## Physics 601 Homework 6---Due Friday October 15

- 1. In class starting with the action  $S = \int d\tau (-m \mathbf{S})$  we used covariance to show that for a particle moving in a Lorentz scalar field  $\frac{d((m+\mathbf{S})u^{\mu})}{d\tau} = \partial^{\mu}\mathbf{S}$  where  $\partial^{\mu} \equiv g^{\mu\nu} \frac{\partial}{\partial x^{\nu}}$ .
  - a. Show that this can be rewritten in the form  $(m + \mathbf{S}) \frac{d(u^{\mu})}{d\tau} = \partial^{\mu} \mathbf{S} (\partial^{\alpha} \mathbf{S}) u_{\alpha} u^{\mu}$
  - b. Show that this equation of motion automatically satisfies the condition  $\frac{d\left(u_{\mu}u^{\mu}\right)}{d\tau}=0.$  This indicates that imposition of covariance yielded a self-consistent result that respects the condition  $u_{\mu}u^{\mu}=1.$
  - c. Show that in the non-relativistic limit where all of the velocities are much less than the speed of light and  $\mathbf{S} << m$  the Lagrangian for the system reduces to  $L = \frac{1}{2}m\dot{\vec{x}}^2 \mathbf{S}$  plus an irrelevant constant and the equation of motion reduces to  $m\ddot{\vec{x}} = -\vec{\nabla}\mathbf{S}$ .
- 2. Start from the action  $S = \int d\tau \left(-m + V^{\mu}u_{\mu}\right)$  where  $A_{\mu}$  is a four vector field that depends on space-time. Show that the equation of motion is  $\frac{d\left(mu_{\mu}\right)}{d\tau} = \left(\frac{\partial V_{\mu}}{\partial x^{\nu}} \frac{\partial V_{\nu}}{\partial x^{\mu}}\right)u^{\nu}.$
- 3. In electro-magnetism, one can write the scalar and vector potentials in a form that looks like a 4-vector:  $A^{\mu} = \begin{pmatrix} \Phi \\ A_x \\ A_y \\ A_z \end{pmatrix}$ . Because one can make arbitrary gauge

transformations  $A^{\mu}$  need not transform as a 4-vector.

- a. Show that a gauge transformation can be written in the form of the form  $A_{\mu} \rightarrow A'_{\mu} = A_{\mu} + \frac{\partial G}{\partial x^{\mu}}$  where G is an arbitrary function of space-time which need not transform as a 4-scalar under Lorentz transformations .
- b. A sufficient condition to show that  $A^{\mu}$  transforms as a 4-vector is to show that  $\frac{\partial A^{\mu}}{\partial x^{\mu}} = 0$  with  $|A^{\mu}| \to 0$  as  $|\vec{x}| \to \infty$  Explain briefly why.
- c. Show that it is always possible to make a gauge transformation (i.e. to choose G) to ensure that  $A'^{\mu}$  does transform as a 4-vector by picking

- choosing  $\Lambda$  to satisfy the condition  $\partial_{\mu}A^{\mu}=-\partial_{\mu}\partial^{\mu}G$ . This is called the Lorentz gauge.
- d. Show that the field-strength tensor  $F_{\mu\nu} \equiv \partial_{\mu}A_{\nu} \partial_{\nu}A_{\mu}$  is gauge invariant. That is show that  $F'_{\mu\nu} \equiv \partial_{\mu}A'_{\nu} \partial_{\nu}A'_{\mu} = F_{\mu\nu}$  for any transformation of the form given in part a.
- e. Show that  $F_{\mu\nu}$  transforms under Lorentz transformations as a 4-tensor:  $F_{\mu\nu} \to \Lambda_{\mu}^{\phantom{\mu}\alpha} \Lambda_{\nu}^{\phantom{\nu}\beta} F_{\alpha\beta}$  where  $\Lambda$  is the matrix which specifies the Lorentz transformation.
- f. Using the result problem 1, plus the preceding parts of this problem show that the equation of motion for a particle of mass m and charge q moving in an electro magnetic field is given by  $\frac{d(mu^{\mu})}{d\tau} = qF^{\mu\nu}u_{\nu}.$
- 4. In the preceding section you showed that  $\frac{d(mu^{\mu})}{d\tau} = qF^{\mu\nu}u_{\nu}$ . Staring with this equation of motion show that of necessity  $\frac{d(u_{\mu}u^{\mu})}{d\tau} = 0$ . This indicates that equation of motion self-consistently respects the condition  $u_{\mu}u^{\mu} = 1$ .
- 5. Suppose that one has a charged particle with mass m and charge q interacting with external electromagnetic fields specified by the potentials  $\Phi = 0$ ,  $A_x = -E_0 t$ ,  $A_y = A_z = 0$ . Note that potentials are independent of special position.
  - a. Verify that these potentials satisfy the Lorentz gauge condition  $\partial_{\mu}A^{\mu}=0$ .
  - b. Construct the field strength tensor  $F_{\mu\nu}$  .
  - c. Suppose that the particle starts from rest at t=0, find the position of the particle as a function of time. (Hint, you may find your solution of problem 4.d of homework 5 to be useful.)
- 6. In class we found the Green's function for the harmonic oscillator. In this problem, I want you to find and use the analogous one for a damped oscillator. The damped driven oscillator satisfied the equation:  $m\ddot{x} + 2\beta m\dot{x} + m\omega_0^2 x = f(t)$  where  $\beta$  is a damping parameter. The solution is  $x(t) = \int_{-\infty}^{\infty} dt' G(t,t') f(t')$  where the Green's function satisfies  $(\partial_t^2 + 2\beta\partial_t + \omega_0^2)G(t,t') = \delta(t-t')$ . A useful first step in constructing this is to exploit the known solution for steady state motion with a harmonic driving force:  $(\partial_t^2 + 2\beta\partial_t + \omega_0^2)x(t) = f_0e^{i\omega t}$  has a solution of the form of the form  $x(t) = \frac{f_0e^{i\Omega(t-t')}}{\omega_0^2 \Omega^2 + 2i\beta\Omega}$ . Thus  $(\partial_t^2 + 2\beta\partial_t + \omega_0^2)\frac{e^{i\Omega(t-t')}}{\omega_0^2 \Omega^2 + 2i\beta\Omega} = e^{i\Omega(t-t')}$ . Let us now integrate both sides with respect to  $\Omega$  and divide by  $2\pi$ :

$$\frac{\int_{-\infty}^{\infty} d\Omega \left(\partial_t^2 + 2\beta \partial_t + \omega_0^2\right) \frac{e^{i\omega(t-t')}}{\omega_0^2 - \omega^2 + 2i\beta \omega}}{2\pi} = \frac{\int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')}}{2\pi}.$$
 We know that 
$$\frac{\int_{-\infty}^{\infty} d\omega e^{i\omega(t-t')}}{2\pi} = \delta(t-t').$$
 Moreover on the left hand side we can pull

 $(\partial_t^2 + 2\beta\partial_t + \omega_0^2)$  out of the integral as it does not depend on  $\omega$ . Thus

$$\left(\partial_t^2 + 2\beta \partial_t + \omega_0^2\right) \left( \frac{\int_{-\infty}^{\infty} d\omega \, \frac{e^{i\omega(t-t')}}{\omega_0^2 - \omega^2 + 2i\beta\omega}}{2\pi} \right) = \delta(t-t') \text{ and the object in the parenthesis}$$

is a Green's function.

- a. Evaluate the integral above using contour integration to find an explicit expression for G(t,t'). Note that the complex exponential implies that the  $\frac{1}{2}$  plane in which the contour is to be closed depends on the sign of t-t'.
- b. Use this Green's function to find a solution of  $m\ddot{x} + 2\beta m\dot{x} + m\omega_0^2 x = f_0 e^{-\Gamma t}\theta(t)$
- c. Consider the result in b. in the regime where  $\Gamma > \infty > \beta$ . In that regime the system should look like an underdamped oscillator getting a delta-functionlike impulse at t=0. Does it?