

1. The relation between the real Fourier coefficients for the sine and cosine terms are obtained with the help of the following identities:

$$\int_{-\pi}^{\pi} \cos(m\theta) \cos(n\theta) d\theta = \pi \delta_{mn} \quad (1)$$

$$\int_{-\pi}^{\pi} \sin(m\theta) \sin(n\theta) d\theta = \pi \delta_{mn} \quad (2)$$

$$\int_{-\pi}^{\pi} \cos(m\theta) \sin(n\theta) d\theta = 0, \quad (3)$$

where  $m$  and  $n$  are assumed to be positive integers. Prove these identities by expressing the cosine and sine in terms of complex exponentials, and using  $\int_{-\pi}^{\pi} e^{ik\theta} d\theta = 2\pi \delta_{k0}$ . (These are equivalent to eqns (15.3-6) in the textbook.

2. Consider the “rectified cosine function” defined by

$$f(x) = \cos(\pi x/2L), \quad L \leq x \leq L, \quad (4)$$

and continued periodically so that  $f(x + 2L) = f(x)$ .

- (a) Sketch the function  $f(x)$  over several periods.
  - (b) Use the symmetry to explain why the Fourier coefficients  $b_n$  vanish.
  - (c) Find the non-vanishing Fourier coefficients. (*Hints:* (i) To clean things up, change variables to  $\theta = \pi x/L$ . (ii) You’ll need to do a probably unfamiliar integral, which you can look up or work out for yourself.)
  - (d) Using a computer program (Mathematica, Maple, Matlab, or something else) plot the sum of the first few terms in the Fourier series, together with (4), for  $\theta \in (-3\pi, 3\pi)$ . Show the result with 1 (just the constant part), 2, 5, and 50 terms included. With 5 terms the sum should already be very close to (4), except near the zeros where the slope is discontinuous.
3. Let’s go back to section 11.5, *Explosion of a nuclear bomb*, and see how Fourier series can be used to evaluate the time development of a general the neutron density function  $N(r, t)$ . (Actually, because of the  $1/r$  factor it’s not quite a Fourier series, but you’ll see the idea is very close.)

In hw7, we assumed a factored form  $N(r, t) = F(r)H(t)$ , and found the equations satisfied by  $F(r)$  and  $H(t)$ . Then we wrote  $F(r) = f(r)/r$  and found that  $f(r)$  must be a sine function. After applying the boundary conditions  $f(0) = 0 = f(R)$  the solution took the form

$$N_n(r, t) = A_n \exp(\mu_n t) \sin(k_n r)/r, \quad (5)$$

where  $n$  is a positive integer,  $A_n$  is an arbitrary constant,  $k_n = n\pi/R$ , and  $\mu_n$  is determined by the diffusion constant  $\kappa$ , the production rate  $\lambda$ , the radius of the sphere  $R$  and the integer  $n$ . A general solution is a linear combination of such solutions,  $N(r, t) = \sum_n N_n(r, t)$ , with different values of the constants  $A_n$ .

Consider the case when the neutron density  $N(r, 0)$  at time  $t = 0$  is equal to  $\bar{N}$  inside a sphere of radius  $a < R$  and zero outside. Find the values of the coefficients  $A_n$  in this case. Given these,  $N(r, t)$  is known for all time. Write out this function explicitly.

*Hint:* To evaluate the coefficients  $A_n$ , I suggest you multiply  $N(r, 0)$  by  $r \sin(k_n r)$  and integrate over  $r$  from 0 to  $R$ . Using the given initial density  $N(r, 0)$  you'll get one value, and using the series expansion you'll encounter integrals very close to (15.5) in the textbook, with  $L$  replaced by  $R$  and with the range of integration cut in half. The latter will be proportional to  $A_n$ .