

4. Let η_i be the displacement of block i from equilibrium. The potential energy of the system is:

$$V = \frac{1}{2} k (\eta_1 - \eta_2)^2 + \frac{1}{2} k (\eta_2 - \eta_3)^2$$

$$T = \frac{1}{2} m (\dot{\eta}_1^2 + \dot{\eta}_2^2 + \dot{\eta}_3^2)$$

\therefore Lagrangian $L = T - V$.

$$\therefore L = \sum_{i=1}^3 \frac{1}{2} m \dot{\eta}_i^2 - \sum_{i=1}^2 (T_{ij} \eta_i \eta_j - V_{ij} \eta_i \eta_j)$$

where $\underline{T} = \begin{pmatrix} m & 0 & 0 \\ 0 & m & 0 \\ 0 & 0 & m \end{pmatrix}$ $\underline{V} = \begin{pmatrix} k & -k & 0 \\ -k & 2k & -k \\ 0 & -k & k \end{pmatrix}$

Using Lagrange's equation, we find the equation of motion:

$$\underline{T} \ddot{\underline{\eta}} + \underline{V} \underline{\eta} = 0$$

here $\underline{\eta} = (\eta_1, \eta_2, \eta_3)^T$.

let $\underline{\eta}(t) = \underline{a}_j e^{i\omega_j t}$

where a_j is time-independent.

$$\Rightarrow (\underline{V} - \omega_j^2 \underline{T}) \underline{a}_j = 0$$

In order for a nontrivial solution to exist, we must have

$$\det [\underline{V} - \omega_j^2 \underline{T}] = 0$$

This leads to a cubic equation of $\omega^2 \Rightarrow \omega_1^2 = 0, \omega_2^2 = k/m,$

and $\omega_3^2 = 3k/m$. Substituting these frequencies into equation $*$, allows us to solve for the three normal modes, for which

we choose the normalization prescription

$$\underline{a}_i^T \underline{T} \underline{a}_i = 1$$

or choose $\underline{a}_i^T \underline{T} \underline{a}_j = \delta_{ij}$.