

PHY 402 - HW2 - SOLUTION

A. i) $\langle p | y \rangle = \int dx \langle p | x \rangle \langle x | y \rangle = \frac{e^{-ipy/\hbar}}{\sqrt{2\pi\hbar}}$

$\langle p | x \rangle$ is the eigenket of \hat{p} with eigenvalue p .
 $\langle x | y \rangle$ is the eigenket of \hat{x} with eigenvalue y .
 $\langle x | p \rangle^* = \frac{e^{-ipx/\hbar}}{\sqrt{2\pi\hbar}}$
 $\delta(x-y)$

$|y\rangle = \int dp |p\rangle \langle p | y \rangle = \int dp \frac{e^{-ipy/\hbar}}{\sqrt{2\pi\hbar}} |p\rangle$

coordinates of y in the $|p\rangle$ basis

ii) matrix element: $\langle p | \hat{x} | p' \rangle = \int dx dx' \langle p | x \rangle \langle x | \hat{x} | x' \rangle \langle x' | p' \rangle$

$\langle x | \hat{x} | x' \rangle = x' \langle x | x' \rangle = x' \delta(x-x')$

$= \int dx \langle p | x \rangle x \langle x | p' \rangle$

$\frac{e^{-ipx/\hbar}}{\sqrt{2\pi\hbar}} \frac{e^{ip'x/\hbar}}{\sqrt{2\pi\hbar}}$

$= \int dx x \frac{e^{-i(p-p')x/\hbar}}{\sqrt{2\pi\hbar}^2}$

$= i\hbar \frac{d}{dp} \int dx \frac{e^{-i(p-p')x/\hbar}}{2\pi\hbar}$

$\delta(p-p')$

$= i\hbar \frac{d}{dp} \delta(p-p')$

$$\text{iii) } \langle x | \frac{\hat{p}^2}{2m} + \frac{M\omega^2}{2} \hat{x}^2 | \psi \rangle = \int dy \langle x | \frac{\hat{p}^2}{2m} + \frac{M\omega^2}{2} \hat{x}^2 | y \rangle \underbrace{\langle y | \psi \rangle}_{\psi(y)}$$

but

$$\langle x | \hat{x} | y \rangle = \langle x | y \rangle = y \langle x | y \rangle = y \delta(x-y)$$

$$\langle x | \hat{x}^2 | y \rangle = \langle x | y^2 | y \rangle = y^2 \langle x | y \rangle = y^2 \delta(x-y)$$

$$\langle x | \hat{p} | y \rangle = -i \frac{d}{dx} \delta(x-y) \quad (\text{as shown in class})$$

$$\langle x | \hat{p}^2 | y \rangle = -i \frac{d}{dx} \langle x | \hat{p} | y \rangle = -\frac{d^2}{dx^2} \langle x | y \rangle = -\frac{d^2}{dx^2} \delta(x-y)$$

definition
of \hat{p} in our
class

$$\text{so } \langle x | \frac{\hat{p}^2}{2m} + \frac{M\omega^2}{2} \hat{x}^2 | \psi \rangle = \int dy \left[-\frac{1}{2m} \frac{d^2}{dx^2} \delta(x-y) + \frac{M\omega^2}{2} y^2 \delta(x-y) \right] \psi(y)$$

$$= -\frac{1}{2m} \frac{d^2}{dx^2} \psi(x) + \frac{M\omega^2}{2} x^2 \psi(x)$$

$$\text{B. i) } \hat{A} | \psi \rangle = \left[E_0 (|A\rangle \langle A| + |B\rangle \langle B|) + T (|A\rangle \langle B| + |B\rangle \langle A|) \right] \frac{(|A\rangle + |B\rangle)}{\sqrt{2}}$$

$$= (E_0 |A\rangle + E_0 |B\rangle + T |A\rangle + T |B\rangle) / \sqrt{2}$$

$$= (E_0 + T) \frac{(|A\rangle + |B\rangle)}{\sqrt{2}} = (E_0 + T) | \psi \rangle$$

ii) The Time-evolution is ^{more} easily computed in the eigenbasis of \hat{H} .
 Since we already used \hat{H} in class we know the eigenvalues/eigenkets

$$\hat{H} |+\rangle = (E_0 + T) |+\rangle, \quad |+\rangle = \frac{|A\rangle + |B\rangle}{\sqrt{2}}$$

$$\hat{H} |-\rangle = (E_0 - T) |-\rangle, \quad |-\rangle = \frac{|A\rangle - |B\rangle}{\sqrt{2}}$$

~~the~~ The Time evolution of $|+\rangle, |-\rangle$ in Time is trivial:

$$i\hbar \frac{d}{dt} |+, t\rangle = \hat{H} |+, t\rangle \Rightarrow |+, t\rangle = e^{-i(E_0+T)t/\hbar} |+\rangle$$

$$i\hbar \frac{d}{dt} |-, t\rangle = \hat{H} |-, t\rangle \Rightarrow |-, t\rangle = e^{-i(E_0-T)t/\hbar} |-\rangle$$

Write $|\psi\rangle$ in the $\{|+\rangle, |-\rangle\}$ basis:

$$|\psi\rangle = |+\rangle \langle +|\psi\rangle + |-\rangle \langle -|\psi\rangle$$

$$= \frac{|A\rangle + |B\rangle}{\sqrt{2}} \left(\frac{\langle A| + \langle B|}{\sqrt{2}} \right) (a|A\rangle + b|B\rangle)$$

$\swarrow a(0)=1$ $\swarrow b(0)=0$

$$+ \frac{|A\rangle - |B\rangle}{\sqrt{2}} \left(\frac{\langle A| - \langle B|}{\sqrt{2}} \right) (a|A\rangle + b|B\rangle)$$

$$= \frac{1}{2} (|A\rangle + |B\rangle) (a+b) + \frac{1}{2} (|A\rangle - |B\rangle) (a-b)$$

$$= \frac{a+b}{\sqrt{2}} |+\rangle + \frac{a-b}{\sqrt{2}} |-\rangle = \frac{1}{\sqrt{2}} |+\rangle + \frac{1}{\sqrt{2}} |-\rangle$$

↓

$$|\psi(t)\rangle = \frac{1}{\sqrt{2}} e^{-i(E_0+T)t/\hbar} |+\rangle + \frac{1}{\sqrt{2}} e^{-i(E_0-T)t/\hbar} |-\rangle$$

$$= e^{-iE_0 t/\hbar} \left[\cos \frac{TE}{\hbar} |A\rangle - i \sin \frac{TE}{\hbar} |B\rangle \right]$$

$$\text{iii) } E = E_0 + T \quad \text{w/ probability } |\langle +1 | \psi(t) \rangle|^2 = \frac{1}{2}$$

$$E = E_0 - T \quad \text{w/ probability } |\langle -1 | \psi(t) \rangle|^2 = \frac{1}{2}$$

$$\text{C. i) } |\psi\rangle = \int dp \langle p | \psi \rangle |p\rangle \Rightarrow \text{probability (density) of measuring } p = p_0 = |\langle p_0 | \psi \rangle|^2$$

$$\text{but } \langle p_0 | \psi \rangle = \int dx \underbrace{\langle p_0 | x \rangle}_{\langle x | p_0 \rangle^*} \underbrace{\langle x | \psi \rangle}_{\psi(x)} = \int dx \frac{e^{-ip_0 x / \hbar}}{\sqrt{2\pi\hbar}} \psi(x)$$

$$\text{For the ground state of the harm. oscillator } \psi(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-\frac{m\omega x^2}{2\hbar}} \quad \text{so}$$

$$\langle p_0 | \psi \rangle = \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-ip_0 x / \hbar - \frac{m\omega x^2}{2\hbar}}$$

$$= \frac{1}{\pi^{1/4}} \frac{1}{(m\omega\hbar)^{1/4}} e^{-\frac{p_0^2}{2m\hbar\omega}}$$

$$\text{so probability density} = \frac{1}{(\pi\hbar m\omega)^{1/2}} e^{-\frac{p^2}{m\hbar\omega}}$$

ii) The eigenfunction of momentum of eigenvalue p_0 , that is, $|p_0\rangle$.
In the x -basis it is:

$$\langle x | p_0 \rangle = \frac{e^{ip_0 x / \hbar}}{\sqrt{2\pi\hbar}}$$

NOTE: This is not a normalizable wavefunction so, strictly speaking, the ~~particle~~ particle can never be on this state and the momentum can never be infinitely well measured. In practice p is known w/ some experimental uncertainty and the wave function after the measurement is very broad but will tend to zero eventually as $|x| \rightarrow \infty$.

D. i)
$$e^{-\frac{i\gamma \hat{p}}{\hbar}} f(x) = e^{-\gamma \frac{d}{dx}} f(x)$$

$$= \sum_{n=0}^{\infty} \frac{(-\gamma)^n}{n!} \frac{d^n}{dx^n} f(x)$$

$$= f(x) - \gamma f'(x) + \frac{\gamma^2}{2} f''(x) - \frac{\gamma^3}{6} f'''(x) + \dots$$

$$= f(x - \gamma)$$

ii) ~~$\psi(x,t) = e^{-i\hat{H}t/\hbar} \psi(x,0)$~~

$$\underbrace{\langle x | \psi(t) \rangle}_{\text{in the } x \text{ basis}} = e^{-i\hat{H}t/\hbar} \underbrace{\langle x | \psi(0) \rangle}_{\text{in the } x \text{ basis}} \Rightarrow \underbrace{|\psi(t)\rangle}_{\text{true in any basis}} = e^{-i\hat{H}t/\hbar} |\psi(0)\rangle$$

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = i\hbar \frac{d}{dt} e^{-i\hat{H}t/\hbar} |\psi(0)\rangle = \hat{H} e^{-i\hat{H}t/\hbar} |\psi(0)\rangle = \hat{H} |\psi(t)\rangle$$

NOTE: how do I know that $\frac{d}{dt} e^{-i\hat{H}t/\hbar} = -\frac{i\hat{H}}{\hbar} e^{-i\hat{H}t/\hbar}$? \rightarrow

Using the definition of exponentials:

$$e^{-i\hat{H}t/\hbar} = \sum_{n=0}^{\infty} \frac{\left(\frac{-i\hat{H}t}{\hbar}\right)^n}{n!} = 1 - \frac{it}{\hbar} \hat{H} - \frac{t^2}{2\hbar^2} \hat{H}^2 + \dots$$

$$\begin{aligned} \frac{d}{dt} e^{-i\hat{H}t/\hbar} &= -\frac{i}{\hbar} \hat{H} - \frac{t}{\hbar^2} \hat{H}^2 + \dots = -\frac{i\hat{H}}{\hbar} \left[1 - \frac{it}{\hbar} \hat{H} + \dots \right] \\ &= -\frac{i}{\hbar} \hat{H} e^{-i\hat{H}t/\hbar} \end{aligned}$$