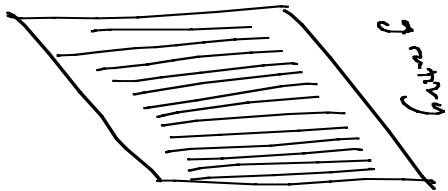
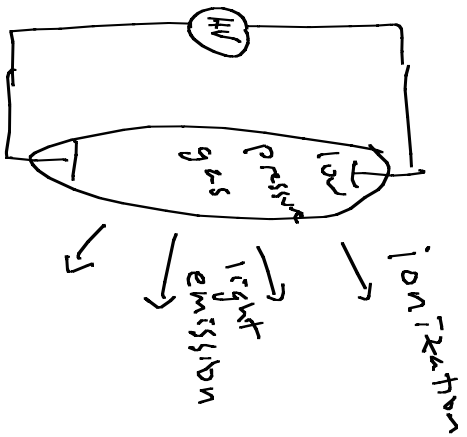
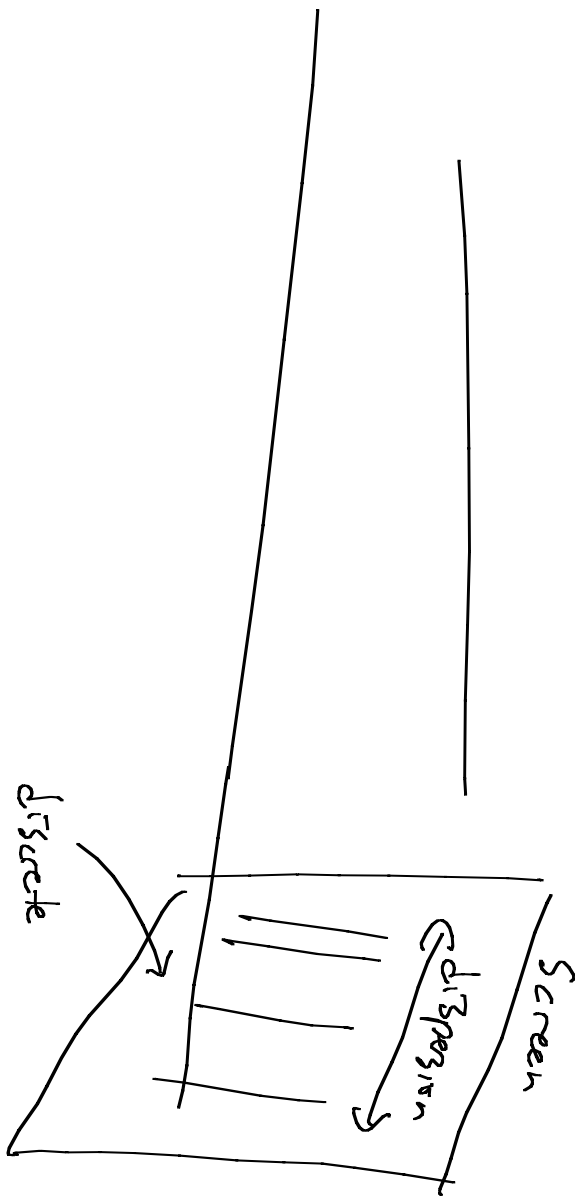


# Atomic discharge



# Spectroscopy



# Bohr model

"quantized" angular momentum

$$(1) \quad mvr = n\hbar$$

Force balance

$$(2) \quad \frac{mv^2}{r} = \frac{e^2}{r^2}$$

[centrifugal = Coulomb]

Subst. (\*) into (1),

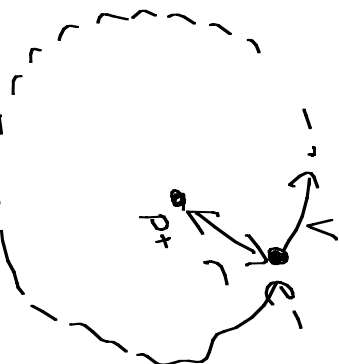
$$\sqrt{e^2 m r} = n\hbar \Rightarrow r = \frac{n^2 \hbar^2}{e^2 m}$$

$$\text{Total Energy} = \text{Kinetic} + \text{potential} = \frac{1}{2}mv^2 - \frac{e^2}{r}$$

$$\text{But, (*) gives Energy} = \frac{1}{2} m \frac{e^2}{m r} - \frac{e^2}{r} = -\frac{e^2}{2r}$$

$n = 1, 2, \dots$   
"principle quantum number"

$$\Rightarrow v = \sqrt{\frac{e^2}{m r}} \quad (*)$$



So,  $E_{\text{energy}} = -\frac{e^2}{2} \frac{e^2 m}{h^2 k^2} = \left(\frac{e^2}{hc}\right)^2 \frac{m c^2}{2 n^2}$

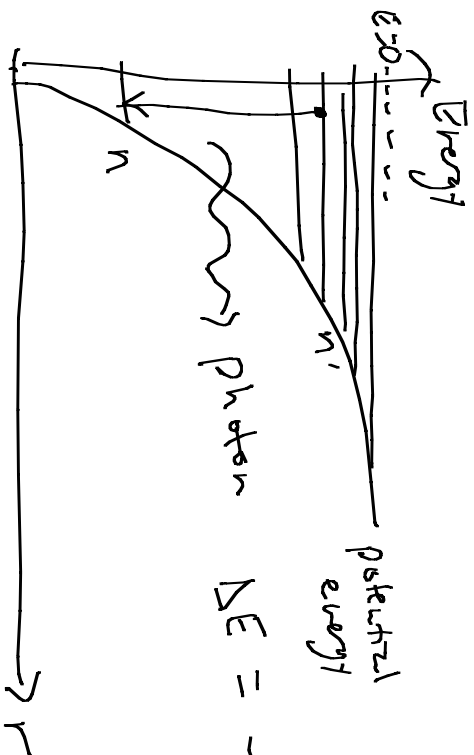
$\alpha \equiv \frac{e^2}{hc}$  "fine structure constant"  $\sim \frac{1}{137}$

$m c^2 \equiv E_0$  "rest mass" 511 keV

$E_{\text{energy}} = -\alpha^2 \frac{E_0}{2 n^2}$

"Radiative transitions"

---



$\Delta E = -\frac{R}{n'^2} - \left(-\frac{R}{n^2}\right) = R \left(\frac{1}{n^2} - \frac{1}{n'^2}\right)$

Rydberg (1888)  
(13.6 eV)

# Fine Structure: electron magnetic moment (spin)

relativity!



$\ominus_e$

$+Ze$

nucleus at rest

$+Ze \rightarrow -v$

electron at rest

Magnetic field at  
electron due to charge  
current from moving nucleus

$$\vec{B} = \frac{\mu_0}{4\pi} \frac{Ze \left( -\vec{v} \times \vec{r}^3 \right)}{r^3} = -\frac{1}{4\pi\epsilon_0 c^2} \frac{Ze \vec{L}}{m r^3}$$

$$\left( \vec{L} = m \vec{v} \times \vec{r} \right)$$

$$E_{\text{energy}} = -\vec{\mu} \cdot \vec{B} = - \left( g \frac{\mu_B}{\hbar} \vec{S} \right) \cdot \left( \frac{-1}{4\pi\epsilon_0 c^2} \frac{Ze \vec{L}}{m r^3} \right) = \frac{1}{4\pi\epsilon_0} \frac{Ze^2}{m r^2} \vec{S} \cdot \vec{L}$$

"Spin-orbit"

Since  $\vec{S}$  and  $\vec{L}$  are in (near-)units of  $\hbar$ , we have

$$\langle E_{\text{spin-orbit}} \rangle \sim \frac{e^2}{m^2 c^2} \frac{\hbar \cdot \hbar}{r^3} \quad \text{and} \quad \langle r^{-3} \rangle \sim \left( \frac{\hbar^2}{e^2 m} \right)^3 \quad \text{so}$$

$$\langle E_{\text{spin-orbit}} \rangle \sim \frac{e^8 m c^2}{\hbar^4 c^4} = \left( \frac{e^2}{\hbar c} \right)^4 m c^2 \rightarrow \text{factor of } \alpha^2 \text{ smaller than electron state splitting!}$$

for fixed  $n$ , there are  $2n-1$  possible values of  $L_z$  ( $L_z = \hbar m_l$ )  $\rightarrow$  spin-orbit splitting of multiple states w/ same  $n$ .