

Reading: See syllabus

Reminder: No class on Tuesday Nov. 14, Thursdays Nov. 16 & 23.

Make-up classes: Fridays Nov. 10 & 17, 2:30-3:45.

Problems to turn in (read the rest):

1. 12-2. To make things easier, you may assume that  $\mathbf{H} = H\hat{z}$  and  $\mathbf{M}$  is diagonal (but with  $M_{xx} \neq M_{yy} \neq M_{zz}$ ). (The solutions will give the general case.)
2. 12-6 and 12-7a. (The pair count as a single problem. Use  $T_R$  for 12-6.)
3. 13-2.
4. 15-1 and 15-4 (The pair count as a single problem. Note that this is an easy extension to 3D of the 2D problem 5' of Set #7.)
5. 31-11.
6. Refer back to pp. 661-663; the spin susceptibility of a conduction electron gas at  $T = 0$  K may be discussed by another method. Let  $n_{\pm} \equiv (n/2)(1 \mp \zeta)$  be the concentration of spin-up (-down) electrons, i.e. parallel (antiparallel) to a magnetic field  $H$ .

a) Show that, in  $H$ , the total energy per volume in the spin-up band in a free-electron gas is

$$E^+ = E_0 (1 - \zeta)^{5/3} + (n/2) \mu_B H (1 - \zeta),$$

where  $E_0 = (3/10)n\varepsilon_F$  in terms of the Fermi energy  $\varepsilon_F$  in zero magnetic field. Find a similar expression for  $E^-$ .

b) Minimize the total energy per volume  $E^+ + E^-$  with respect to  $\zeta$  and solve for the equilibrium value of  $\zeta$  in the approximation  $\zeta \ll 1$ . Show then that the magnetization  $M = (3/2) n\mu_B^2 H/\varepsilon_F$ , as in the class discussion of Pauli paramagnetism.

c) We now consider the effect of exchange interactions among the conduction electrons. As a viable first approximation, we assume that electrons with parallel spins interact with each other with energy  $-V$  (with  $V > 0$ ), while electrons with antiparallel spins do not interact with each other. Show that the additional term  $-(1/8) V n^2 (1 - \zeta)^2$  is added to  $E^+$  and find a similar expression for  $E^-$ .

d) Minimize the total energy and solve for  $\zeta$  again in the limit  $\zeta \ll 1$ . Show that the magnetization is

$$M = \frac{3n\mu_B^2}{2\varepsilon_F - \frac{3}{2}Vn} H$$

Notice that there is peculiar behavior for  $V > 4\varepsilon_F/3n$ . One can easily show that at  $H=0$  the total energy for the paramagnetic state with  $\zeta = 0$  is unstable relative to a ferromagnetic state with finite  $\zeta$ . This is called the Stoner criterion for ferromagnetism. (adapted from Kittel, ISSP) This kind of phase transition at  $T=0$  is now glamorized with the label "quantum phase transition."

Read 15-3. Since I derive this result in undergraduate modern physics courses, I presume you have already seen it. Read also 15-5, which I might have assigned had time permitted. To help visualize why  $\gamma$  is negative, you should look at Fig. 9.13.