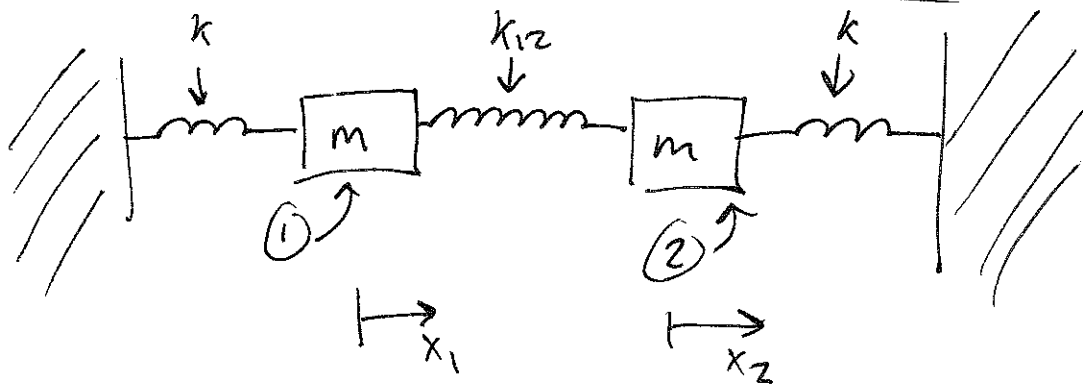


Two Coupled Mechanical Oscillators



3 springs, but just 2 spring constants.

$x_1$ : displacement of ① from equilibrium

$x_2$ : " " ② " "

Newton's 2nd Law  
↓

Force on ①:  $F_1 = -kx_1 - k_{12}(x_1 - x_2) = m\ddot{x}_1$

Force on ②:  $F_2 = -kx_2 + k_{12}(x_1 - x_2) = m\ddot{x}_2$

Equations of Motion:

$$\begin{cases} m\ddot{x}_1 + (k+k_{12})x_1 - k_{12}x_2 = 0 \\ m\ddot{x}_2 + (k+k_{12})x_2 - k_{12}x_1 = 0 \end{cases}$$

Coupled  
Differential  
Equations.

$x_1$  and  $x_2$  appear in both equations.

We must solve for  $x_1(t)$  &  $x_2(t)$  ~~at the~~ simultaneously.

~~Let's~~ Our strategy: Let's look for solutions where both masses execute harmonic motion at the same frequency.

Guessed Solution:

$$x_1 = A_1 e^{i(\omega t + \delta_1)}$$

$$x_2 = A_2 e^{i(\omega t + \delta_2)}$$

$\omega$  is unknown. But the

key feature of our guess is that both

$x_1$  AND  $x_2$  oscillate

at the same frequency.

Normal Mode

\* If all the masses in a system oscillate at the same frequency, then we call the motion a "Normal Mode".

Normal Frequency: The frequency of a normal mode.

Let's Simplify Notation: re-write the guessed solution:

$$x_1 = A_1 e^{i(\omega t + \delta_1)} = \underbrace{(A_1 e^{i\delta_1})}_{\equiv B_1} e^{i\omega t} \equiv B_1 e^{i\omega t}$$

Also:  $x_2 = B_2 e^{i\omega t}$   
 $\uparrow$   
 complex,  
 2 free parameters

$B_1$  is complex,  
 it has two  
free parameters.  
 (real & imaginary parts).

Week 6

Phys 273

③

Substitute the guessed solution into the equation of motion:

$$\left. \begin{aligned} m(-\omega^2 B_1) + (k+k_{12})B_1 - k_{12}B_2 &= 0 \\ m(-\omega^2 B_2) + (k+k_{12})B_2 - k_{12}B_1 &= 0 \end{aligned} \right\} e^{i\omega t} \text{ has cancelled everywhere.}$$

Gather Terms:

$$\left. \begin{aligned} (k+k_{12}-m\omega^2)B_1 - k_{12}B_2 &= 0 \quad \text{Eq. 1} \\ -k_{12}B_1 + (k+k_{12}-m\omega^2)B_2 &= 0 \quad \text{Eq. 2} \end{aligned} \right\}$$

To have a solution, the determinant must equal zero:

$$(k+k_{12}-m\omega^2)^2 - k_{12}^2 = 0$$

$$k+k_{12}-m\omega^2 = \begin{matrix} + \\ - \end{matrix} k_{12} \quad \text{2 possibilities.}$$

$$\omega = \sqrt{\frac{k+k_{12} \pm k_{12}}{m}}$$

We have found 2 normal mode frequencies:

Let's call them:

$$\left. \begin{aligned} \omega_1 &= \text{smaller frequency} = \sqrt{\frac{k}{m}} \\ \omega_2 &= \text{larger frequency} = \sqrt{\frac{k+2k_{12}}{m}} \end{aligned} \right\}$$

So the small frequency solution is:

$$x_1 = B_1 e^{i\omega_s t}$$

$$x_2 = B_2 e^{i\omega_s t} \quad , \quad \omega_s = \sqrt{\frac{k}{m}}$$

But we are not done. We can show that for this solution, we must have  $B_1 = B_2$ . To see this, substitute  $\omega_s = \sqrt{k/m}$  into (Eq. 1) & (Eq. 2):

$$\begin{cases} (k + k_{12} - k) B_1 - k_{12} B_2 = 0 \\ -k_{12} B_1 + (k + k_{12} - k) B_2 = 0 \end{cases}$$

or

$$\begin{cases} k_{12} (B_1 - B_2) = 0 \\ -k_{12} (B_1 - B_2) = 0 \end{cases}$$

or

$$\boxed{B_1 = B_2} \text{ for the small frequency solution.}$$

Let's call it  $B_1 = B_2 \equiv \underline{B_s}$  = "small frequency case."

Then our solution is

$$\begin{cases} x_1 = B_s e^{i\omega_s t} \\ x_2 = B_s e^{i\omega_s t} \end{cases}$$

$B_s = \text{complex} = 2 \text{ free parameters,}$   
(real & imaginary parts)

This is called the "symmetric mode", because both oscillators have exactly the same motion.  $\Rightarrow$  Amplitude, phase, and frequency are identical.

Large Frequency mode: Exactly the same methods

leads to

$$B_1 = -B_2 \quad \text{for } \omega_L = \sqrt{\frac{k+2k_{12}}{m}}$$

( $-$ ) sign means

that  $x_1$  &  $x_2$  are out-of-phase by  $180^\circ$ .

Call  $B_1 = B_L$ . Then the large frequency solution is

$$\begin{array}{l} x_1 = B_L e^{i\omega_L t} \\ x_2 = -B_L e^{i\omega_L t} \end{array}$$

}  $B_L$  has 2 free parameters: real & imaginary parts.

We call this the "anti-symmetric mode" because the two oscillators are  $180^\circ$  out-of-phase with each other.

## General Solution

The 2 normal mode solutions are the simplest type of motion that the system may execute. But we can find a general solution by adding the normal mode solutions. This works because the equations of motion are linear.

$$\begin{aligned} x_1 &= B_S e^{i\omega_S t} + B_L e^{i\omega_L t} \\ x_2 &= B_S e^{i\omega_S t} - B_L e^{i\omega_L t} \end{aligned}$$

The most general solution.

Note that we have 4 free parameters: the real & imaginary parts of  $B_S$  &  $B_L$ . We need 4 initial conditions to specify them:

position and velocity of ① at  $t=0$   
 & position and velocity of ② at  $t=0$ .

Let's take the real part and apply one particular set of initial conditions:

$$x_1 = b_S \cos(\omega_S t + \delta_S) + b_L \cos(\omega_L t + \delta_L)$$

$$x_2 = b_S \cos(\omega_S t + \delta_S) - b_L \cos(\omega_L t + \delta_L)$$

$b_S, \delta_S, b_L, \delta_L$  are free parameters.

Suppose that

$$x_1(t=0) = a$$

$$\dot{x}_1(t=0) = 0$$

$$x_2(t=0) = 0$$

$$\dot{x}_2(t=0) = 0$$

Then the  $\dot{x}_1$  &  $\dot{x}_2$  requirements are:

$$\begin{cases} \dot{x}_1 = -\omega_S b_S \sin(\delta_S) - \omega_L b_L \sin(\delta_L) = 0 \\ \dot{x}_2 = -\omega_S b_S \sin(\delta_S) + \omega_L b_L \sin(\delta_L) = 0 \end{cases}$$

Add these equations:  $-2\omega_S b_S \sin(\delta_S) = 0 \Rightarrow \delta_S = 0$

Subtract these equations:  $-2\omega_L b_L \sin(\delta_L) = 0 \Rightarrow \delta_L = 0$

And the  $x_1 = a$  and  $x_2 = 0$  requirements are

$$x_1 = b_S \cos(\delta_S) + b_L \cos(\delta_L) = a$$

$$x_2 = b_S \cos(\delta_S) - b_L \cos(\delta_L) = 0$$

or

$$\begin{cases} b_S + b_L = a \\ b_S - b_L = 0 \end{cases} \Rightarrow \begin{cases} b_S = \frac{a}{2} \\ b_L = \frac{a}{2} \end{cases}$$

So the complete solution for these initial conditions is

$$x_1 = \frac{a}{2} \cos(\omega_S t) + \frac{a}{2} \cos(\omega_L t)$$
$$x_2 = \frac{a}{2} \cos(\omega_S t) - \frac{a}{2} \cos(\omega_L t)$$

What does it look like?

