

QUANTUM PHYSICS I SUMMARY

(UP TO 1ST MIDTERM)

GENERAL FORMALISM

- state of the particle at a particular time described by a complex function of the position: the wave function $\Psi(x)$.
- every classical observable corresponds to an operator

position: $x \Rightarrow \hat{x}, \hat{x} \Psi(x) = x \Psi(x)$
 momentum: $p \Rightarrow \hat{p}, \hat{p} \Psi(x) = -i\hbar \frac{d}{dx} \Psi(x)$
 kinetic energy: $T = \frac{p^2}{2M} \Rightarrow \hat{T} = \frac{\hat{p}^2}{2M}, \hat{T} \Psi(x) = -\frac{\hbar^2}{2M} \frac{d^2}{dx^2} \Psi(x)$
 total energy (hamiltonian): $H = \frac{p^2}{2M} + V(x) \Rightarrow \hat{H} = \frac{\hat{p}^2}{2M} + V(\hat{x}), \hat{H} \Psi(x) = \left[-\frac{\hbar^2}{2M} \frac{d^2}{dx^2} + V(x) \right] \Psi(x)$
 :

- these operators have a set of eigenvalue/eigenfunction pairs

$$\hat{A} \psi_n(x) = a_n \psi_n(x)$$

(sometimes n 's form a discrete set, sometimes a continuous set)

Examples:

momentum: $\hat{p} \psi_p(x) = p \psi_p(x) \Rightarrow \psi_p(x) = \frac{e^{iPx/\hbar}}{\sqrt{2\pi\hbar}}, p = \text{real number}$

position: $\hat{x} \psi_{x_0}(x) = x_0 \psi_{x_0}(x) \Rightarrow \psi_{x_0}(x) = \delta(x-x_0), x_0 = \text{real number}$

hamiltonian

for free particle: $\hat{H} \psi_k(x) = E_k \psi_k(x) \Rightarrow \psi_k(x) = \frac{e^{ikx}}{\sqrt{2\pi\hbar}}, E_k = \frac{\hbar^2 k^2}{2M}$ for any real k

Hamiltonian of the infinite square well : $\hat{H} \psi_n(x) = E_n \psi_n(x) \Rightarrow \psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right), E_n = \frac{n^2 \hbar^2 \pi^2}{2mL^2}$
 $\frac{\hat{p}^2}{2m} + V(x)$
 \uparrow
 0 for $0 < x < L$
 ∞ otherwise
 for $n=1, 2, \dots$

in all these examples the eigenfunctions satisfy an orthonormality condition:

discrete spectrum $\int_{-\infty}^{\infty} dx \psi_n^*(x) \psi_m(x) = \delta_{nm} = \begin{cases} 1 & \text{for } n=m \\ 0 & \text{for } n \neq m \end{cases}$
 OR

continuous spectrum $\int_{-\infty}^{\infty} dx \psi_p^*(x) \psi_{p'}(x) = \delta(p-p')$

- The wave function evolves in time according to the (time-dependent) Schrödinger equation:

$$i\hbar \frac{\partial}{\partial t} \psi(x,t) = \underbrace{-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x,t) + V(x) \psi(x,t)}_{\hat{H} \psi(x,t)}$$

The general solution of the Schrödinger eq. is given by

$$\psi(x,t) = \sum_n c_n e^{-iE_n t/\hbar} \psi_n(x),$$

where $E_n, \psi_n(x)$ are the eigenvalues/eigenfunctions of \hat{H}

$$\hat{H} \psi_n(x) = E_n \psi_n(x)$$

The c_n 's are determined by the initial condition

$$\psi(x,0) = \sum_n c_n \psi_n(x) \Rightarrow c_n = \int_{-\infty}^{\infty} dx \psi_n^*(x) \psi(x,0)$$

- the only outcomes of a measurement of an observable $A(x,p)$ at time t are the eigenvalues of \hat{A} . The probability of getting one particular eigenvalue a_n is given by $|c_n|^2$ where c_n is the coefficient of the expansion of Ψ in terms of ψ_n 's.

$$\hat{A} \psi_n(x) = a_n \psi_n(x)$$

PROBABILITY OF GETTING THIS EIGENVALUE IS THE MAGNITUDE SQUARE OF THIS

$$\Psi(x,t) = \sum_n c_n \psi_n(x) \Rightarrow c_n = \int_{-\infty}^{\infty} dx \psi_n^*(x) \Psi(x,t)$$

Examples:

- 1) measurement of energy on a infinite square well when the wave function is

$$\Psi(x,t) = \underbrace{\sqrt{\frac{3}{5}}}_{c_1} \underbrace{\sqrt{\frac{2}{L}} \sin \frac{\pi x}{L}}_{\psi_1(x)} - i \underbrace{\sqrt{\frac{2}{5}}}_{c_3} \underbrace{\sqrt{\frac{2}{L}} \sin \frac{3\pi x}{L}}_{\psi_3(x)}$$

$$\text{probability for } E_1 = |c_1|^2 = 3/5$$

$$E_1 = \frac{\hbar^2 \pi^2}{2mL^2}$$

$$\text{probability for } E_3 = |c_3|^2 = 2/5$$

$$E_3 = 9 \frac{\hbar^2 \pi^2}{2mL^2}$$

$$\text{any other } E_n = |c_n|^2 = 0 \quad (n \neq 1, 3)$$

2) measurement of momentum
when the wave function
is $\Psi(x,t) = \psi(x)$

$$\psi(x) = \int_{-\infty}^{\infty} dp \frac{e^{i \frac{p}{\hbar} x}}{\sqrt{2\pi\hbar}} c(p) \Leftrightarrow c(p) = \int_{-\infty}^{\infty} dx \frac{e^{-i \frac{p}{\hbar} x}}{\sqrt{2\pi\hbar}} \psi(x) = \text{"Fourier Transform of } \psi(x)\text{"}$$

$\psi(x)$ written as a
linear combination of momentum
eigenfunctions

probability (density)
of measuring p is

$$= |c(p)|^2 = \left| \int_{-\infty}^{\infty} dx \frac{e^{-i \frac{p}{\hbar} x}}{\sqrt{2\pi\hbar}} \psi(x) \right|^2$$

3) measurement of position
when the wave function
is $\Psi(x,t) = \psi(x)$

$$\psi(x) = \int_{-\infty}^{\infty} dy c(y) \delta(x-y) \Leftrightarrow c(y) = \int_{-\infty}^{\infty} dx \delta(x-y) \psi(x) = \psi(y)$$

$\psi(x)$ written as a
linear combination of position
eigenfunctions

probability (density)
of finding the particle at y

$$= |c(y)|^2 = |\psi(y)|^2$$

It follows from the rule above that the average value of the
measurement of A is given by

$$\langle A \rangle = \int_{-\infty}^{\infty} dx \psi^*(x,t) \hat{A} \psi(x,t)$$