

Reading: See syllabus

Problems to turn in (read the rest):

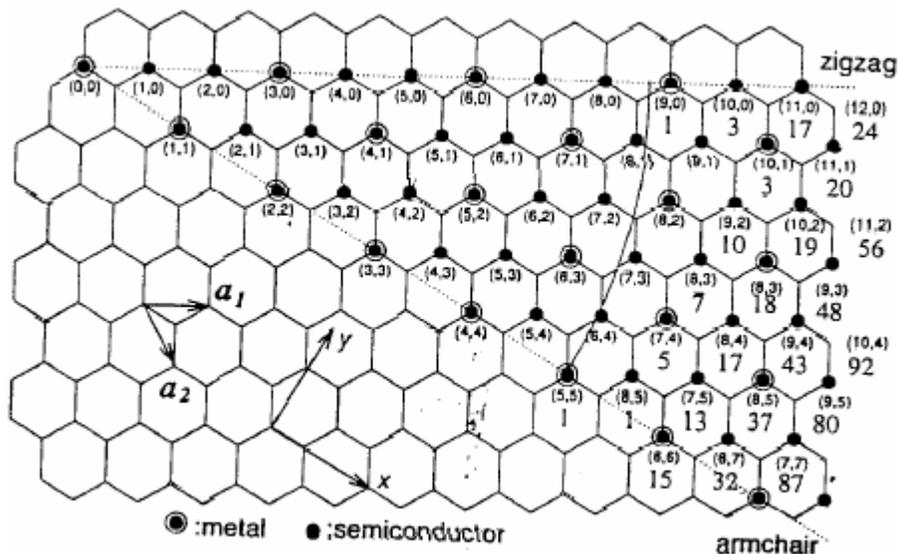
- 12-2. To make things easier, you may assume that $\mathbf{H} = H\hat{z}$ and \mathbf{M} is diagonal (but with $M_{xx} \neq M_{yy} \neq M_{zz}$). (The solutions will give the general case, which is an unpleasant mess.)
- 12-6 and 12-7a. (The pair count as a single problem. Use T_R for 12-6.)
- 13-2 is worked out on p. 407 of Grosso and Paravicini, as discussed in class (and also in the solutions). If we define kinetic coefficients, for $n = 0, 1, 2$:

$$e^2 K_n \equiv \frac{e^2}{4\pi^3} \int \tau(\hat{\mathbf{E}} \cdot \vec{v}) \varepsilon^n \left(-\frac{\partial f_0}{\partial \varepsilon} \right) d^3k = \int \left(-\frac{\partial f_0}{\partial \varepsilon} \right) \varepsilon^n \sigma(\varepsilon) d\varepsilon$$

- Show that the second equality above is true.
 - Show explicitly that $e^2 K_n = \mu^{n-1} \sigma(\mu) + ((\pi k_B T)^2 / 6) h_n(\mu)$, to order T^2 , and find $h_n(\mu)$.
 - Again to order T^2 write $e^2 K_n$ in terms of ε_F rather than μ (which, curiously, Grosso and Paravicini do not do).
4. 15-1 and 15-4 (The pair count as a single problem. Note that this is an easy extension to 3D of the 2D problem 5' of Set #7.)

Read 15-3. Since I derive this result in undergraduate modern physics courses, I presume you have already seen it. Read also 15-5, which I might have assigned had time permitted. To help visualize why γ is negative, you should look at Fig. 9.13.

- 5 (20 points). Consider the band structure of carbon nanotubes. The figure from a Dresselhaus² review, pasted below, will be helpful



Recall that in the nearest-neighbor, single-band, tight-binding model, the energy of a honeycomb is given by

$$\varepsilon(k_x, k_y) = \pm \gamma \left[1 + 4 \cos\left(\frac{\sqrt{3}k_x a}{2}\right) \cos\left(\frac{k_y a}{2}\right) + 4 \cos^2\left(\frac{k_y a}{2}\right) \right]^{1/2}$$

in Schönemberger notation, with a the lattice constant of the hexagonal Bravais net, so $\sqrt{3}$ times the nearest-neighbor (NN) spacing a_0 . This notation assumes that the points of the hexagons are in the x direction.

- a) In the armchair orientation, the wrapping is by (n,n) in terms of the primitive vectors, with the unit-cell spacing (along the circumference of the tube) being the NN spacing plus the distance across a hexagon tip-to-tip. (Cf. Fig. 19.2.) i) Show then that this unit distance $|\mathbf{a}_1 + \mathbf{a}_2|$ is 3 times the NN spacing, or $a\sqrt{3}$. By analogy with finding the allowed values of k for periodic boundary conditions, find the n allowed values of k_x here and iii) write the dispersion relation for each as a function of $k_y \rightarrow k$, the 1D wavevector in the "infinite" direction. (Only one equation need be written, with k_x replaced by its discrete values.)
- b) Draw the 2D Brillouin zone for graphene (making use of your results in earlier homework sets) and on it draw the lines corresponding to these solutions for $n = 5$ (solid lines) and $n=6$ (dashed lines).
- c) Now consider the zigzag orientation, with periodicity $(n,0)$. i) Show that the unit distance is a (so that the circumference of the tube is na) and that the wrapping can be viewed as being in the y direction. As above ii) find the allowed values of k_y and iii) write the dispersion relation.
- d) Again draw the 2D Brillouin zone for graphene with solid and dashed lines for $n= 5$ and 6 , respectively.
- e) Comment on the difference between the two and why the figure says armchair tubes are always metallic but zig-zag ones are metallic only when n is a multiple of 3.