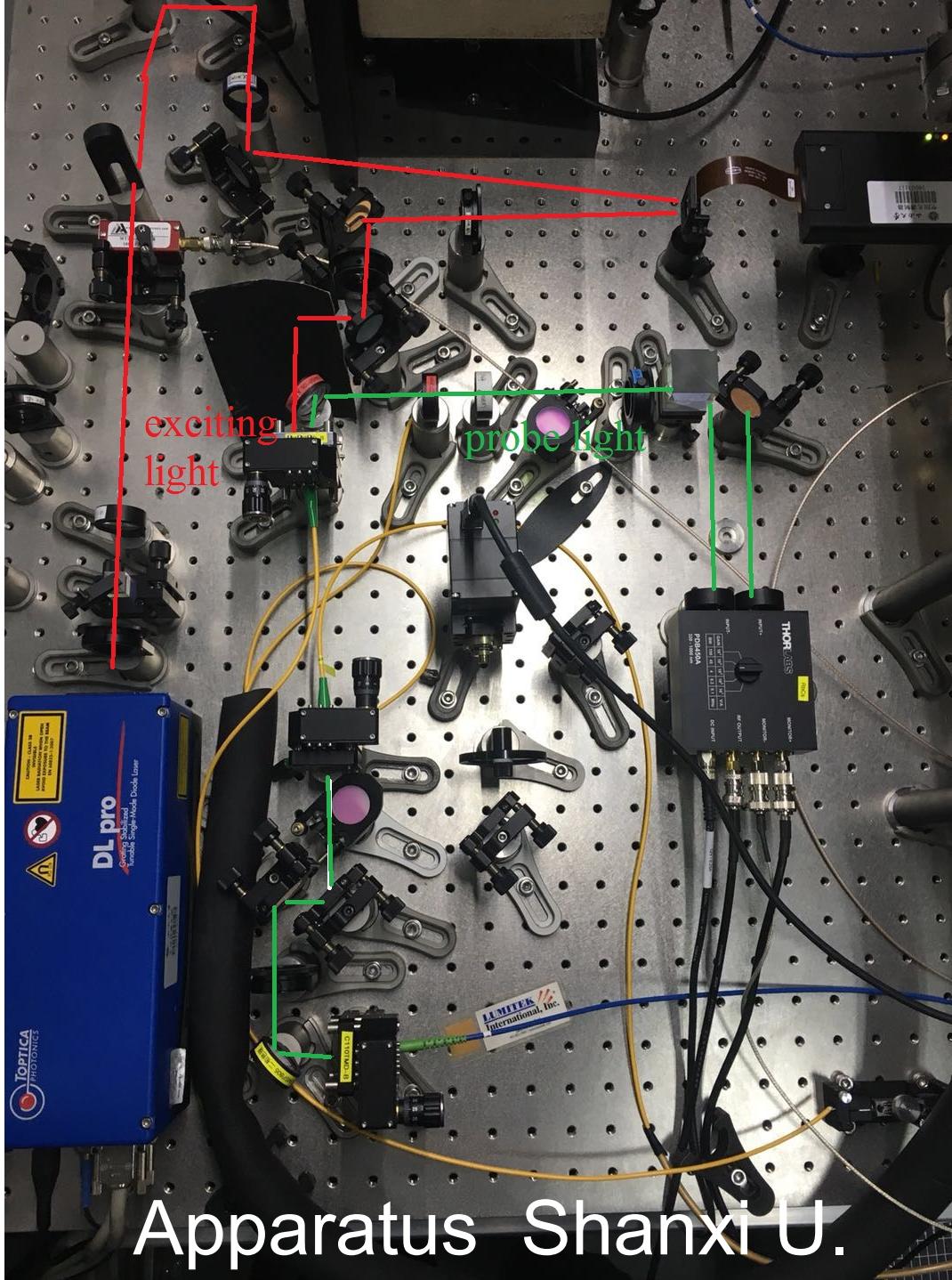
Change of temperature causes changes to Young and Poisson moduli causing a change in the wave velocity.

Excitation with linearly polarized light with light
with orbital angular momentum e. g. LG01

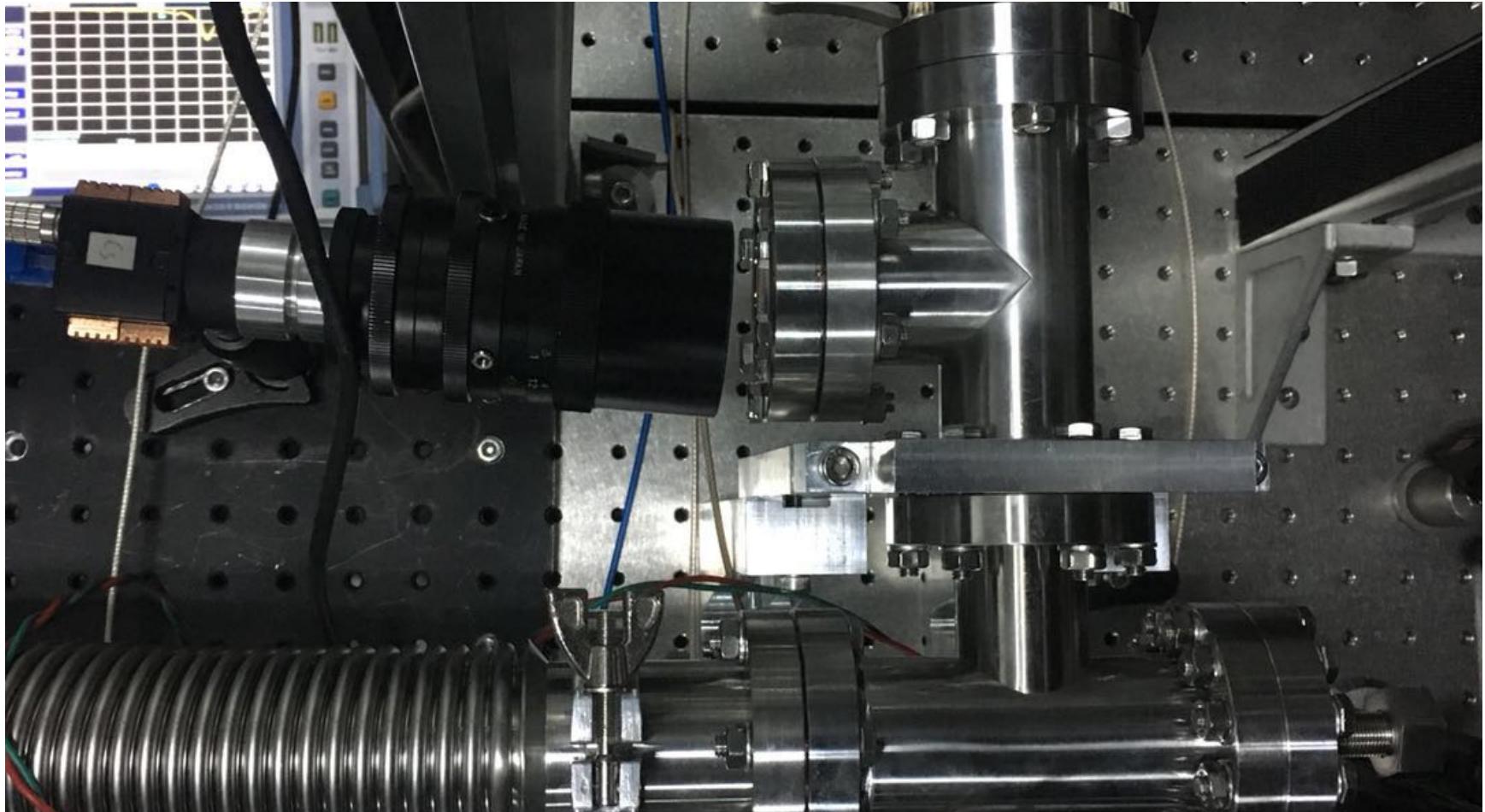


The Q factor of the first mode is greater than 2×10^4 ; Increases by some 40 dB; Acopla

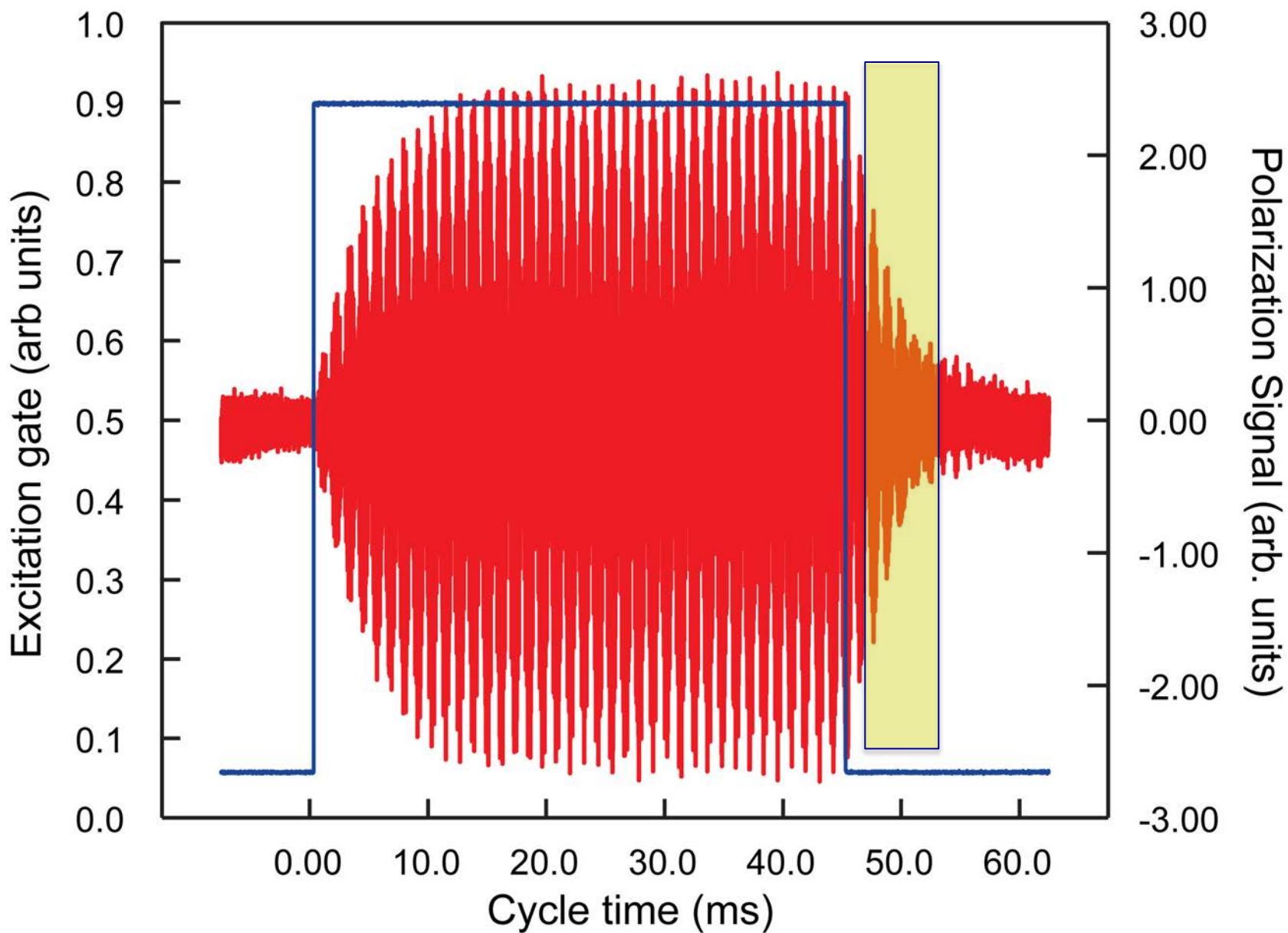
Preliminary results



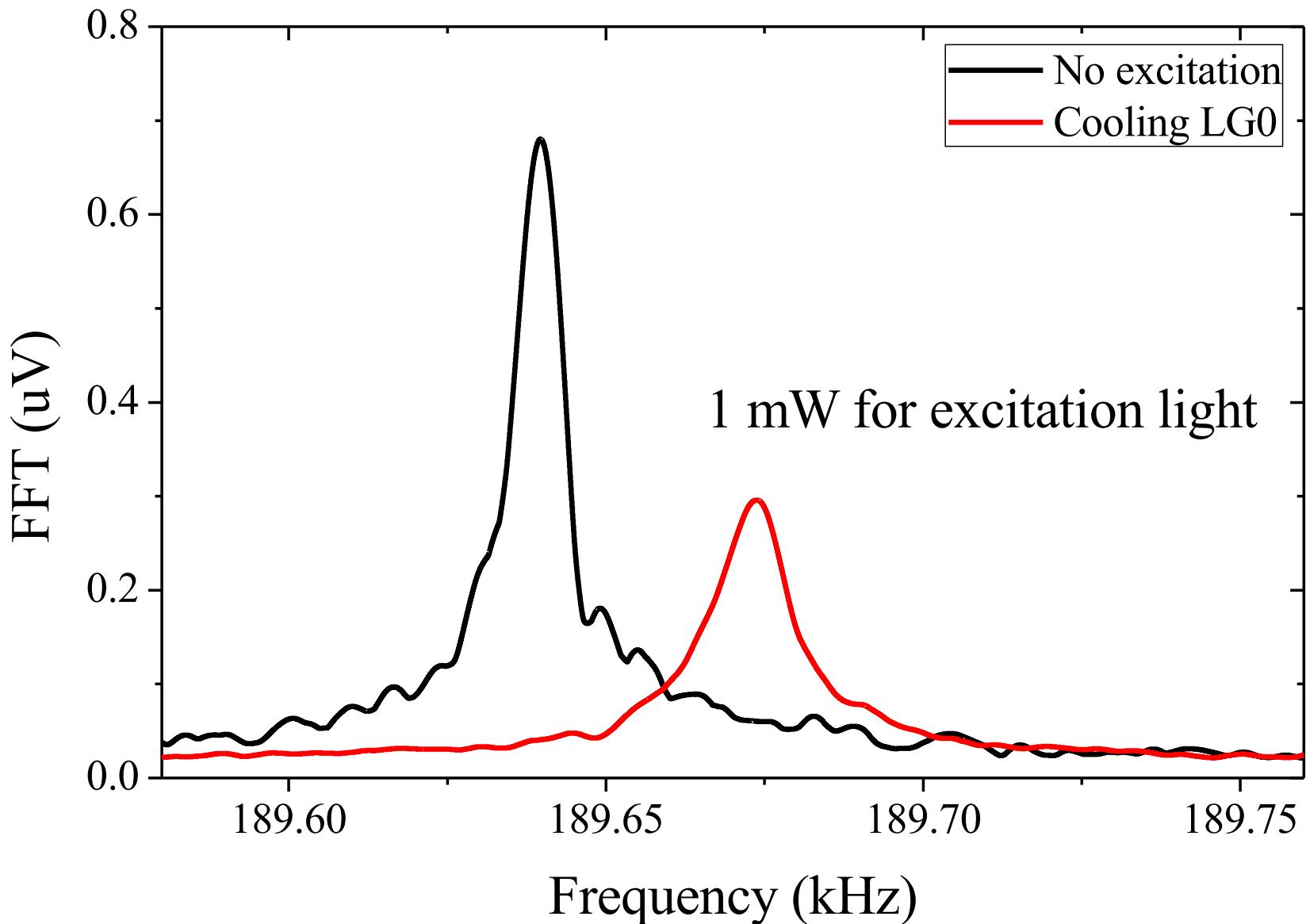
Measurement of Rayleigh Scattering

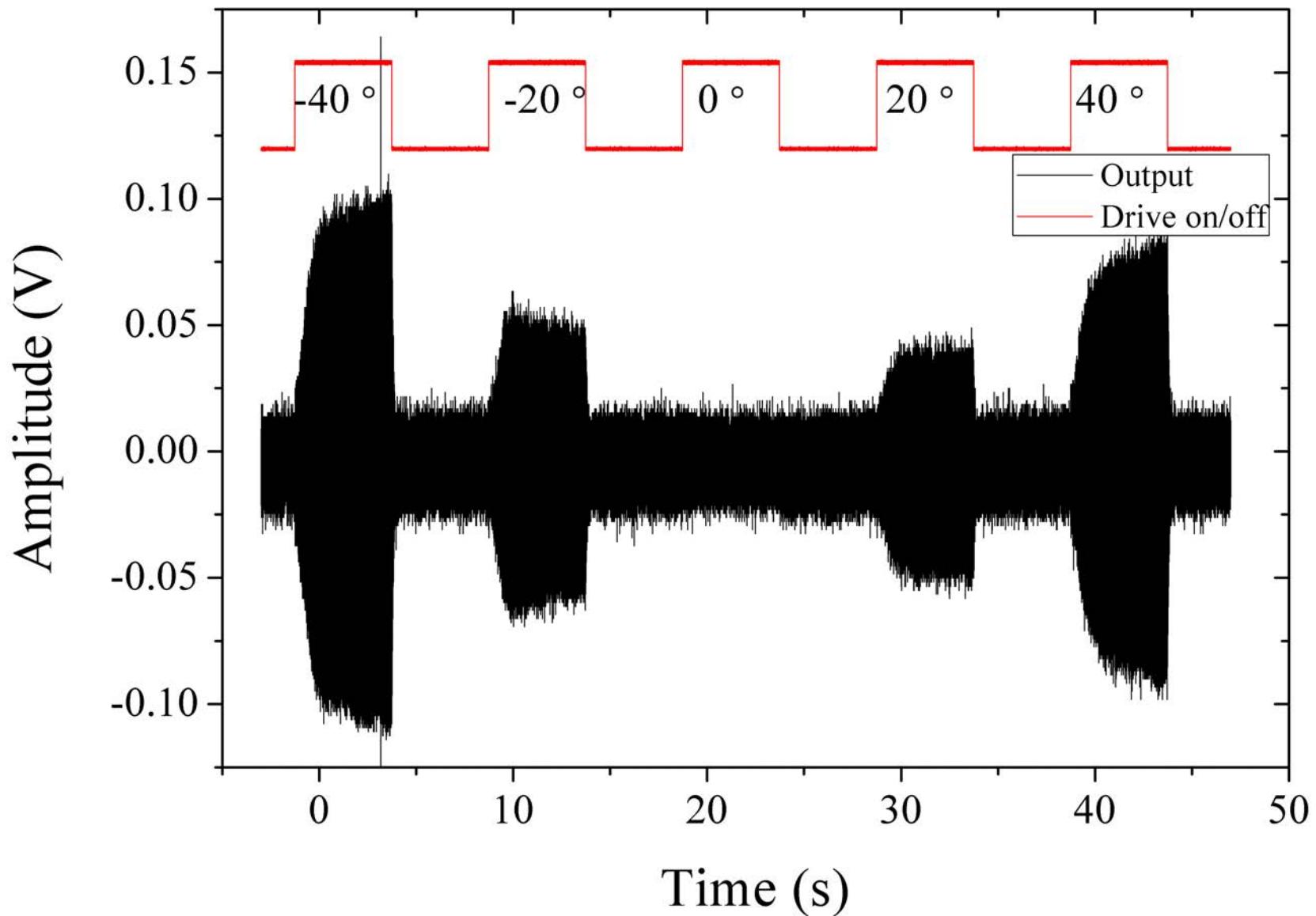


Excitation and response.



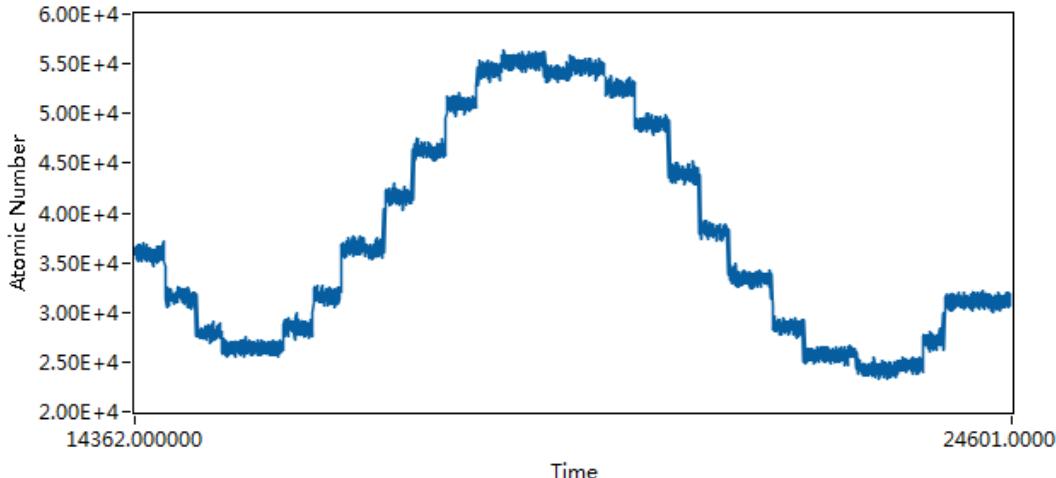
The surprise



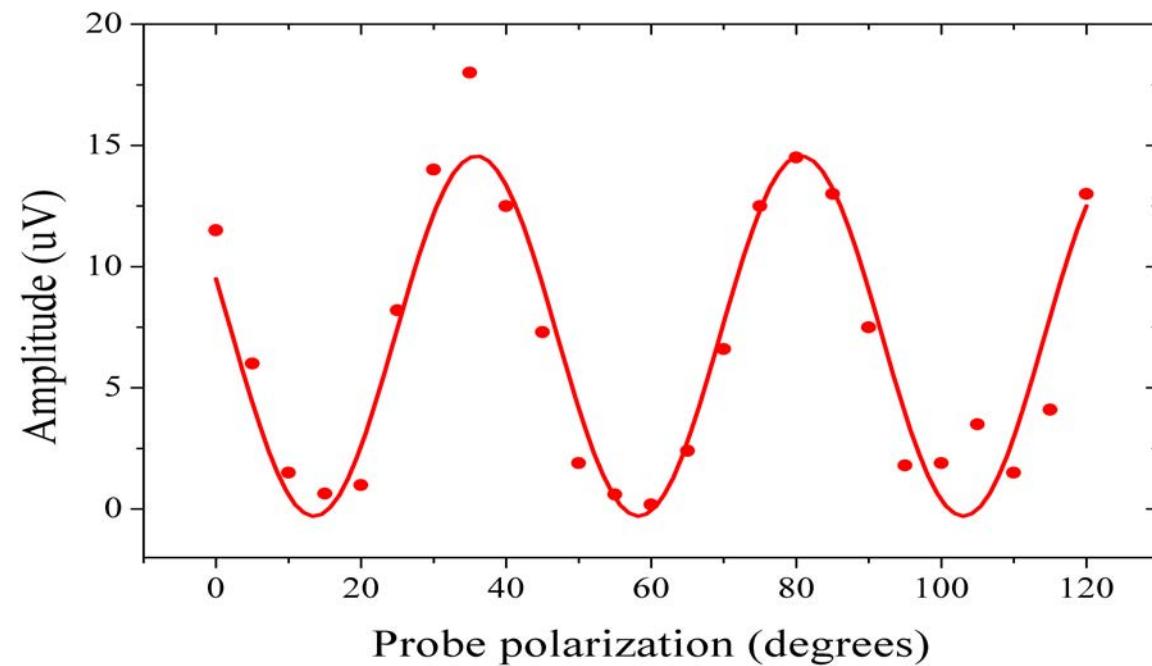


Amplitude (resonant excitation) of the probe for five different polarizations of the drive.

ROI intensity



Dispersión de Rayleigh



Study of the
amplitude of the
thermal noise as a
function of the
probe polarization

The electromagnetic field exerts a torque if the medium has a polarizability

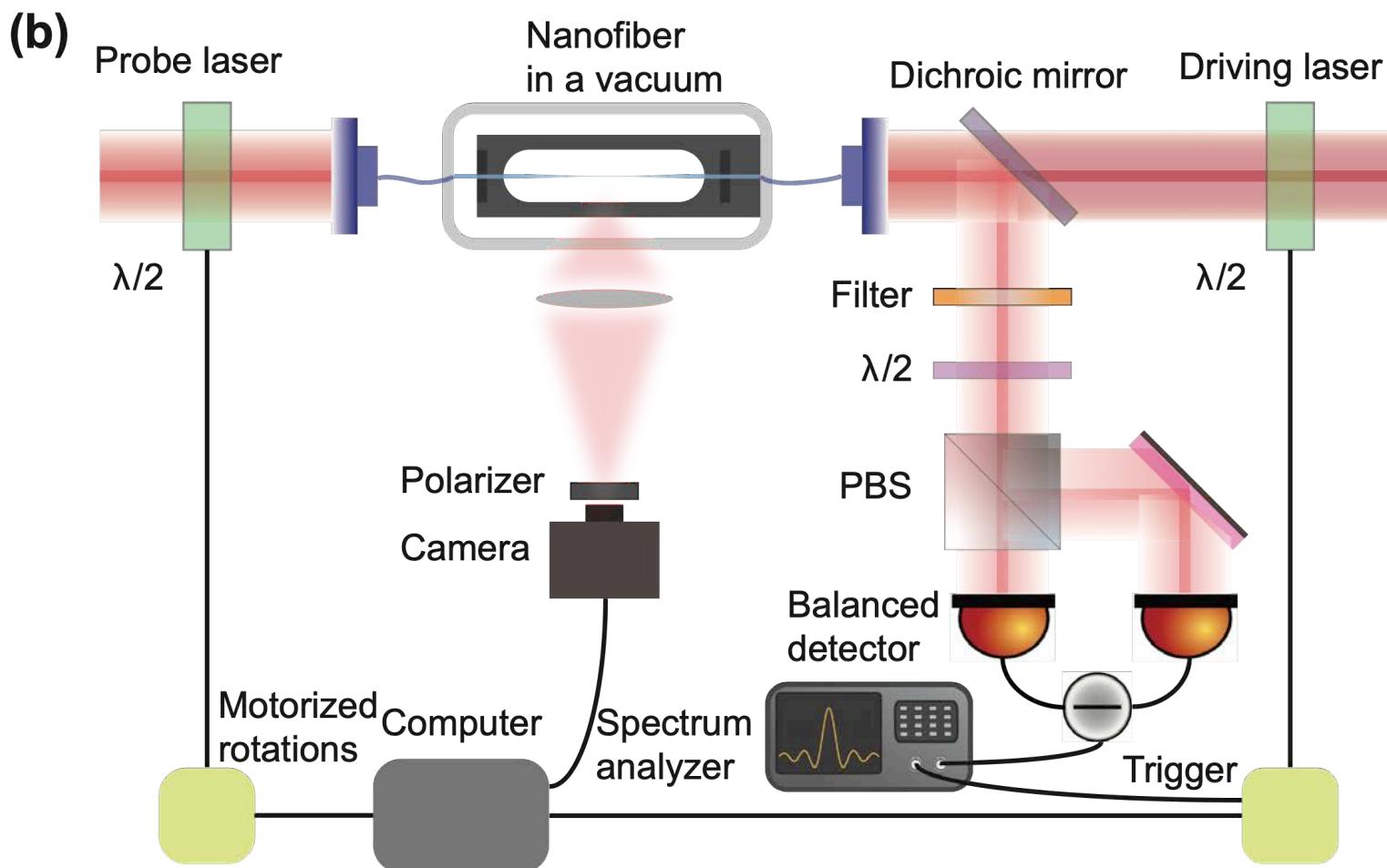
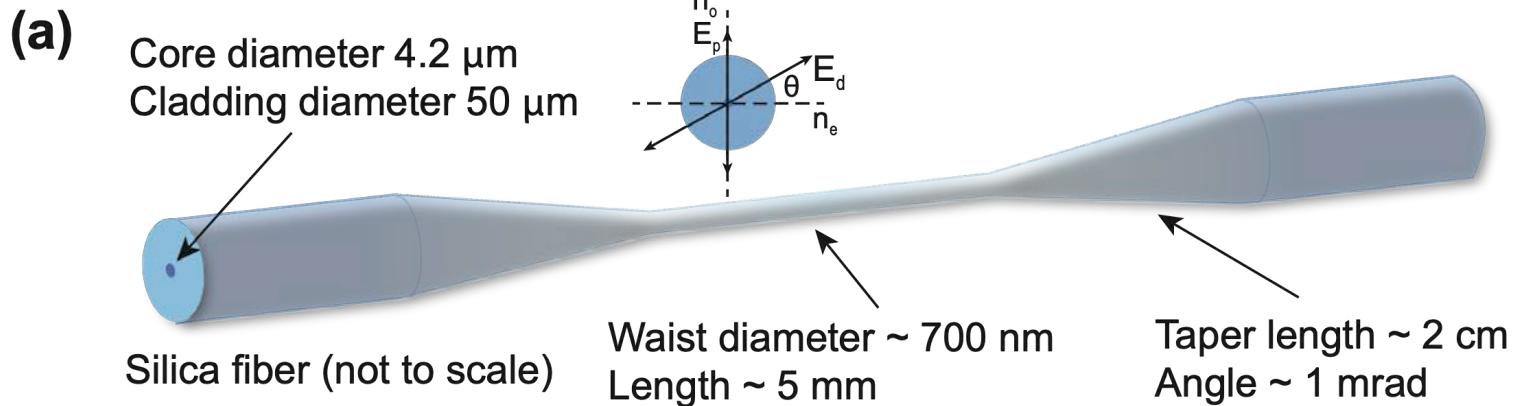
$$\vec{T} = \vec{P} \times \vec{E}$$

$$T(t) = \frac{\epsilon E^2}{2\omega_l} \sin \Gamma \sin 2\theta(t)$$

$$\Gamma = kd(n_o - n_e)$$

The intensity (A^2) is very large in the nanofiber, $n_o - n_e \sim 10^{-8}$

Richard A. Beth, "Mechanical Detection and Measurement of the Angular Momentum of Light" Phys. Rev. 50, 115 (1936) and PhD Thesis Christian Wuttke.



The Torque

$$T(t) = \frac{\epsilon E^2}{2\omega_l} \sin \Gamma \sin 2\theta(t) \quad \Gamma = kd(n_o - n_e)$$

Torsion of a disk

$$I\ddot{\theta}(t) + \gamma\dot{\theta}(t) + \kappa\theta(t) - A_0 \sin(2\theta(t)) = T_{th}.$$

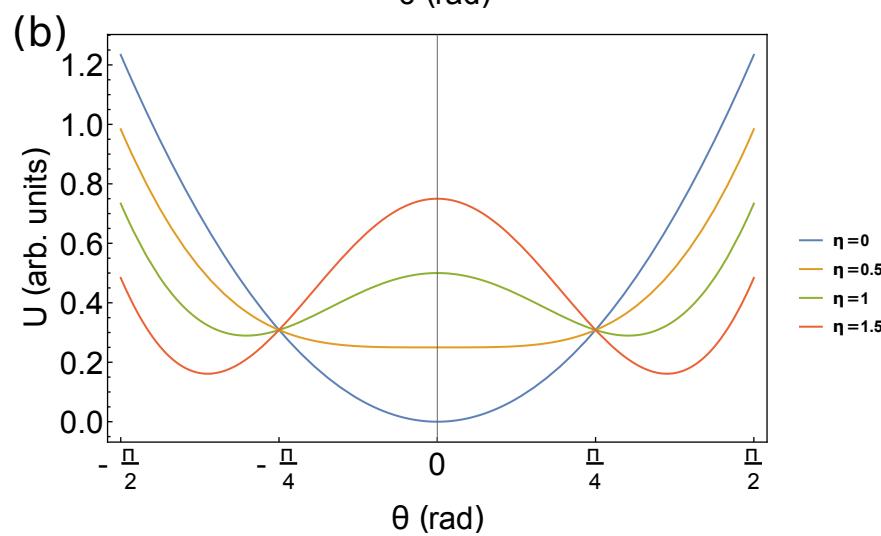
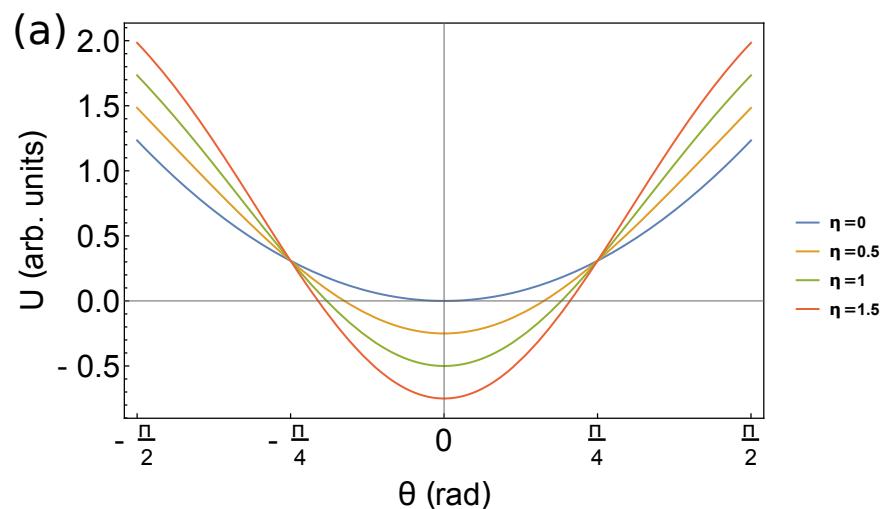
Spectral density of the rotation

$$S_{\delta\theta} = \frac{4k_B T \gamma}{((\kappa - A_0 \cos 2\theta(\omega)) - I\omega^2))^2 + \gamma^2 \omega^2}$$

$$A_0 = \frac{\epsilon E^2}{2\omega_l} \sin \Gamma$$

$$U = U_0 + \frac{1}{2}\kappa\theta^2 + \frac{1}{2}A_0 \cos 2\theta.$$

$$\kappa_{eff} = \kappa + 2A_0$$



Perturbation treatment around ss

$$I\ddot{\delta\theta}(t) + \gamma\dot{\delta\theta}(t) + \delta\theta(t) [\kappa - A_0 \cos(2\theta_{ss})] = T_{th}.$$

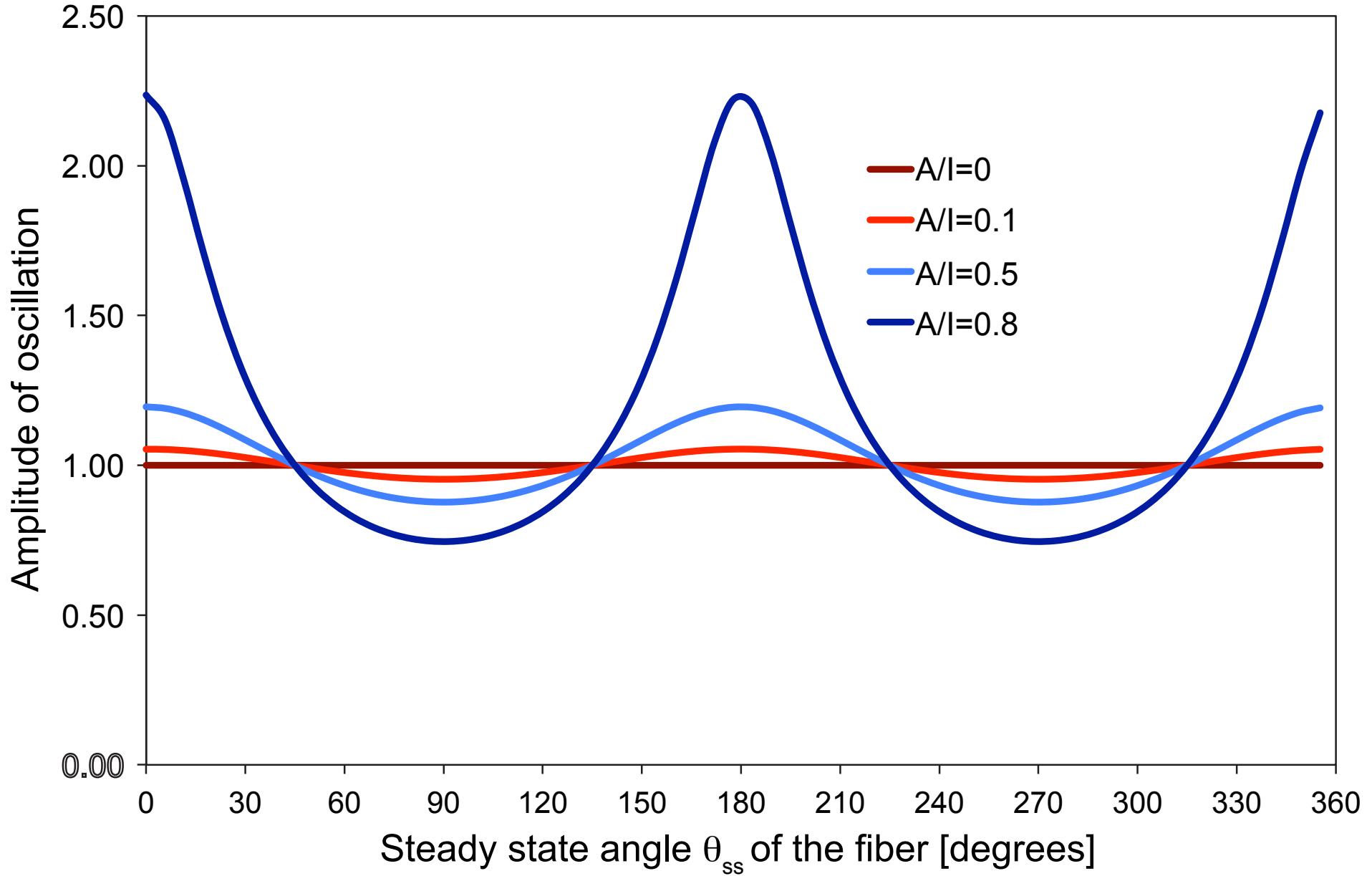
$$\omega_{ss} = \sqrt{\frac{\kappa - A_0 \cos 2\theta_{ss}}{I}}$$

Frequency shift

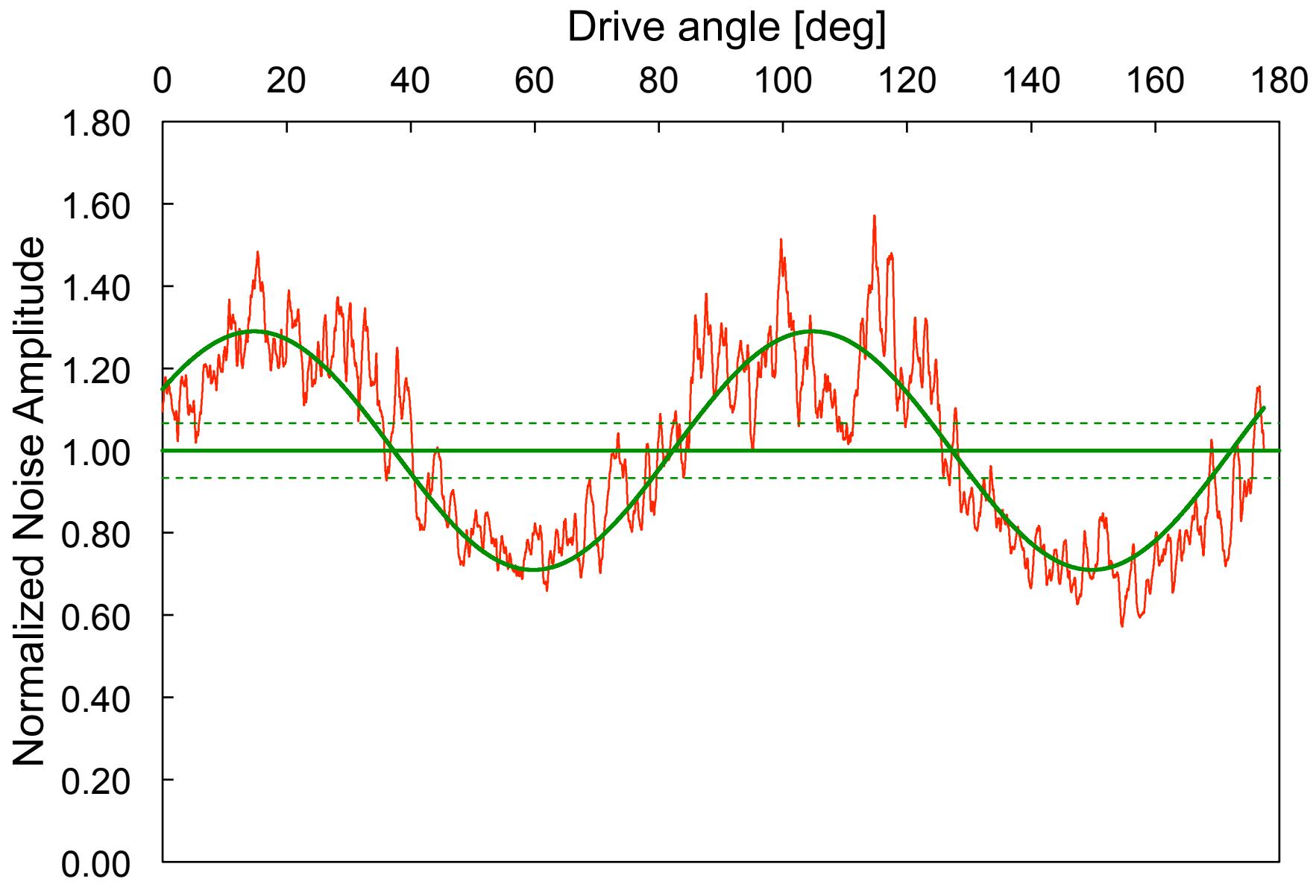
$$\delta\omega = -\sqrt{\frac{\kappa}{I}} \left(\frac{A \cos 2\theta}{2\kappa} \right)$$

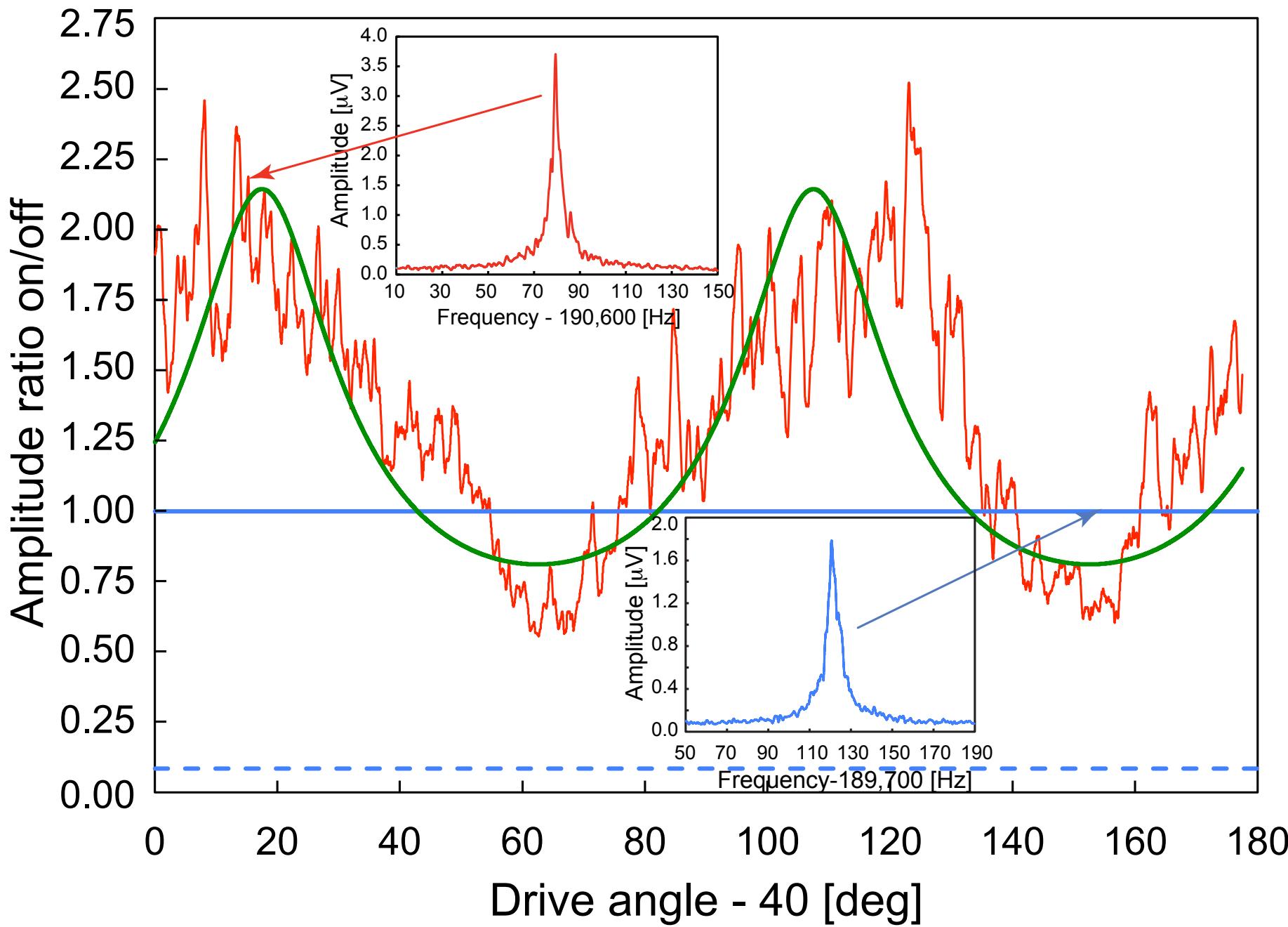
Amplitude of oscillation

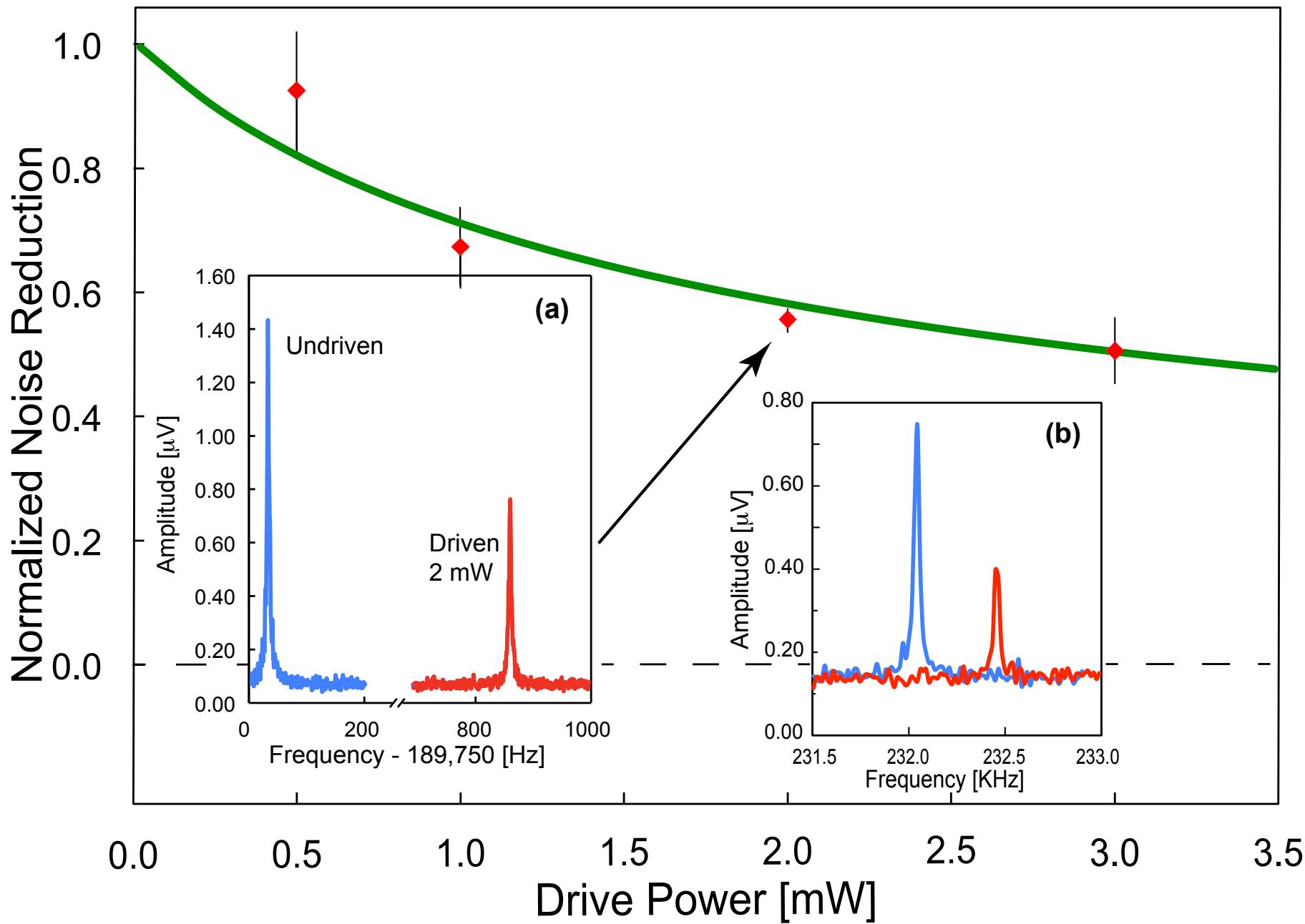
$$\delta\theta = \frac{F_{th}/I}{\tilde{\gamma}\sqrt{(\kappa - A_0 \cos 2\theta_{ss})/I}}$$



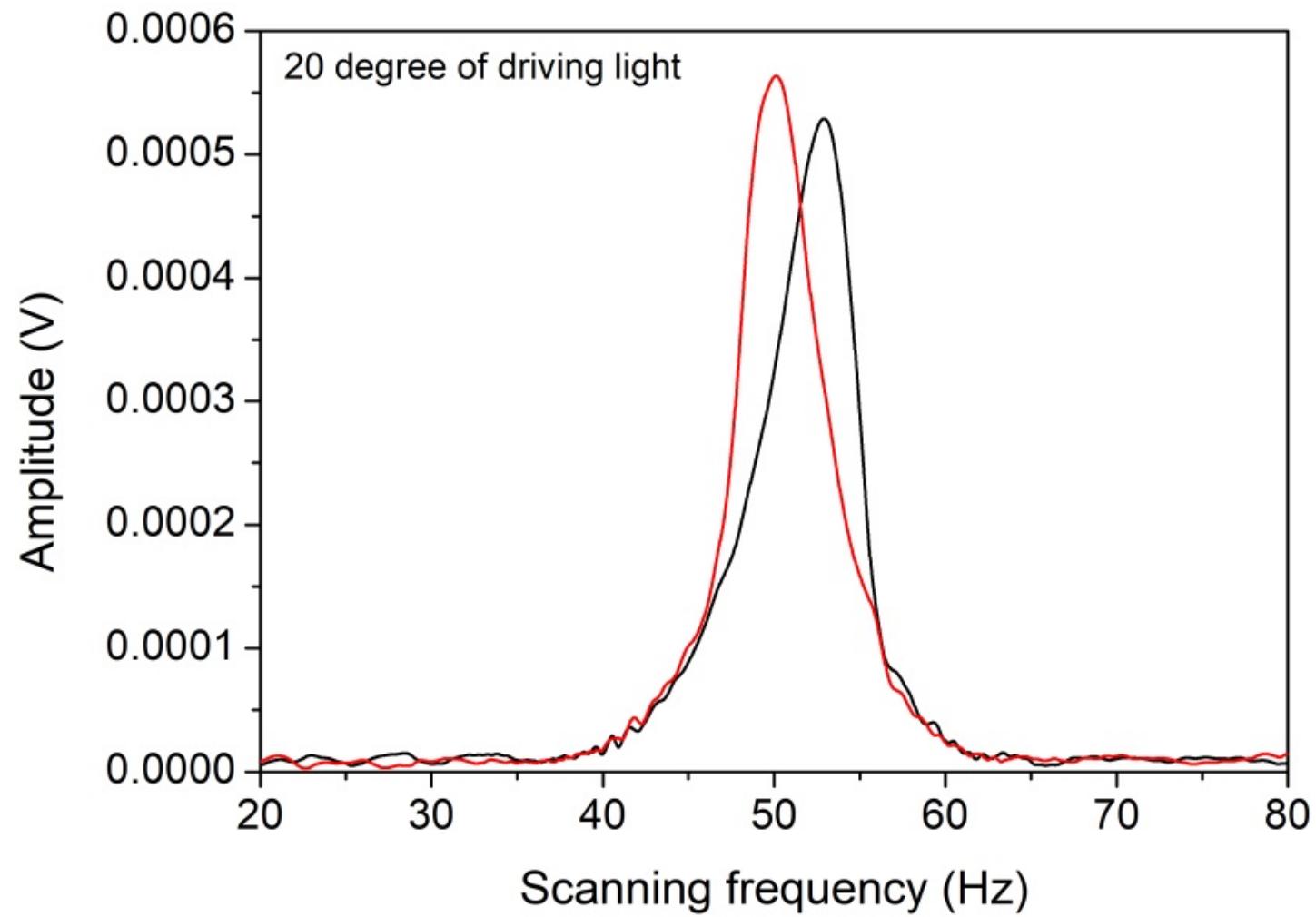
Results

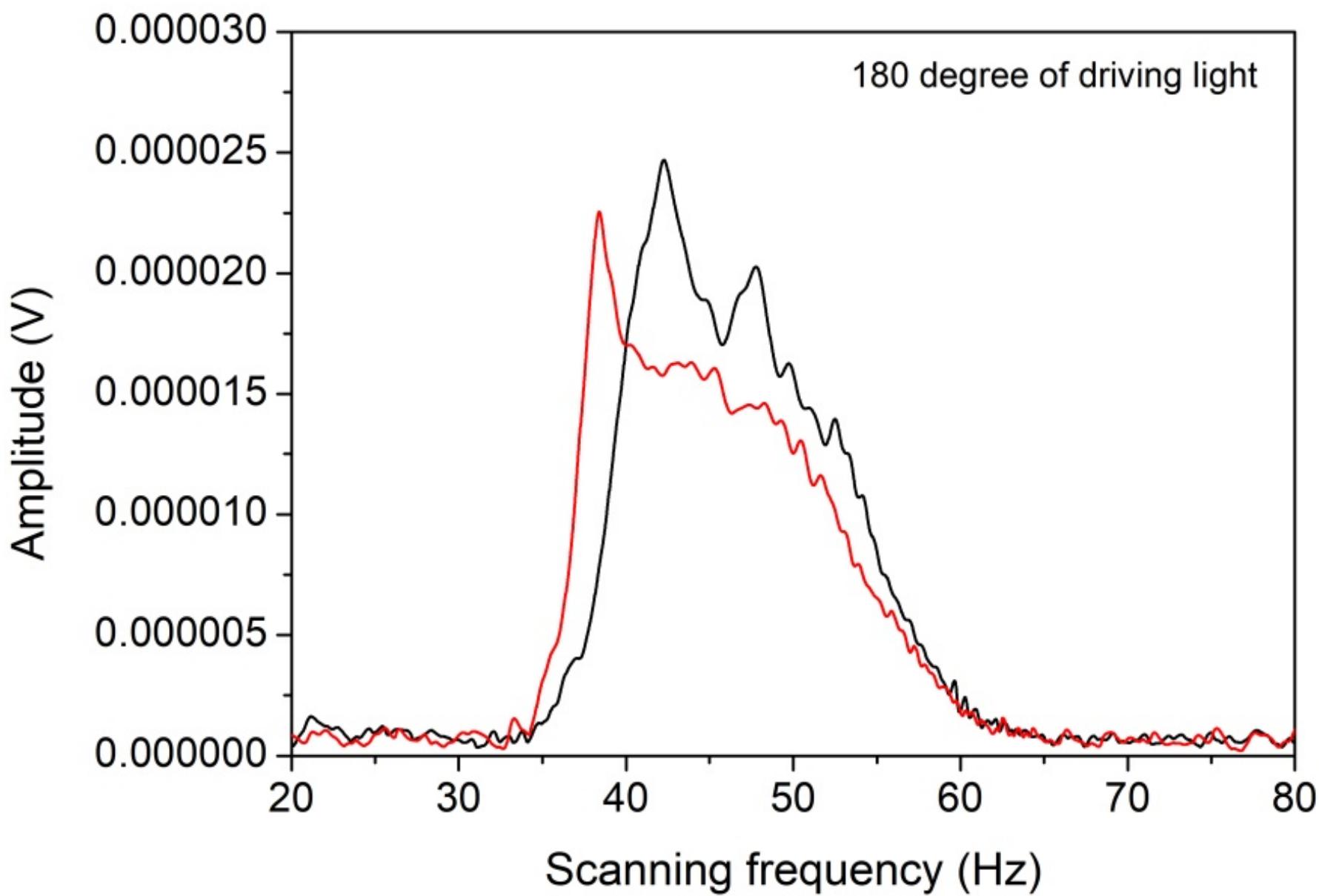


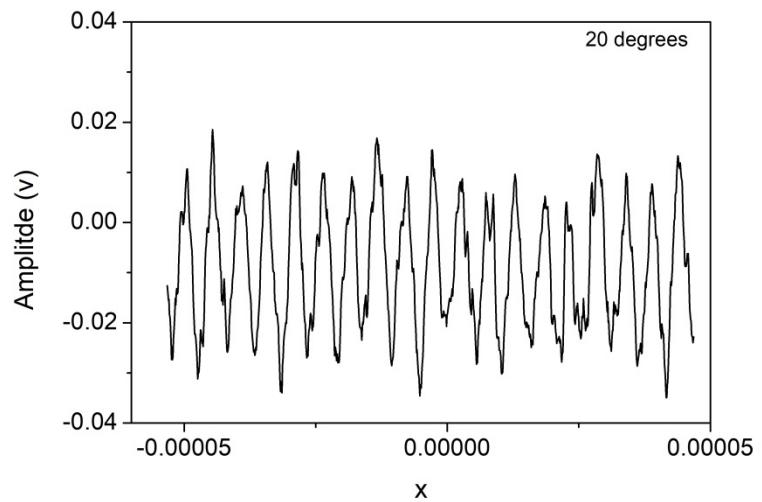
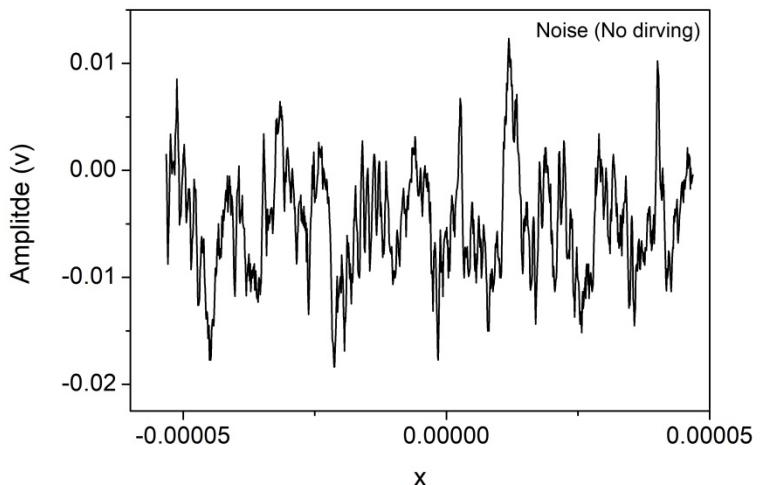
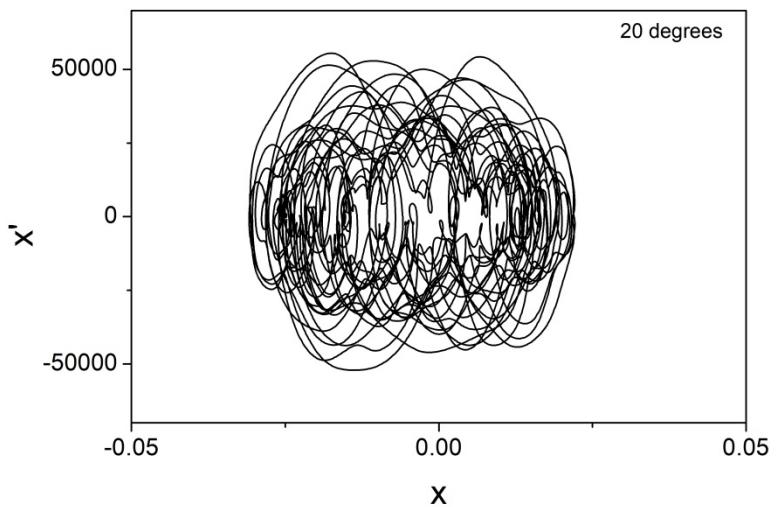
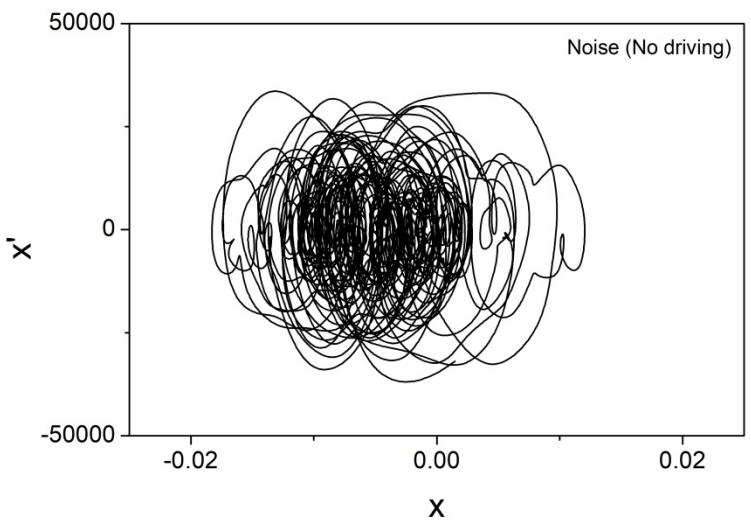




Forcing the resonance







Summary

- We have observed reduction of the thermal noise with the presence of a drive. The reduction depends on the angle of the polarization of the drive.
- The reduction of noise is on many torsional modes and it can be greater than a factor of two.
- $k_B T/h\nu = 10^{-8}$ but the quality factor is at least $\approx 10^4$ so it should be possible to think at some point of quantum effects.
- There is a lot of room for applications.

Thanks