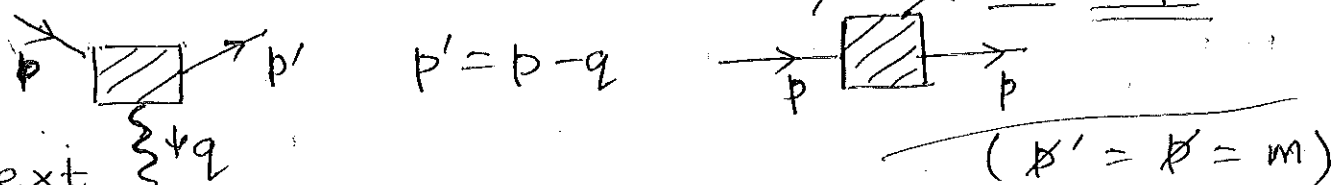


# Proof of $Z_1 = Z_2$

- Start with the general (i.e., fermions & photon can be off-shell) Ward-Takahashi <sup>(WT)</sup> identity:

$$q_\mu \Pi^{\mu, loop} = Q [\Sigma(p') - \Sigma(p)] \dots (1)$$



Next

- Take the fermions to be on-shell, where

$$\Pi_\mu^{loop}(p, p') = \gamma_\mu F_1^{loop}(q^2) + i\sigma_{\mu\nu} q^\nu F_2^{loop}(q^2) \dots (2)$$

(as shown in LP sec. 11.1 or PS sec. 6.2) <sup>assuming parity invariance (no  $\gamma_5$ )</sup>

- Furthermore, take limit of small momentum transfer between fermions, i.e.,  $q^\nu \rightarrow 0$

(as when measuring charge in a static electric field: see LP sec. 11.2.1)

- So, LHS of Eq. (1) becomes  $q_\mu \gamma^\mu F_1^{loop}(0) \dots (3)$

- On RHS of Eq. (1), do "Taylor expansion" [given  $(p'-p)^\nu \rightarrow 0$ ]:

$$Q \times \frac{\partial \Sigma}{\partial \not{p}} \Big|_{\not{p}=m} \times (-\not{q}) \dots (4)$$

- So far, we are at level of loop calculation.

Next, relate coefficients of counterterms to above loop result:

- Eq. (3) becomes  $Q(-Q)(Z_1 - 1)$  (see other note <sup>see LP Eq. 12.85</sup> on choice of  $Z_1$ )

- Eq. (4) becomes  $\oplus Q \not{q} \frac{\partial \Sigma_{CT}}{\partial \not{p}} \Big|_{\not{p}=m}$  given the choice  $\frac{\partial (\Sigma + \Sigma_{CT})}{\partial \not{p}} \Big|_{\not{p}=m} = 0$

$$= +Q \not{x} a_{CT} \quad \left( \text{since } \Sigma_{CT} = a_{CT} \not{x} + b_{CT} \right) \quad \leftarrow \text{see LP Eq. 12.84}$$

$$= -Q \not{x} (z_2 - 1) \quad \left( \text{see just below } \frac{\text{independent of } p}{\text{LP Eq. 12.84}} \right)$$

- Equating these final LHS & RHS gives

$$z_1 = z_2$$

- In HW, you are to calculate  $z_{1,2}$  explicitly (as opposed to general argument above), i.e., using

(i) DIMREG : you should indeed find  $z_1 = z_2$  since this regularization preserves gauge invariance (hence WT identity)

(ii) hard cut-off : it does not respect gauge invariance so that  $z_1 \neq z_2$  (actually, divergent pieces are equal, but not finite terms)