

Energy

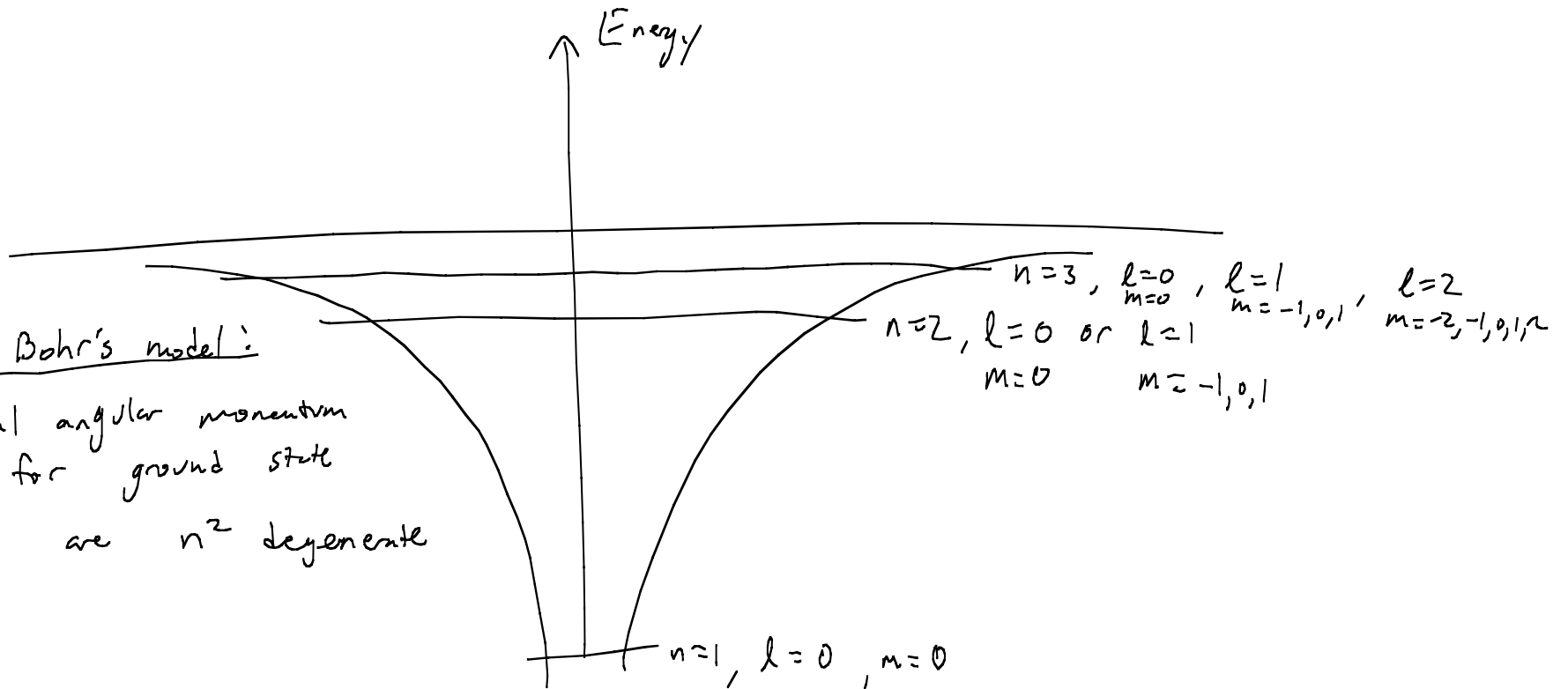
$$2n - p_0 = 0 \rightarrow 2n - \left( \frac{2me^2}{4\pi\epsilon_0\hbar^2 K} \right) = 0 \rightarrow K = \frac{me^2}{4\pi\epsilon_0\hbar^2 n}$$

So

$$E = -\frac{\hbar^2 K^2}{2m} = -\frac{\hbar^2}{2m} \left( \frac{me^2}{4\pi\epsilon_0\hbar^2 n} \right)^2 = -\frac{m}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \frac{1}{n^2}$$

$$n = 1, 2, 3, \dots$$

(Bohr's formula!)



Unlike Bohr's model:

1. Orbital angular momentum  $l=0$  for ground state
2. States are  $n^2$  degenerate

## Wave functions

$$\Psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_{lm}(\theta, \phi)$$

$$R_{nl}(r) = \frac{u_{nl}(r)}{r} = \frac{\rho^{l+1} e^{-\rho} V_{nl}(\rho)}{r},$$

$$\rho = Kr = \frac{me^2 r}{4\pi\epsilon_0 \hbar^2 n} = \frac{r}{a_0 n}$$

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{me^2} \quad \text{Bohr radius}$$

$V_{nl}(\rho)$  is determined by sum of polynomial terms given by recursion

## Splitting degeneracy

electron with orbital angular momentum

$$|\vec{L}| = \sqrt{l(l+1)} \hbar$$

has a classical magnetic moment

$$\begin{aligned} \mu &= I \cdot \text{Area} = -ef \cdot \pi r^2 = -\frac{eV}{2\pi r} \pi r^2 = -\frac{eVr}{2} = -\frac{e(mvr)}{2m} = -\frac{e\hbar}{2m} \frac{|\vec{L}|}{\hbar} \\ &= -\mu_B \frac{|\vec{L}|}{\hbar} \end{aligned}$$

↑  
"Bohr magneton"

Energy in magnetic field

$$E = -\vec{\mu} \cdot \vec{B} = -\mu_z B_z \quad \left( \vec{\mu} = \mu_x \hat{x} + \mu_y \hat{y} + \mu_z \hat{z}, \text{ and } \vec{B} = B_z \hat{z} \right)$$

So if states w/ same  $n$  have different  $\langle \mu_z \rangle$ , their energy eigenvalues will shift differently in a magnetic field and degeneracy will be broken!

## Orbital Angular momentum in z-direction

We expect  $\langle M_z \rangle = -\mu_B \frac{\langle L_z \rangle}{\hbar}$

Based on our knowledge  $\hat{p}_x = \frac{\hbar}{i} \frac{\partial}{\partial x}$ ,  $\hat{p}_y = \frac{\hbar}{i} \frac{\partial}{\partial y}$ ,  $\hat{p}_z = \frac{\hbar}{i} \frac{\partial}{\partial z}$

By analogy,  $\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$

$$\langle L_z \rangle = \langle \psi_{nlm} | \frac{\hbar}{i} \frac{\partial}{\partial \phi} | \psi_{nlm} \rangle$$

$$= \frac{\hbar}{i} \langle e^{im\phi} | \frac{\partial}{\partial \phi} | e^{im\phi} \rangle = \hbar m$$

"magnetic quantum number"

So "Zeeman energy" is

$$E = -\vec{\mu} \cdot \vec{B} = -\langle M_z \rangle B_z = \mu_B \frac{\langle L_z \rangle}{\hbar} B_z = \mu_B m B_z$$

Note:  $\mu_B = 5.8 \times 10^{-5} \frac{\text{eV}}{\text{T}}$  (small in comparison to electronic transitions)

# "Normal Zeeman Effect"

In a magnetic field,  $l \neq 0$  levels split!

Energy ↑

$n=3$   $l=0$   $m=0$

$l=1$   $m$   
1  
0  
-1

$l=2$   $m$   
2  
1  
0  
-1  
-2

$n=2$   $l=0$   $m=0$

$l=1$   $m$   
1  
0  
-1

$n=1$   $l=0$   $m=0$

What effect does this have on optical spectra?  
Not all transitions are allowed!