

$$\rightarrow T_{\alpha\beta} = (\rho + p) u_\alpha u_\beta + p g_{\alpha\beta}$$

• vacuum: $T_{\alpha\beta} = -\rho_v g_{\alpha\beta}$, $\rho_v = \text{vac energy density}$

$$= T_{\alpha\beta} u^\alpha u^\beta$$

$$= \text{const.}$$

← How do I know this?

$$T_{\alpha\beta} = \frac{1}{4\pi} [F_{\alpha\mu} F_\beta{}^\mu - \frac{1}{4} g_{\alpha\beta} F_{\mu\nu} F^{\mu\nu}]$$

$$\left. \begin{aligned} \frac{E}{\Phi}{}^\alpha{}_\alpha &= -4\pi G T_{\mu\nu} u^\mu u^\nu \\ &= -R_{\mu\nu} u^\mu u^\nu \end{aligned} \right\} \Rightarrow \boxed{R_{\mu\nu} = 4\pi G T_{\mu\nu}} \quad \text{Guess \#1.}$$

Does not conserve energy consistently.

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\partial_t \vec{B}$$

$$\vec{\nabla} \times \vec{B} = \vec{J} 4\pi$$

$$\text{But Maxwell knew } \vec{\nabla} \cdot (\vec{\nabla} \times \vec{f}) = 0$$

(mix partials commute, reason why anything is true).

So $\vec{\nabla} \cdot \vec{J} = 0$, but this is not true, only true

in steady state. In general, $\vec{\nabla} \cdot \vec{J} = -\partial_t \rho$

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho, \quad \vec{\nabla} \cdot \vec{J} = -\frac{1}{4\pi} \vec{\nabla} \cdot (\partial_t \vec{E}).$$

$$\vec{\nabla} \times \vec{B} = 4\pi \vec{J} + \partial_t \vec{E}$$

Relativistic notation for charge conservation & ME.

$$F_{[\alpha\beta, \gamma]} = 0 \quad \text{identity if } F_{\alpha\beta} = 2 \partial_{[\alpha} A_{\beta]}$$

$$\partial_\alpha F^{\alpha\beta} = 4\pi j^\beta \Rightarrow \text{charge conservation } \partial_\beta \partial_\alpha F^{\alpha\beta} = 0$$

$\Rightarrow$ consistency of $F^{\alpha\beta}$ w/ charge cons.

Energy-momentum conservation $\nabla_\alpha T^{\alpha\beta} = 0$, local equation; not global total

Then it'd better be true that $\nabla_\alpha R^{\alpha\beta} = 0$, but it is not!

But I can not modify it in a way that it messes up the Newtonian limit.

$$\nabla^\alpha (R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R) \equiv 0, \quad \text{Bianchi Identity.}$$

$$\nabla^\alpha G_{\alpha\beta} = 0 \quad \text{for any metric.}$$

Answer: $R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R = 8\pi G T_{\alpha\beta}$

$$R - \frac{1}{2} \cdot 4 \cdot R = 8\pi G T$$

$$R = -8\pi G T$$

I knew

$$R_{\alpha\beta} = 4\pi G T_{\alpha\beta}$$

$$R_{\alpha\beta} = 8\pi G T_{\alpha\beta} + \frac{1}{2} g_{\alpha\beta} (-8\pi G T)$$

$$= 4\pi G [2T_{\alpha\beta} - g_{\alpha\beta} T]$$

So we need to compare $T_{\alpha\beta} = 2T_{\alpha\beta} - g_{\alpha\beta} T$

for non-relativistic $T_{\alpha\beta} = \begin{pmatrix} \rho & 0 \\ 0 & 0 \end{pmatrix}$, so $T = -\rho$

$$\text{So } 2T_{00} - g_{00} T = 2\rho - (-1)(-\rho) = \rho = T_{00}$$

Dec 9th

Nature of Einstein Eqs.

Some of them are constraint eqs.

$$\text{Tr } \ddot{\Phi}_i^i = -4\pi G \rho_{\text{mass}}$$

$$\text{Tr} \rightarrow -R_{\mu\nu} u^\mu u^\nu,$$

$$\rho_{\text{mass}} \rightarrow P_{\text{ENERGY}} \rightarrow \begin{cases} T_{ab} u^a u^b \\ -T = -T_{ab} g^{ab} \end{cases}$$

These two are close, so in guessing, some linear combination. Energy & momentum conservation tells us what appears.

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Conservation of energy: $\nabla^\mu T_{\mu\nu} = 0$.

The field equation does not do this for me, with only $R_{\mu\nu}$.

$$\nabla^\mu R_{\mu\nu} = \frac{1}{2} \nabla_\nu R \leftarrow \text{contracted w/ Bianchi}$$

$$R - \frac{1}{2} R \cdot 4 = 8\pi G T$$

$$R = -8\pi G T$$

$$\begin{aligned} R_{\mu\nu} &= 8\pi G T_{\mu\nu} + \frac{1}{2} (-8\pi G T) g_{\mu\nu} \\ &= 4\pi G (2T_{\mu\nu} - T g_{\mu\nu}) \end{aligned}$$

$T_{\mu\nu}$ for a perfect fluid: $-T_{\mu\nu} = \rho u_\mu u_\nu + p (g_{\mu\nu} + u_\mu u_\nu)$

$$T = -\rho + 3p$$

$$(2T_{\mu\nu} - T g_{\mu\nu}) u^\mu u^\nu = 2\rho + (-\rho + 3p) = \rho + 3p$$

energy density

$$\begin{pmatrix} p & & \\ & p & \\ & & p \end{pmatrix}$$

$$T_{\alpha\beta} = \frac{1}{4\pi} [F_{\alpha\mu} F_{\beta}{}^{\mu} - \frac{1}{4} g_{\alpha\beta} F_{\mu\nu} F^{\mu\nu}]$$

$$T = \frac{1}{4\pi} [F_{\alpha\mu} F^{\alpha\mu} - \frac{1}{4} \cdot 4 F_{\mu\nu} F^{\mu\nu}] = 0.$$

Taking trace, $-\rho + 3p$, that should be 0, tells us that $p = \frac{1}{3}\rho$ for thermal radiation.

What it means for "source of gravity of attraction"

How does this mean source of gravity? $\rightarrow (2T_{\mu\nu} - T g_{\mu\nu}) u^{\mu} u^{\nu} = \rho + 3p = 2\rho$

We are on a roll! So what about "vacuum energy"?

$$T_{\mu\nu} = -p_r g_{\mu\nu}$$

- Locally Lorentz Inv't, $-T_{\mu\nu} u^{\mu} u^{\nu} = \rho$

But here $p = -\rho$ from perfect fluid form $T_{\mu\nu} = \rho u_{\mu} u_{\nu} + p(g_{\mu\nu} + u_{\mu} u_{\nu})$

So the source of attraction $= \rho + 3p = -2\rho_r$.

This is "anti-gravity!"

one technical thing: contracted Bianchi identity

Bianchi identity $R_{\mu\nu}[\alpha\beta; \gamma] = 0$

In a local inertial coordinate system

$$R_{\mu\nu}[\alpha\beta; \gamma]$$

$$= \frac{1}{2} \partial_{[\mu} g_{\nu]}[\alpha, \beta], \gamma$$

holy God!
anti-symmetrizer on two partials.

der of der of metric, it could be a problem.

$$R = \partial\Gamma + \Gamma\Gamma$$

$$R_{,\alpha} = \partial\partial\Gamma + \partial\Gamma\Gamma +$$

$$R_{\mu\nu\alpha\beta; \gamma} + R_{\mu\nu\gamma\alpha; \beta} + R_{\mu\nu\beta\gamma; \alpha} = 0$$

$$\begin{matrix} \mu \leftrightarrow \alpha \\ \nu \leftrightarrow \beta \end{matrix}$$

$$R_{,\gamma\alpha} + -R_{\nu\alpha; \gamma} - R_{\mu\gamma; \alpha} = 0$$

$$R_{;\gamma} = 2 R_{\mu\alpha; \gamma} \rightarrow 2 \nabla^{\mu} R_{\mu\nu} = \nabla_{\nu} R.$$

contracted Bianchi.

of Eq and # of Things to Be Determined

$G_{\mu\nu} = 8\pi T_{\mu\nu}$, $g_{\mu\nu}$ has 10 quantities, 4 coordinate freedom.

4/10 of E.E. must not specify independent

Indeed, $G_{0\mu} = 8\pi T_{0\mu}$ involves no 2nd deri of $g_{\mu\nu}$.

$$\nabla_\mu G^{\mu\nu} = 0 \rightarrow \partial_\mu G^{\mu\nu} + \Gamma G = 0$$

$$\rightarrow \underbrace{\partial_0 G^{0\nu}}_{\substack{\text{no } \partial_t^2 \\ \text{?}}} + \underbrace{\partial_i G^{i\nu}}_{\substack{\text{no } \partial_t^3 \\ \partial_i \partial^2}} + \underbrace{\Gamma G}_{\substack{\text{no } \partial^3}} = 0$$

Initial Data

$$g_{\mu\nu} \mid g_{\mu\nu,\alpha} \mid$$

so this must not have ∂_t^2 in it!
because the entire equation adds to zero.

(How come "no ∂_t^2 " $\Rightarrow$ "does not evolve"?)

$\partial_\mu F^{\mu\nu} = 4\pi J^\nu$, must not determine all 4 components! Gauge Inv.

$$\partial_\mu F^{\mu 0} = \partial_0 F^{00} + \underbrace{\partial_i F^{i0}}_{\substack{\uparrow \\ \text{at most } \partial_t}}$$

Gauss' Law $\vec{\nabla} \cdot \vec{E} = 4\pi \rho$, does not evolve, but constrains initial data. This equation has only one time derivative of A^μ .

Since $\vec{\nabla} \cdot \vec{E} = 4\pi \rho$ is independent of time, so this constraint is satisfied for all times.

[Take $\partial_0 (G_{0\mu} - 8\pi T_{0\mu})$ using $\nabla^\mu G_{\mu\nu} = 0$]

$$\begin{aligned} & \partial_0 G_{0\mu} - 8\pi \partial_0 T_{0\mu} \\ &= -\partial_i G^{i0} - \Gamma G + 8\pi \partial_i T_{i0} \end{aligned}$$

For numerical GR, I need initial data that satisfy $G_{0\mu} = 8\pi T_{0\mu}$.

The Action Principle

use Lagrangian, vary paths, variation of action is 0 $\Rightarrow$ field eq.

Einstein-Hilbert Action

$$S[g] = \underbrace{\int d^4x \sqrt{-\det g}}_{\text{proper volume}} R \frac{1}{16\pi G}, \quad \frac{\delta S}{\delta g^{\mu\nu}} \propto R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}$$

$$S[g] = \int d^4x \sqrt{-\det g} R + S[\text{matter fields}, A_\mu]; \quad \frac{\delta R_{\mu\nu}}{\delta g_{\alpha\beta}} \Rightarrow \text{pure divergence.}$$

$$\frac{\delta S}{\delta g^{\mu\nu}} = 0 \Rightarrow \text{Einstein Eq.} \quad \text{"God is merciful!"}$$