

## 1. Sampling Theorem

Exact reconstruction of a continuous-time signal from its discrete-time samples is possible if the signal is band-limited and the sampling frequency is greater than twice the signal bandwidth.

Consider a signal  $f(t)$  whose Fourier transform  $\tilde{f}(\omega)$  is zero for  $|\omega| > \Omega$ ,

$$f(t) = \int_{-\Omega}^{\Omega} \tilde{f}(\omega) e^{-i\omega t} d\omega. \quad (1)$$

This is called a *band-limited* signal. Evaluating (1) at the discrete times  $t = nt_s$ , where the *sampling time*  $t_s$  is defined by  $t_s = \pi/\Omega$ , yields

$$f(nt_s) = \int_{-\Omega}^{\Omega} \tilde{f}(\omega) e^{-in\pi\omega/\Omega} d\omega. \quad (2)$$

The right hand side of (2) is recognized as  $2\Omega$  times the  $n$ th coefficient in the Fourier series for  $\tilde{f}(\omega)$ . Being limited to the finite range  $-\Omega < \omega < \Omega$ , the function  $\tilde{f}(\omega)$  is determined by its Fourier series coefficients, and therefore by the discrete “samples”  $f(nt_s)$ . The sample values thus determine  $f(t)$  via (1). The sampling frequency  $1/t_s = \Omega/\pi$  is twice the bandwidth  $\Omega/2\pi$ .

Show that  $f(t)$  can be reconstructed explicitly from the samples  $f(nt_s)$  via

$$f(t) = \sum_{n=-\infty}^{\infty} f(nt_s) \frac{\sin(\Omega t - n\pi)}{\Omega t - n\pi}. \quad (3)$$

## 2. Electron tunneling between neighboring layers: an application of Dirac delta functions

When electrons tunnel (make transitions) between two closely spaced semiconducting layers in the presence of a magnetic field in the  $y$  direction, the transition probability is proportional to the integral

$$D = \iint \delta\left(\frac{p_x^2 + p_y^2 - p_F^2}{2m}\right) \delta\left(\frac{(p_x - q)^2 + p_y^2 - p_F^2}{2m}\right) dp_x dp_y, \quad (4)$$

where the  $\delta$ 's are Dirac delta functions. (See Chapter 14. Section 14.3 discusses delta function of a function, which I explained in class Friday.)  $p_F$  is the *Fermi momentum*,  $m$  is the electron mass, and  $q = eBd$  is the product of electric charge  $e$ , magnetic field  $B$ , and inter-layer distance  $d$ . The factor  $D$  determines interlayer conductivity, which can be measured experimentally as a function of the magnetic field.

- (a) Make a sketch in the  $(p_x, p_y)$  plane to show geometrically where the two delta functions are non-zero for different values of  $q$ .
- (b) Show using dimensional analysis that  $D$  has dimensions of  $[m/p_F]^2$ . (This will require figuring out what are the dimensions of the delta functions. *Hint*: Think about the normalization condition defining the delta function.)
- (c) Calculate  $D$ . For practice, and to simplify writing, adopt units with  $p_F = 1$  and  $m = 1$ , and restore the factors of  $p_F$  and  $m$  at the end using dimensional analysis. (*Hint*: First do one integral, then the other.)
- (d) Sketch a plot of  $D$  as a function the dimensionless parameter  $q/p_F$ , and explain the qualitative properties of the dependence on  $q$  in terms of your sketch in part 2a. (*Guidance*: Imagine the delta functions as having a small width, before taking the limit, and consider how the area of the region in which both delta functions are non-zero depends on  $q$ .)

### 3. Wavepackets and group velocity for a relativistic quantum particle

In relativistic quantum mechanics, the (complex) wave function  $\Phi$  for a particle of mass  $m$  satisfies the partial differential equation

$$\partial_t^2 \Phi = c^2 \partial_x^2 \Phi - (m^2 c^4 / \hbar^2) \Phi, \quad (5)$$

where  $c$  is the speed of light and  $\hbar$  is Planck's constant. This is called the *Klein-Gordon* equation. For simplicity it is assumed here that the wave function depends on only one space coordinate  $x$ .

- (a) Assume  $\Phi(x, t)$  has the form of a complex exponential  $e^{-i\omega t} e^{ikx}$  with definite frequency and wave vector, and determine the dispersion relation  $\omega(k)$  between the frequency and wave vector by imposing the equation (5).
- (b) Using the quantum mechanical relation  $E = \hbar\omega$  between energy and frequency and  $p = \hbar k$  between momentum and wavevector, rewrite the dispersion relation as a relation between energy and momentum. Planck's constant should not appear. (See problem 4 of Exam 1 if you want to check yourself.)
- (c) Write an expression for a completely general positive frequency ( $\omega > 0$ ) wavepacket solution to (5). What free function determines your general solution?
- (d) What is the group velocity of a purely positive frequency wavepacket composed of a narrow range of wavevectors centered on some value  $k_0$ ? Does it depend on  $k_0$ ?
- (e) Now suppose the particle is massless,  $m = 0$ . What is then the group velocity of a purely positive frequency wavepacket composed of a narrow range of wavevectors centered on some value  $k_0$ ? Does it depend on  $k_0$ ?