

Diffraction Gratings, Atomic Spectra

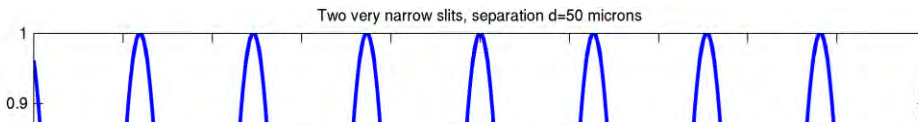
Prof. Shawhan (substituting for Prof. Hall)

November 14, 2016

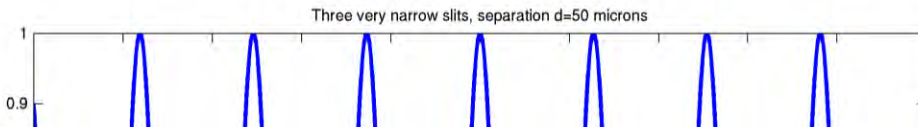
Visual Comparisons

Increase number
of slits:

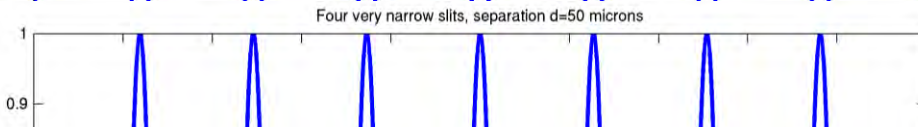
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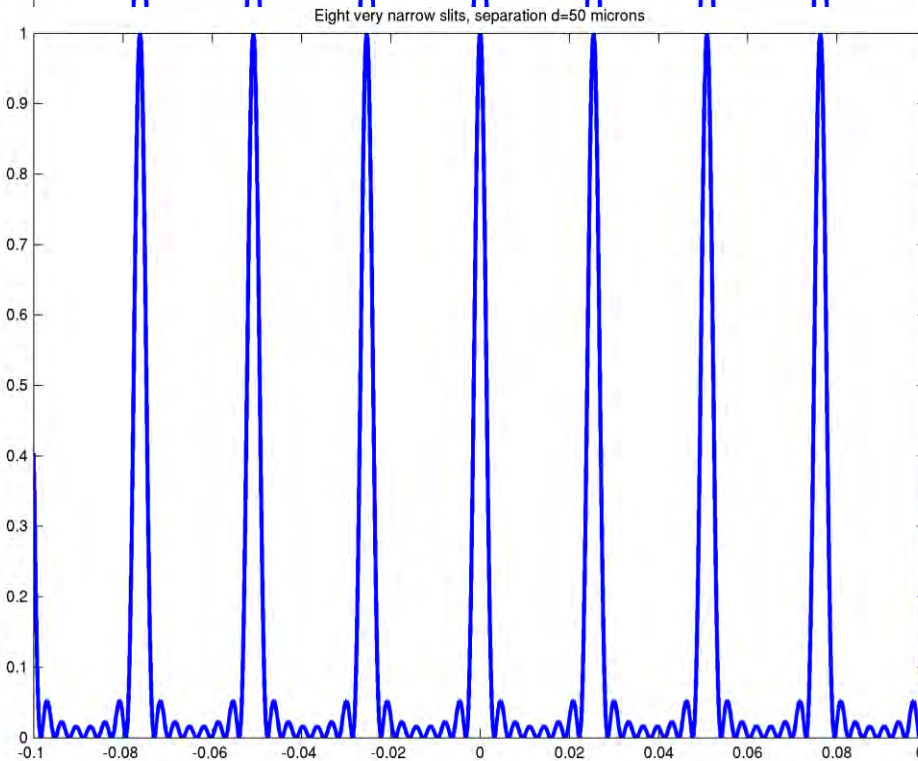
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4



8



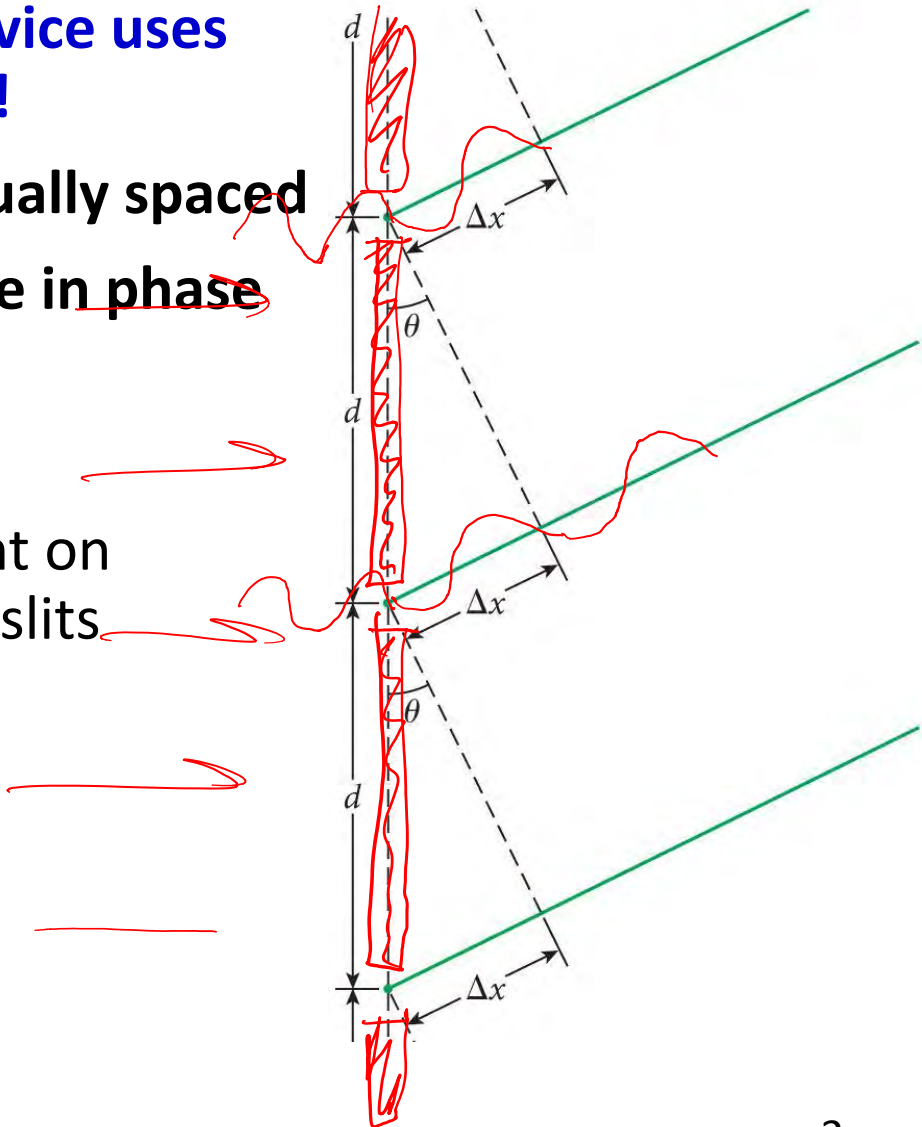
Diffraction Grating *(large N_{slit})*

Note: despite the name, this device uses interference, not diffraction!

Many slits (or obstructions), equally spaced

Need light from *all* of them to be in phase to get a bright spot

Ideal picture: plane wave incident on grating, so same phase at all slits

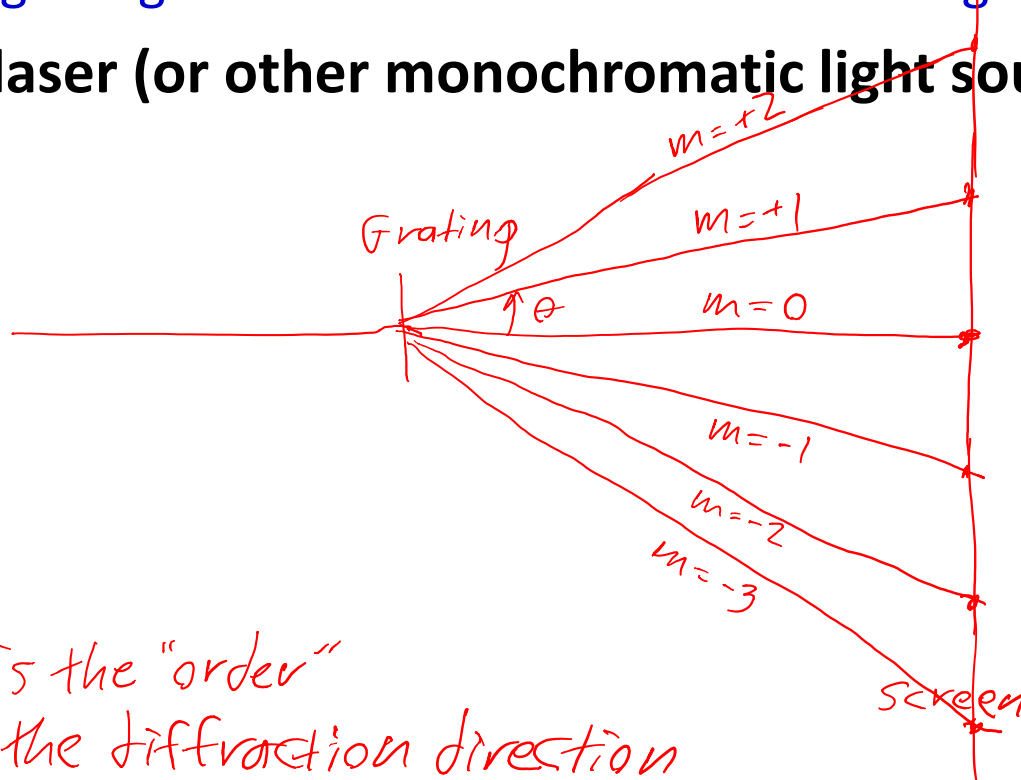


Pattern from a Diffraction Grating

Constructive interference condition: $a \sin \theta = m\lambda$ spacing between "slits" $\sin \theta = m \frac{\lambda}{a}$

Note: gratings can be made with small $a \rightarrow$ large θ 's

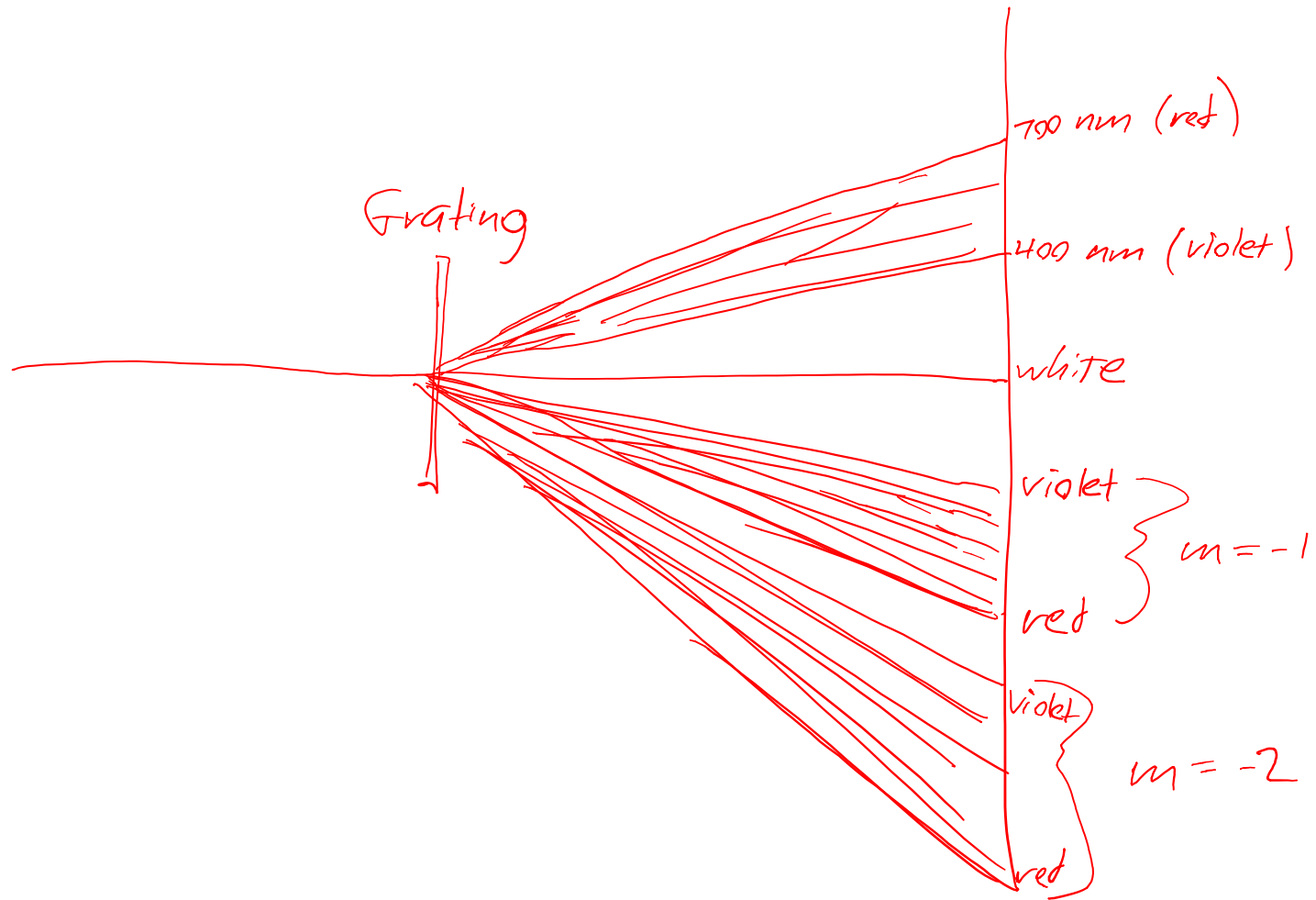
From a laser (or other monochromatic light source):



m is the "order"
of the diffraction direction

Pattern from a Diffraction Grating

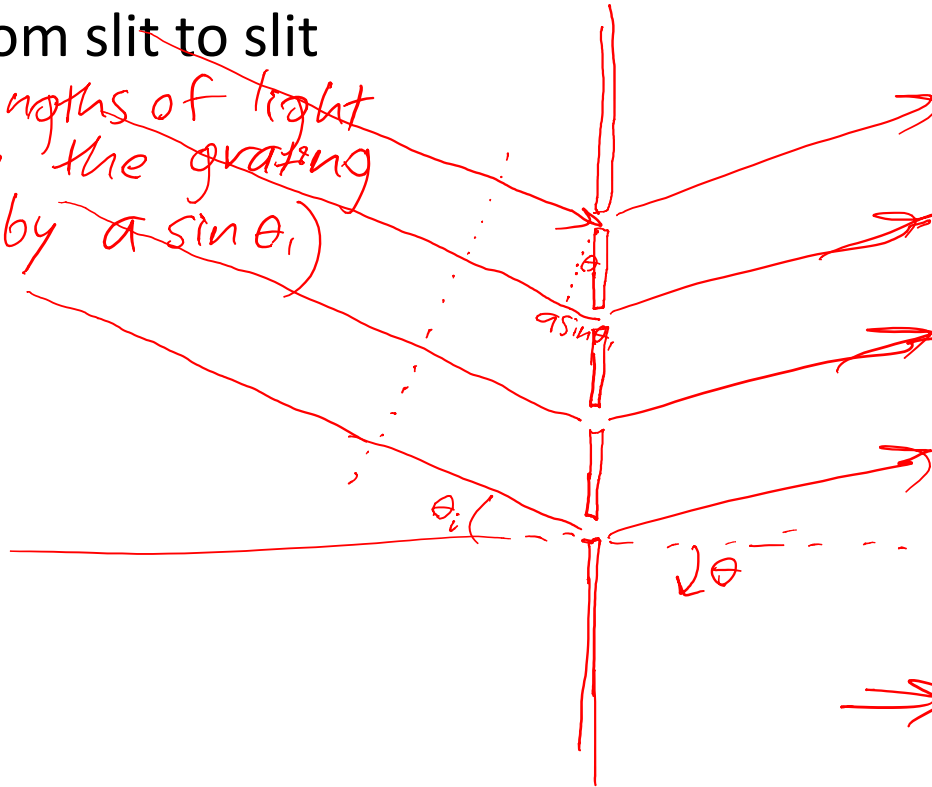
From a white-light (continuous spectrum) source:



What if light is incident on grating at an angle?

Plane wave is now coming in at an angle, so there is a phase shift from slit to slit

(path lengths of light before the grating differ by $a \sin \theta_i$)



⇒ shift in the bright spot positions, relative to $\theta_i = 0$

Still need outgoing light to all be in phase to get a bright spot, i.e. when $a(\sin \theta_i + \sin \theta) = m\lambda$

Consequences for your experiment

Can you count on the grating in your spectrometer to be perfectly aligned, normal to the light beam? *No*

How you can align it:

Retro-reflect

Adjust so that diffracted lines are at symmetric angles

How you can take data intelligently to minimize systematic error from mis-alignment:

*Measure diffraction in both directions,
both positive & negative θ 's,
and average*

Transmission vs. Reflection Gratings

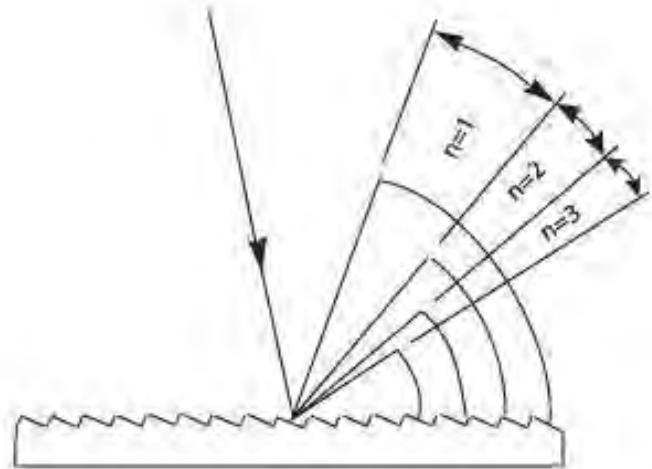
Transmission: slits, or scratches, or a fine mesh of wires

Reflection: Reflective surface with interruptions
or surface height changes

Note: angles of diffracted beams are typically not small, so you can't make the approximation $\sin \theta \approx \theta$

Tuned reflective surface:

To improve the “efficiency”
for a certain refraction order

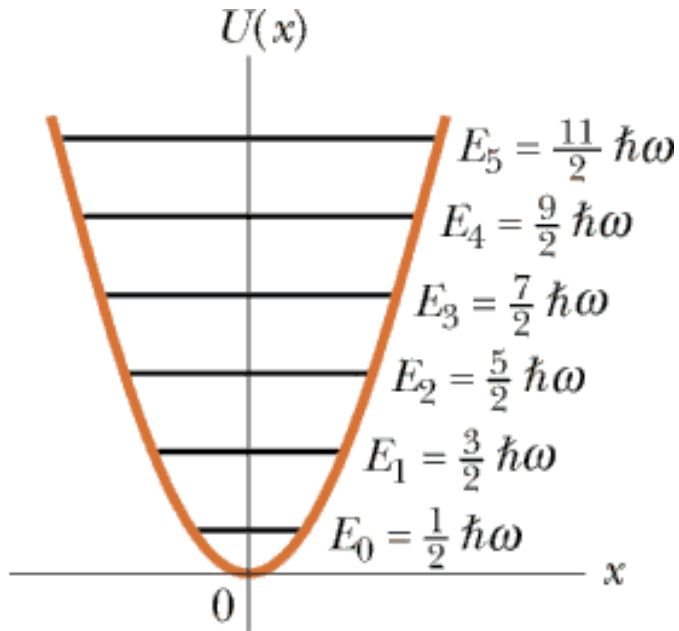


Energy Levels and Transitions

It's all about the potential!

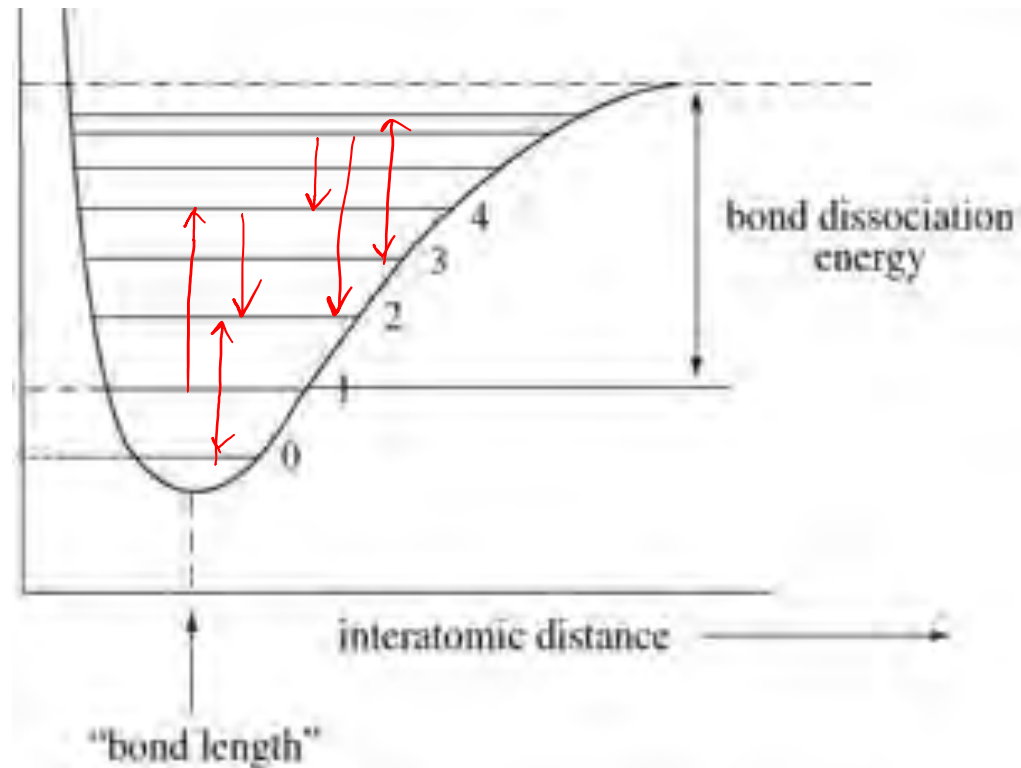
A quantum state describes a *system*, e.g. an electron in a potential

Harmonic Oscillator



©2004 Thomson - Brooks/Cole

Bond between atoms



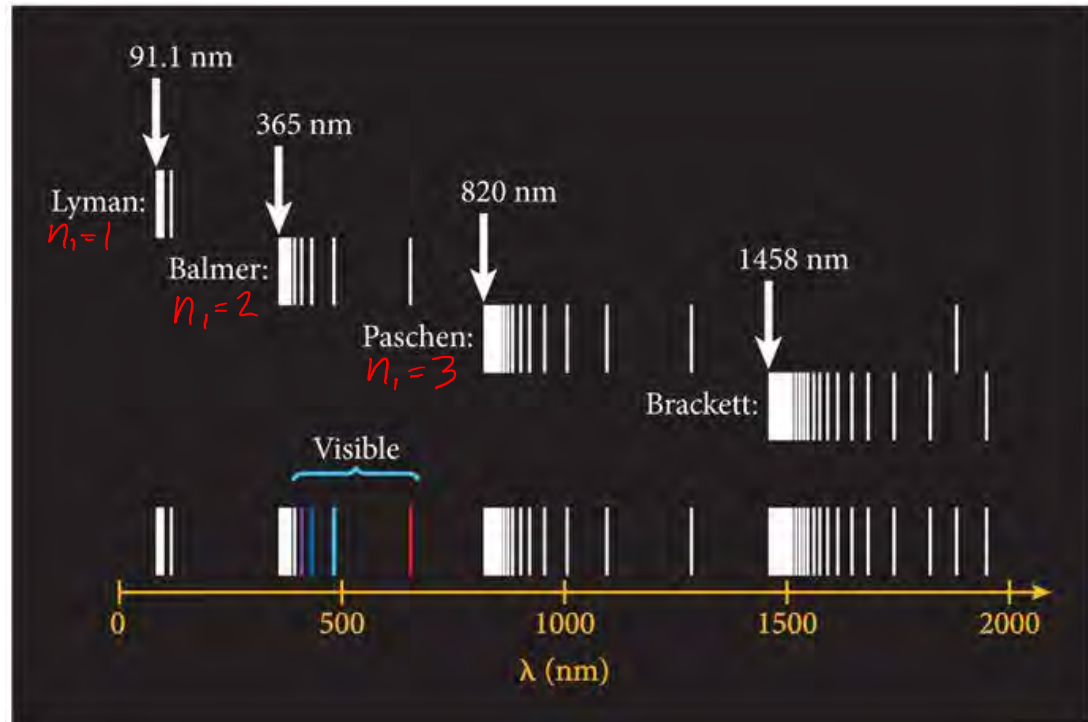
Spectrum of Hydrogen Lamp

Spectrum spread out using a diffraction grating
(Better than using dispersion in a glass prism)

Empirical formula by Balmer: $\lambda = (364.56 \text{ nm}) \frac{n^2}{(n^2 - 4)}$ *for integers 1* $= \left(\frac{\quad}{4}\right) \frac{1}{\left(\frac{1}{4} - \frac{1}{n^2}\right)}$

Full spectrum of hydrogen emission lines:

Includes UV and infrared
Must be from transitions between energy levels



Bohr Model for the Atom

Picture electrons orbiting the nucleus

Problems with that, from classical theory:

- Electron should be able to have any energy level
- Charged particle in orbit should radiate energy and collapse

Bohr's model:

Assume that electrons can only occupy discrete orbits with angular momentum equal to a multiple of \hbar

Solving the circular motion problem gives

orbit radii:

$$r_n = a_0 n^2$$

$$\text{with } a_0 = \hbar^2 / \mu k e^2 = 5.295 \times 10^{-11} \text{ m}$$

μ (reduced) mass

$$E_n = -E_0 / n^2 \quad \text{with } E_0 = k e^2 / 2 a_0 = 13.6 \text{ eV}$$

(Neglecting fine structure from electron spin-orbit coupling, and hyperfine structure from nuclear spin couplings)

Hydrogen Atom Transitions

Alternatively,

$$E_n = \frac{-R_\infty hc}{n^2}$$

R_∞ is the “Rydberg constant”, $1.09737 \times 10^7 \text{ m}^{-1}$

$R_\infty hc$ is the “Rydberg energy”, $\sim 13.6 \text{ eV}$

But for a hydrogen atom, we should use the reduced mass

→ R_H is Rydberg constant for hydrogen, $1.09678 \times 10^7 \text{ m}^{-1}$

$$\mu = \frac{m_p m_e}{m_p + m_e}$$

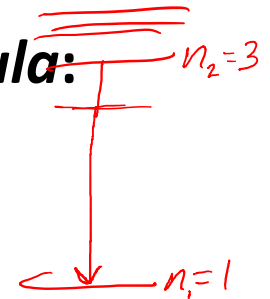
Starting from $E_n \propto -1/n^2$...

A photon emitted or absorbed in a transition must have energy equal to the difference of two energy levels

Photon wavelengths are given by the *Rydberg formula*:

$$\frac{1}{\lambda} = R_H \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

Example: $\frac{1}{\lambda} = R_H \left(\frac{1}{1^2} - \frac{1}{3^2} \right)$



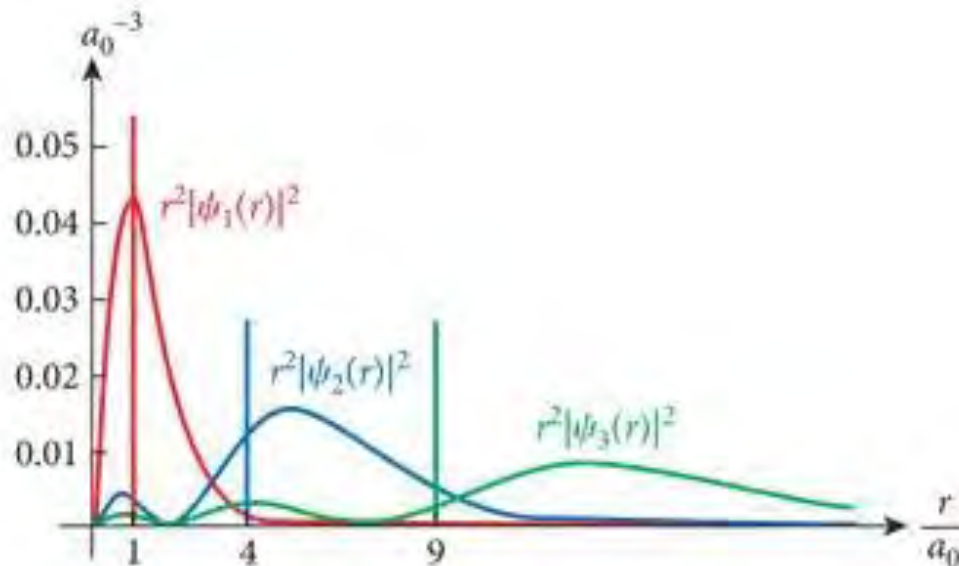
Quantum Mechanics Solution

Quantum mechanical system with one electron in Coulomb
(electrostatic) potential

3-D system

Exactly solvable, but the math is complicated

$$\frac{-\hbar^2}{2\mu} \left[\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right] + \left(\frac{-ke^2}{r} \right) \psi = E \psi$$



Extending to Other Atoms

Single-electron atoms

Simple!

Change e^2 to Ze^2 and use appropriate reduced mass μ

*↑
charge of nucleus*

Multi-electron atoms

Complicated!

Multi-particle quantum state with interacting electrons

Notes about Atomic Spectra Experiment

Manual equipment and data recording

Uses a glass diffraction grating

Figure out what the knobs do

Vernier scale for angles – do you know how to read it?

Grating needs to be aligned (might be OK already, or might not)

Suggest using Matlab scripts for data analysis calculations

Evaluate measurement uncertainties

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Strong

Intensity	Vacuum Wavelength (Å)	Spectrum	Reference
20	893.0847	Hg II	SR01
12	915.819	Hg II	SR01
20	942.630	Hg II	SR01
25	962.711	Hg II	SR01
25	969.142	Hg II	SR01
20	1039.6315	Hg II	SR01
20	1062.7802	Hg II	SR01
1000 P	1649.9373	Hg II	SR01
1000 P	1849.499	Hg I	WA63
1000 P	1942.273	Hg II	SR01
15	1973.794	Hg II	SR01
10	1987.841	Hg II	SR01

Intensity	Air Wavelength (Å)	Spectrum	Reference
20	2026.860	Hg II	SR01
400 P	2052.828	Hg II	SR01
20	2224.711	Hg II	SR01
10	2252.786	Hg II	SR01
60	2260.294	Hg II	SR01
400 P	2262.223	Hg II	SR01

60	3131.839	Hg I	BAL50
12	3208.169	Hg II	SR01
10	3532.594	Hg II	SR01
10	3605.762	Hg II	SR01
600 P	3650.153	Hg I	BAL50
70	3654.836	Hg I	BAL50
50	3663.279	Hg I	BAL50
1000 P, c	3983.931	Hg II	SR01
400 P	4046.563	Hg I	BAL50
60	4339.223	Hg I	BAL50
100	4347.494	Hg I	BAL50
1000 P	4358.328	Hg I	BAL50
12 c	5128.442	Hg II	SR01
15	5204.768	Hg II	SR01
80 P	5425.253	Hg II	SR01
500 P	5460.735	Hg I	BAL50
200 P	5677.105	Hg II	SR01
50	5769.598	Hg I	BAL50
60	5790.663	Hg I	BAL50
12	5871.279	Hg II	SR01
20 c	5888.939	Hg II	SR01
15	6146.435	Hg II	SR01
250 P, c	6149.475	Hg II	SR01
25	7081.90	Hg I	F54
6	7346.508	Hg II	SR01
250 P	7944.555	Hg II	SR01
6 h	9520.198	Hg II	SR01
200 P	10139.76	Hg I	BAL50
50	13570.21	Hg I	H53
40	13673.51	Hg I	H53
50	15295.82	Hg I	H53
50	17072.79	Hg I	H53
25	23253.07	Hg I	PBT55