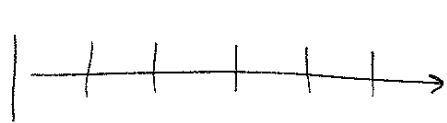
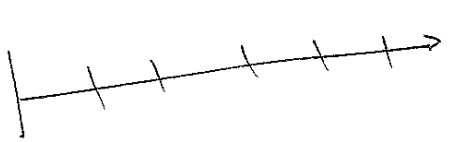


Lecture 8 - Two Beam Interference Phys 375 10/23/07



$$E_1 = A e^{i(kx_1 - \omega t)}$$

P
- waves overlap here



$$E_2 = A e^{i(kx_2 - \omega t)}$$

$$E_p = E_1 + E_2$$

$$I \propto E_p E_p^*$$

$$= 4A^2 \cos^2 \frac{k\delta}{2}, \quad \delta = x_2 - x_1 = \text{path difference.}$$

More generally, $E_p E_p^* = 4A^2 \cos^2 \frac{\Delta\phi}{2}$,

$\Delta\phi$ = phase difference between beams.

For above case, $\Delta\phi = 2\pi \left(\frac{x_2 - x_1}{\lambda} \right) = k\delta$, $k = \frac{2\pi}{\lambda}$

To observe interference we need

- same polarization
 - coherent waves
- } use the same light source to meet these conditions

Real light waves:

Not perfect cosine wave.



"coherence time" = length of time wave remains predictable

"coherence length" = coherence time \times speed of light

He-Ne lasers: $\tau_c \sim 1 \text{ ns} \sim 30 \text{ ns}$

~~l_c~~ $l_c \sim 20 \text{ cm to } 10 \text{ m}$

White light:

~~τ_c~~ $l_c \sim \frac{1}{2} \text{ wavelength} \sim 200 \text{ nm}$

$\tau_c \sim 1 \times 10^{-15} \text{ sec.}$

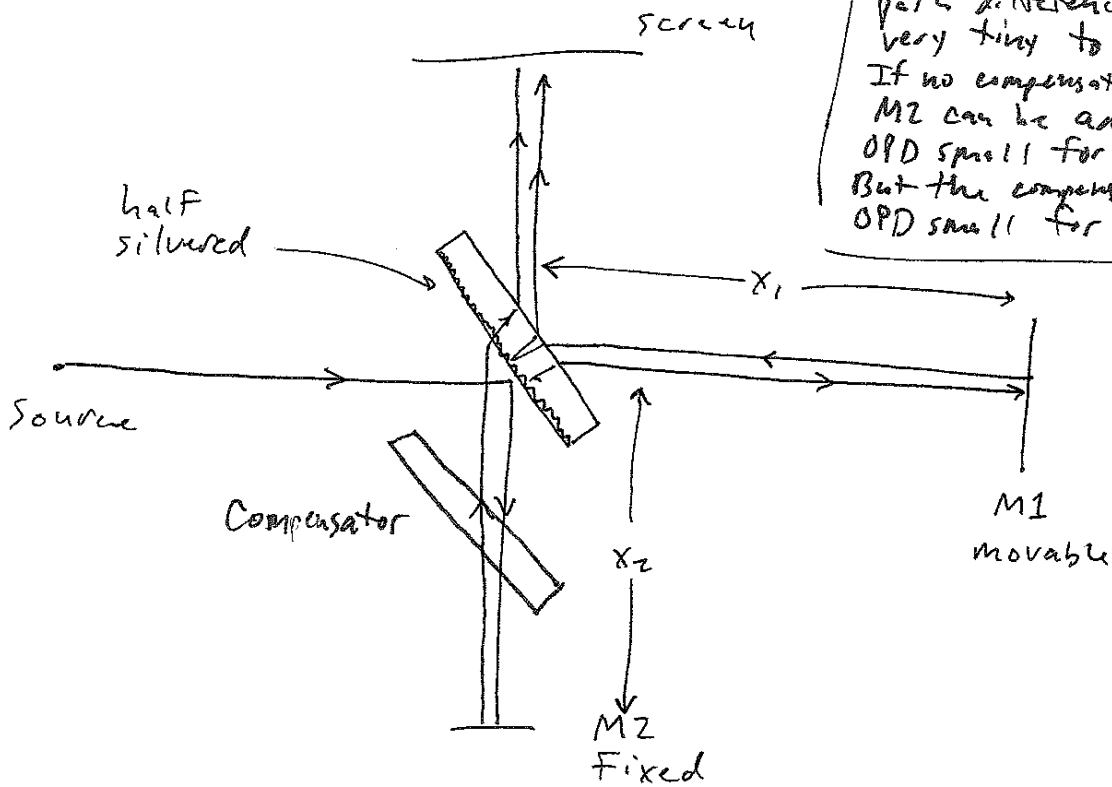


Michelson - 1852-1931

First American Nobel Prize - 1907

(2)

Michelson Interferometer



Compensator:
 For white light, the optical path difference must be very tiny to see interference. If no compensator exists, the MZ can be adjusted to make OPD small for one wavelength. But the compensator makes OPD small for all wavelengths.

$$E_p E_p^* = 4A^2 \cos^2 \frac{\Delta\phi}{2}$$

path difference = $z(x_1 - x_2) = \delta$

phase difference due to path difference = $k\delta = 2k(x_2 - x_1)$

phase difference due to ^{external} internal reflection of beam along M2 = π

total phase difference = $2k(x_2 - x_1) + \pi = \Delta\phi$

For constructive interference,

$$\frac{\Delta\phi}{2} = m\pi, \quad m = 0, \pm 1, \pm 2, \dots$$

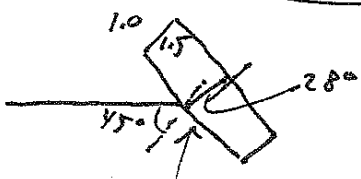
$$k(x_2 - x_1) + \frac{\pi}{2} = m\pi$$

$$\frac{2\pi}{\lambda}(x_2 - x_1) + \frac{\pi}{2} = m\pi$$

"fringe order"

$$\frac{2(x_2 - x_1)}{\lambda} = m + \frac{1}{2}$$

constructive interference



@ 45°, both s & p have a π phase shift for an external reflection

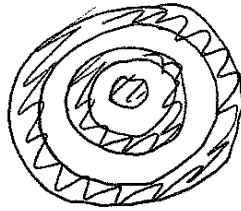
@ 28°, both s & p have no phase shift for an internal reflection

$$\frac{z(x_2 - x_1)}{\lambda} = m$$

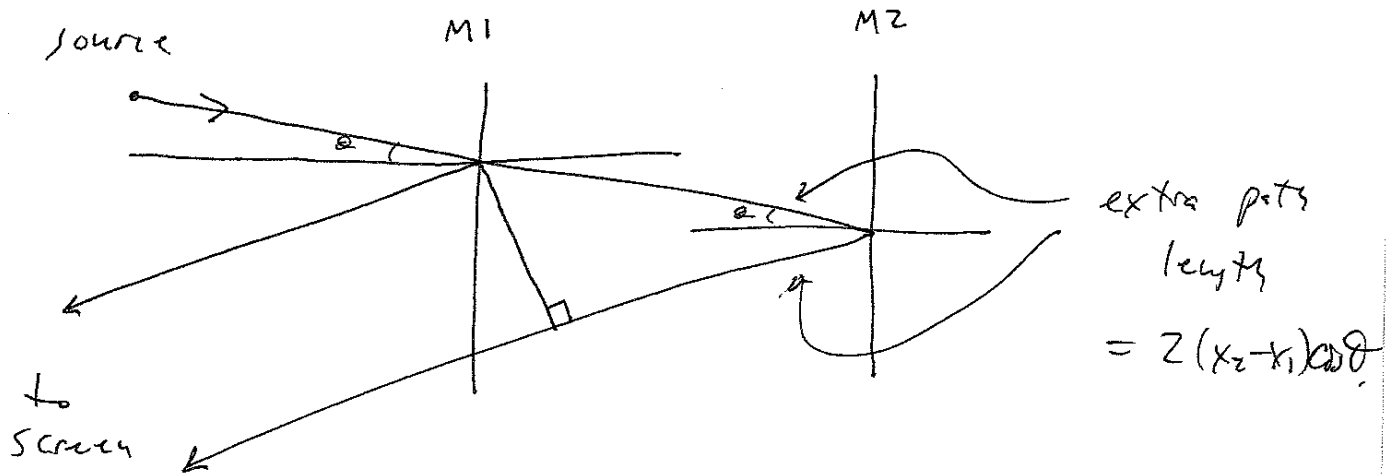
destructive interference.

(3)

In practice you don't see one spot, you see:



Because off-axis points have an extra path difference



If $x_2 \equiv x_1$, then only one fringe is observed for all θ (no off axis phase difference)

Otherwise many rings will be seen.

Usually, we don't care about which order of fringe we are observing. Instead, move M2, and count the number of fringes that pass by:

$$\Delta m = \frac{z(x_2 - x_1)}{\lambda}$$

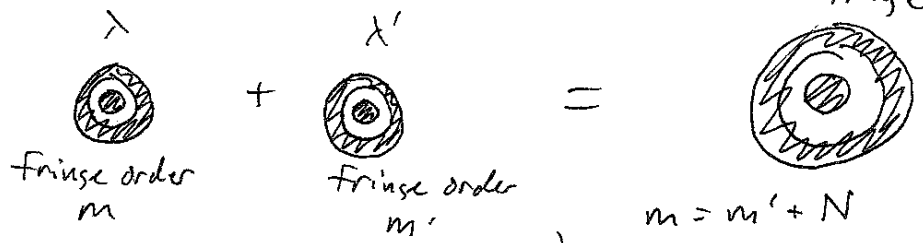
Measure $x_2 - x_1$,
 Count Δm , and
 calculate the
 wavelength λ .

Sodium Lamp

Two closely spaced wavelengths: λ, λ'
Each makes its own fringe pattern.

Suppose both patterns have a maximum:

$m = 1000$
 $m' = 1500$



Then
$$\frac{zd_1}{\lambda} = \frac{zd_1}{\lambda'} + N \quad d_1 = x_2 - x_1$$

Now move the mirror to increase d . After a while, the two patterns look like

$m = 1300$
 $m' = 1800.5$



Then move the mirror more, and the fringes re-appear:

Now $x_2 - x_1 = d_2$

$$\frac{zd_2}{\lambda} = \frac{zd_2}{\lambda'} + \underbrace{N+1}_{\substack{m \text{ and } m' \text{ differ} \\ \text{by } N+1 \text{ now}}}$$

$$\frac{z(d_2 - d_1)}{\lambda} = \frac{z(d_2 - d_1)}{\lambda} + 1$$

(5)

$$\frac{z \Delta d}{\lambda} - \frac{z \Delta d}{\lambda'} = 1$$

$$\lambda - \lambda' = \frac{\lambda \lambda'}{z \Delta d}$$

or, since $\lambda \approx \lambda'$,

$$\lambda - \lambda' = \boxed{\Delta \lambda \approx \frac{\lambda^2}{z \Delta d}}$$

Δd is the change in the mirror distance between the two "high fringe visibility" regions.