

Solutions to hw 12

- 1) According to the Bohr model for hydrogen, the wavelength emitted in a transition between levels n and m is given by

$$\lambda_{nm}^{-1} = R_H \left(\frac{1}{n^2} - \frac{1}{m^2} \right) \quad R_H = 109.737 \text{ nm}^{-1}$$

$$E_{nm} = \frac{hc}{\lambda_{nm}} = E_0 \left(\frac{1}{n^2} - \frac{1}{m^2} \right) \quad E_0 = 13.6 \text{ eV}$$

The shortest wavelength and highest energy for a series with fixed n is then obtained with $m \rightarrow \infty$, such that

$$\min \lambda_n = \frac{n^2}{R_H}, \quad \max E_n = \frac{E_0}{n^2}$$

	n	$\min \lambda_n$	$\max E_n$
Lyman	1	91.1 nm	13.6 eV
Balmer	2	365	3.40
Paschen	3	820	1.51
Brackett	4	1458	0.85

- 2) $E_{nm} = E_0 \left(\frac{1}{n^2} - \frac{1}{m^2} \right)$ emitted for $m \rightarrow n$ or absorbed for $n \rightarrow m$
- A) $(m, n) = (3, 5) \Rightarrow E_{nm} = -0.071 E_0$ gains the most energy
- B) $(m, n) = (5, 3) \Rightarrow E_{nm} = +0.071 E_0$ emits shortest wavelength
- C) $(m, n) = (8, 4) \Rightarrow E_{nm} = +0.047 E_0$
- d) $(m, n) = (4, 8) \Rightarrow E_{nm} = -0.047 E_0$
- loses energy for B, C ; gains for A, D

- 3) The kinetic energy for each atom is $K = \frac{hc}{\lambda} = 10.2 \text{ eV}$ for $\lambda = 121.6 \text{ nm}$.
 Because $K \ll mc^2$, we can use the nonrelativistic kinetic energy to find $\beta = (2K/mc^2)^{1/2} = 1.47 \times 10^{-4}$ where $mc^2 = 938 \text{ MeV}$ for hydrogen. $\therefore v = \beta c = 4.42 \times 10^4 \text{ m/s}$.

4) The Bohr radius for a single-electron atom or ion is given by $a = a_0/Z$ where $a_0 = 0.0529 \text{ nm}$ and where Z is the nuclear charge in units of e .

a) $\text{He}^+ \Rightarrow Z=2 \Rightarrow a = 0.0265 \text{ nm}$

b) $\text{Li}^{2+} \Rightarrow Z=3 \Rightarrow a = 0.0176 \text{ nm}$

c) $\text{Be}^{3+} \Rightarrow Z=4 \Rightarrow a = 0.0132 \text{ nm}$

5) If an electron is confined to diameter D with $\lambda \leq D$, it must have momentum $p = h/\lambda \geq h/D$. For this problem D is so small and p so large that we should use relativistic kinematics with $\gamma \gg 1 \Rightarrow K \approx pc \geq hc/D$. Thus, $D \sim 10^{-5} \text{ nm} \Rightarrow K \geq 1.2 \times 10^8 \text{ eV}$. The electric potential energy is of order $U \sim -\frac{Ze^2}{4\pi\epsilon_0 D} \sim -0.14 \text{ MeV}$ for $Z=1$. Hence, $K \gg -U$ and we would expect the electron to escape almost immediately; it cannot be bound within a nucleus.

b) my numbers: $m = 60 \text{ kg}$, $w = 0.81 \text{ m}$ width, $d = 12 \text{ cm}$ thickness

a) $w < 10\lambda$, $\lambda = h/mv \Rightarrow v < 10h/mw = 1.36 \times 10^{-34} \text{ m/s}$

b) $t = d/v = 8.8 \times 10^{32} \text{ s} = 2.2 \times 10^{15}$ times age of Universe

c) not in your lifetime!

7) Assume that nonrelativistic kinematics is adequate, such that $K_{\max} = (eBR)^2/2m_e$ and $\phi = K_{\max} - hc/\lambda$.

my numbers: $\lambda = 328 \text{ nm}$, $B = 2.10 \times 10^{-5} \text{ T}$, $R = 0.2 \text{ m}$

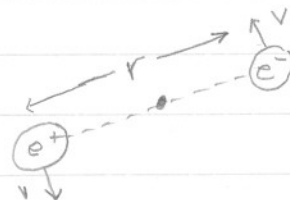
$K_{\max} = 1.55 \text{ eV}$, $E_\gamma = 3.78 \text{ eV} \Rightarrow \phi = 2.23 \text{ eV}$

Note that $K_{\max} \ll m_e c^2 \Rightarrow$ nonrelativistic kinematics are fine.

67) Assume that nonrelativistic kinematics can be used.

energy: $E = 2\left(\frac{1}{2}mv^2\right) - \frac{k_E e^2}{r}$

angular momentum: $L = 2mv\left(\frac{r}{2}\right) = mvr = n\hbar$



circular motion: $\frac{mv^2}{r/2} = \frac{k_E e^2}{r^2} \Rightarrow mv^2 = \frac{1}{2} \frac{k_E e^2}{r} \Rightarrow E = -\frac{1}{2} \frac{k_E e^2}{r}$

$$m\left(\frac{n\hbar}{mr}\right)^2 = \frac{1}{2} \frac{k_E e^2}{r} \Rightarrow r = 2 \frac{(n\hbar)^2}{k_E m e^2}$$

$$\therefore E_n = -\frac{E_0}{n^2}, \quad E_0 = \frac{\alpha^2 m_e c^2}{4} = 6.80 \text{ eV}$$

$$r_n = n^2 a_0, \quad a_0 = 2 \frac{\hbar c}{\alpha m_e c^2} = 0.106 \text{ nm}$$

9) Bohr $\Rightarrow E_n = -E_0 n^{-2} \Rightarrow E_n - E_{n-1} = E_0 \frac{2n-1}{n^2(n-1)^2} = \Delta E_n$

$$n \rightarrow \infty \Rightarrow \Delta E_n \approx E_0 \frac{2n-1}{n^4} \approx 2E_0 n^{-3}$$

classically: $f = \frac{v}{2\pi r}$ with $v = \frac{n\hbar}{m_e r}$, $r = n^2 a_0 \Rightarrow f = \frac{\hbar}{2\pi m_e a_0^2 n^{-3}}$

$$E_0 = \frac{\alpha^2 m_e c^2}{2}, \quad a_0 = \frac{\hbar c}{\alpha m_e c^2} \Rightarrow \Delta E_n = \alpha^2 m_e c^2 n^{-3}, \quad f = \hbar^{-1} \alpha^2 m_e c^2 n^{-3}$$

(note $\alpha = k_E e^2 / \hbar c$)

$$\therefore \Delta E_n = hf$$

The classical orbital frequency for large n agrees with quantum mechanical transition energy, divided by h , for adjacent orbits.

Can a free stationary particle absorb a photon without creating one or more additional particles in the final state?
Energy and momentum conservation require

$$E' = E_\gamma + mc^2$$

$$p' = p_\gamma$$



but the invariance of rest mass requires

$$(mc^2)^2 = E'^2 - (p'c)^2 = (E_\gamma + mc^2)^2 - p_\gamma^2 c^2$$

Finally, using $E_\gamma = p_\gamma c$ we find that this equation requires $E_\gamma = 0$. Therefore, the process is impossible.