

①

Deep Inelastic Electron-Proton Scattering

(for cross-section)

from general formula to more convenient (and specific to QED) one (based on Peskin & Schroeder: section 17.3)

— Recall that the general formula for cross-section for electron-proton inelastic scattering is

$$\sigma [e^-(k) \overset{\text{proton}}{p}(P) \rightarrow e^-(k') + \overset{\text{any hadronic state}}{X}]$$

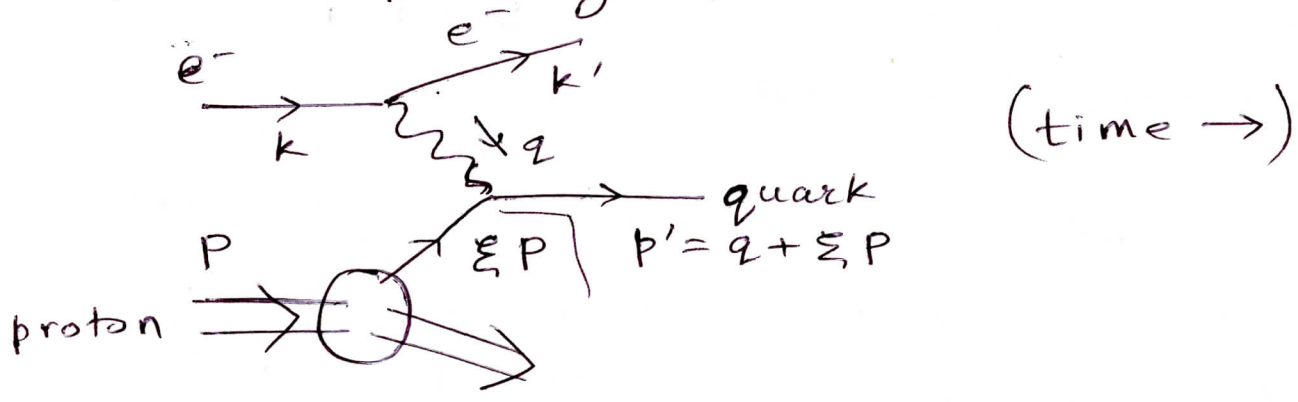
$$= \int_0^1 d\xi \sum_f f_f(\xi) \sigma [e^-(k) q_f(\xi P) \rightarrow e^-(k') + q_f(P')]$$

over all constituents of proton (in general), but for QED (photon exchange) over quarks and antiquarks only (at leading order)

where $f_f(\xi) d\xi =$ probability of finding constituent

f with longitudinal fraction ξ (f_f is parton distribution function or PDF)

— The corresponding "picture" is



— Cross section for parton-level process is

$$\frac{d\sigma}{d\hat{t}} (e^- q \rightarrow e^- q) = \frac{2\pi\alpha^2 Q_f^2}{\hat{s}^2} \left[\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} \right]$$

borrowing result of electron-muon scattering: see

Peskin & Schroeder, section 5.4 or HW10.2 from Phys 624

— Here Q_f is electric charge of quark and $\hat{s}, \hat{t}, \hat{u}$ are the Mandelstam variables (see Peskin & Schroeder, section 5.4 or Lahiri & Pal, section 7.6) for $2 \rightarrow 2$ parton-level process, i.e.,

$$\hat{s} = (k + \underbrace{\xi P}_{\text{quark momentum}})^2 \approx 2 \xi P \cdot k = \xi S \quad \rightarrow \text{at proton-level}$$

(neglecting all masses)

$$\hat{t} = (k' - k)^2 \equiv q^2 \equiv -Q^2$$

Using $\hat{s} + \hat{t} + \hat{u} = 0$, we have $\hat{u} = Q^2 - \xi S \Rightarrow$

$$\frac{d\sigma}{dQ^2} = \int_0^1 d\xi \sum_f f_f(\xi) Q_f^2 \frac{2\pi\alpha^2}{Q^4} \left[1 + \left(\frac{1 - Q^2}{\xi S} \right)^2 \right] \theta(\xi S - Q^2)$$

assume $Q^2 \gg (1 \text{ GeV})^2$

due to constraint:
 $\hat{s} \geq |\hat{t}|$

... upto $O[\alpha_s(Q^2)] \ll 1$
corrections from hard
gluon exchange/emission

— Note that the \hat{s}^2 term in $[\hat{s}^2 + \hat{u}^2]$ ^{factor} for parton-level cross-section comes from LH electron scattering from LH quark and RH... RH... whereas \hat{u}^2 is from LH electron scattering from RH quark and RH... LH...
(obviously, we have summed/averaged over spins above)

— Finally, cross-section for $e^- \bar{q} \rightarrow e^- \bar{q}$ is same as for $e^- q \rightarrow e^- q$ (this fact has already been used in last formula)

- Next, we convert above formula into a more convenient / standard form

- First, we note that ξ is actually fixed by measurement of scattered electron momentum k' (and thus momentum transfer q):

ϵ mass of quark / parton

$$0 \approx (\xi P + q)^2 = 2 \xi P \cdot q + q^2 + \underbrace{(\xi P)^2}_0 \Rightarrow$$

4-momentum of scattered quark

$$\xi = \frac{Q^2}{2 P \cdot q} \equiv \boxed{x} \text{ (clearly, } x \text{ is directly measurable / observable)}$$

- Express above formula for $d\sigma/dQ^2$ as doubly differential cross-section in x and Q^2 :
 then it is a simple product of parton-level cross section and sum of PDF's evaluated at $\xi = x$ (cf. $\int d\xi$ in earlier formula)

- It is convenient to use dimensionless combination of kinematic variables: x and (instead of Q^2) $\boxed{y} \equiv \frac{2 P \cdot q}{2 P \cdot k} = \frac{2 P \cdot q}{s}$ \leftarrow momentum of proton

also directly measurable / observable

(physical picture: in proton rest frame, $y = \frac{q^0}{k^0}$, i.e., fraction of incident electron's energy that is transferred to hadronic system)

Rewriting y in terms of partonic variables (using

4-momentum of initial quark/parton = ξP : (4)

$$y = \frac{2(\xi P) \cdot (k - k')}{2(\xi P) \cdot k} = \frac{\hat{s} + \hat{u}}{\hat{s}}, \text{ i.e., } \frac{\hat{u}}{\hat{s}} = -(1-y)$$

- Kinematically allowed region: $0 \leq x, y \leq 1$

Since it is a fraction ... use above formula(e)

- Finally, $Q^2 = xys$ and $d\xi dQ^2 = dx dQ^2 = \frac{dQ^2}{dy} dx dy$
 $= xs dx dy$

(using above relations) to give

$$\frac{d\sigma}{dx dy} (e^- p \rightarrow e^- X) = \left(\sum_f x f_f(x) Q_f^2 \right) \frac{2\pi\alpha^2 s}{Q^4} [1 + (1-y)^2]$$

again, observable

- Once photon propagator factor = $1/Q^4$, which involves both x & y , is removed, the dependence on x and y factorize: former is from PDF's, latter reflects helicities of interacting particles (see note at bottom of page 2): evidence that partons involved in deep inelastic scattering were fermions

- One can't determine separately the various PDF's from just electron scattering experiments ... but deep inelastic neutrino scattering does the job: formulae are similar to above, except that only LH particles (and RH antiparticles) ...