

HW7

1.

$$(u_n \cdot u_n) = A^2 \int_{-L}^L \cos(n\pi x/L)^2 dx = A^2 \int_{-L}^L \frac{1+\cos(2n\pi x/L)}{2} dx = A^2 L = 1.$$

So $A = 1/\sqrt{L}$.

$$\begin{aligned} (u_n \cdot u_m) &= \frac{1}{L} \int_{-L}^L \cos(n\pi x/L) \cos(m\pi x/L) dx \\ &= \frac{1}{L} \int_{-L}^L \frac{1}{2} \left(\cos\left(\frac{(n+m)\pi x}{L}\right) + \cos\left(\frac{(n-m)\pi x}{L}\right) \right) dx = 0 \text{ for } n \neq m. \end{aligned}$$

Therefore $(u_n \cdot u_m) = \delta_{nm}$.

2.

$f(x) = -x^2 + 1$ is an even function of x so only Fourier cosine terms survive.

$$f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos(n\pi x) \text{ with } a_n = \int_{-1}^1 f(x) \cos(n\pi x) dx = -\frac{4(-1)^n}{n^2\pi^2}. \text{ (integrate by part twice)}$$

$$\begin{aligned} 3. \psi(\vec{k}) &= \frac{1}{(2\pi)^3} \int \lambda e^{-r/a} e^{-i\vec{k} \cdot \vec{r}} d^3r = \frac{\lambda}{(2\pi)^3} \int dr d\theta d\phi r^2 \sin\theta e^{-r/a} e^{-ikr \cos\theta} \\ &= \frac{\lambda 2\pi}{(2\pi)^3} \frac{1}{ik} \int dr r e^{-r/a} (e^{ikr} - e^{-ikr}) \\ &= \frac{\lambda 2\pi}{(2\pi)^3} \frac{1}{ik} \left(\frac{4ia^3 k}{(1+a^2 k^2)^2} \right) = \frac{\lambda a^3}{\pi^2(1+a^2 k^2)^2}. \text{ (integrate by part)} \end{aligned}$$

4.

If $k > 0$, then close the contour in the upper half plane. (the pole is ia)

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{x^2+a^2} e^{ikx} dx = \frac{1}{2\pi} 2\pi i \left(\frac{e^{ik(ia)}}{2ia} \right) = \frac{1}{2a} e^{-ka}$$

If $k < 0$, then close the contour in the lower half plane. (the pole is $-ia$)

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{x^2+a^2} e^{ikx} dx = \frac{1}{2\pi} (-2\pi i) \left(\frac{e^{ik(-ia)}}{-2ia} \right) = \frac{1}{2a} e^{ka}.$$

$$\text{Therefore } \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{x^2+a^2} e^{ikx} dx = \frac{1}{2a} e^{-|k|a}.$$

5.

$$(1) e^z = e^{x+iy} = e^x (\cos y + i \sin y) = (e^x \cos y) + i(e^x \sin y) = g + ih$$

It satisfies Cauchy-Riemann condition at all points in the complex plane.

$$\partial_x g = e^x \cos y = \partial_y h, \partial_y g = -e^x \sin y = -\partial_x h. \text{ So } e^z \text{ is analytic at } z = 0.$$

$$(2) \text{ Let } f(z) = \sqrt{z}. f'(z_0) = \lim_{\Delta z \rightarrow 0} \frac{f(z_0 + \Delta z) - f(z_0)}{\Delta z}.$$

$$f'(0) = \lim_{\Delta z \rightarrow 0} \frac{f(\Delta z) - f(0)}{\Delta z} = \lim_{\Delta z \rightarrow 0} \frac{1}{\sqrt{\Delta z}}.$$

If we choose $\Delta z = \Delta x$, $f'(0) = \lim_{\Delta x \rightarrow 0} \frac{1}{\sqrt{\Delta x}} = \infty$. (approach the origin from real axis)

If we choose $\Delta z = i\Delta y$, $f'(0) = \lim_{\Delta y \rightarrow 0} \frac{1}{\sqrt{i\Delta y}} = \frac{1}{\sqrt{i}} \infty$. (approach the origin from the imaginary axis)

The two limits are different. Therefore $\sqrt{z}$ is not analytical at $z = 0$.

6.

$$F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) e^{i\omega t} dt, H(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} h(t) e^{i\omega t} dt.$$

$$\int_{-\infty}^{\infty} F(\omega) H^*(\omega) e^{-i\omega t} d\omega = \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} dt_1 \int_{-\infty}^{\infty} dt_2 f(t_1) h^*(t_2) \int_{-\infty}^{\infty} d\omega e^{-i\omega(t-t_1+t_2)}$$

$$= \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} dt_1 \int_{-\infty}^{\infty} dt_2 f(t_1) h^*(t_2) 2\pi \delta(t-t_1+t_2)$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau f(t_1) h^*(t+\tau)$$

$$\begin{aligned}
7. \quad \int_0^\infty \left(\frac{\sin x}{x}\right)^2 dx &= \frac{1}{2} \int_{-\infty}^\infty \left(-\frac{1}{4}\right) \frac{e^{i2x}-2+e^{-i2x}}{x^2} dx = -\frac{1}{8} \left(\int_{-\infty}^\infty \frac{e^{i2x}-1}{x^2} dx + \int_{-\infty}^\infty \frac{e^{-i2x}-1}{x^2} dx \right) \\
&= -\frac{1}{8} 2\pi i \left(\frac{1}{2}\right) \operatorname{res} \left[\frac{e^{i2x}-1}{x^2} \right] - \frac{1}{8} 2\pi i \left(-\frac{1}{2}\right) \operatorname{res} \left[\frac{e^{-i2x}-1}{x^2} \right] \\
&= -\frac{1}{8} \pi i (2i) - \frac{1}{8} \pi i (2i) = \frac{\pi}{2}
\end{aligned}$$

Note when calculating $\int_{-\infty}^\infty \frac{e^{i2x}-1}{x^2} dx$ you will close the contour in the upper half plane ($\lim_{x \rightarrow i\infty} e^{i2x} = 0$). When calculating $\int_{-\infty}^\infty \frac{e^{-i2x}-1}{x^2} dx$ you will close the contour in the lower half plane ($\lim_{x \rightarrow -i\infty} e^{-i2x} = 0$). The $(\pm \frac{1}{2})$ factor is because you detour around the pole at $z = 0$ with counterclockwise/clockwise half circle.

8.

$$\phi(x) = \frac{1}{2\pi} \int_{-\infty}^\infty \phi(k) e^{ikx} dk, \delta(x) = \frac{1}{2\pi} \int_{-\infty}^\infty e^{ikx} dk.$$

Inserting these to the diffusion equation, you will get $(Dk^2 + K^2) \phi(k) = Q$.

So $\phi(k) = \frac{Q}{(Dk^2 + K^2)}$.

$$\phi(x) = \frac{1}{2\pi} \int_{-\infty}^\infty \left(\frac{Q}{(Dk^2 + K^2)} \right) e^{ikx} dk = \frac{Q}{D} \frac{1}{2\pi} \int_{-\infty}^\infty \frac{1}{k^2 + (K/\sqrt{D})^2} e^{ikx} dk = \frac{Q}{D} \frac{1}{2K/\sqrt{D}} e^{-|x|K/\sqrt{D}}.$$

(From Prob.4)