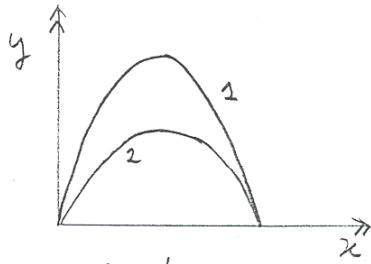


Homework 3 Solutions - Physics 161, SPG 03

s1/a). Since $h_{\max} = \frac{v_{oy}^2}{2g}$, we can definitely say that $v_{oy1} > v_{oy2}$.



b). No. Flight time depends on v_{oy} - as we saw $t_{\text{flight}} = \frac{2v_{oy}}{g}$ - so $t_{\text{flight } 1} > t_{\text{flight } 2}$.

c). since $t_{\text{flight } 1} > t_{\text{flight } 2}$ but $x_{\max 1} = x_{\max 2}$ & $x_{\max} = v_{ox} t_{\text{flight}}$ $\Rightarrow v_{ox1} < v_{ox2}$.

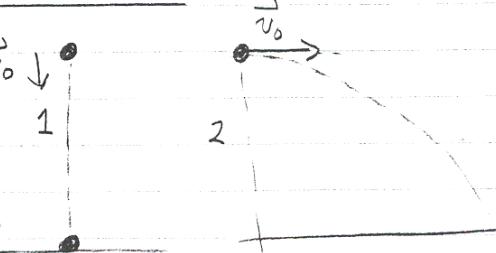
It is possible for the two projectiles to have been thrown with the same initial speed.

$$x_{\max 1} = x_{\max 2} \Rightarrow \frac{v_{oy1} v_{ox1}}{g} = \frac{v_{oy2} v_{ox2}}{g}$$

$v_{o1}^2 \sin 2\theta_1 = v_{o2}^2 \sin 2\theta_2$. If $v_{o1} = v_{o2}$, we find $\sin 2\theta_1 = \sin 2\theta_2$.

Since $v_{oy1} \neq v_{oy2}$, $\theta_1 \neq \theta_2$. It turns out that there is a solution for $\theta_1 \neq \theta_2$. If $\theta_1 = 45^\circ + \alpha$ and $\theta_2 = 45^\circ - \alpha$, we can still have $x_{\max 1} = x_{\max 2}$. (Also see P2, part c).

s2/a) The ball that is thrown vertically downwards wins the race because as far as the vertical motion is concerned only velocity in the y direction matters. Therefore, the first ball has an advantage because it starts out with some



initial velocity in the y -direction and picks up more speed due to acceleration due to gravity. As far as the vertical motion is concerned, ball 2 starts out with $v_{oy} = 0$. So it loses the race.

(b). Ball 2: $v_{oy} = 0$.

$$y = y_0 + v_{oy}t + \frac{1}{2}a_y t^2$$

$$h = 0 + 0 + \frac{1}{2}gt^2$$

$$\Rightarrow t = \sqrt{\frac{2h}{g}}$$

$$y=0 \rightarrow v_0$$

$$\Downarrow \vec{a} = +g\hat{j}$$

Ball 1:

$$y = y_0 + v_{oy}t + \frac{1}{2}a_y t^2$$

$$h = 0 + v_{oy}t + \frac{1}{2}gt^2$$

$$\Rightarrow 0 = -h + v_{oy}t + \frac{1}{2}gt^2$$

$$\Rightarrow t^2 + \frac{2v_{oy}}{g}t - \frac{2h}{g} = 0$$

$$y=0$$

$$v_{oy}$$

$$\Downarrow \vec{a} = +g\hat{j}$$

$$\Rightarrow t = -\frac{2v_{oy}}{g} \pm \sqrt{\frac{4v_{oy}^2}{g^2} + \frac{8h}{g}}$$

2

$$= -\frac{v_{oy}}{g} \pm \frac{v_{oy}}{g} \sqrt{1 + \frac{2hg}{v_{oy}^2}}$$

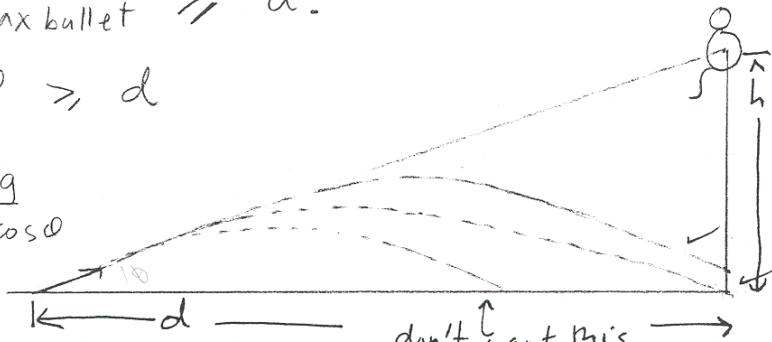
The question is, which sign should we keep. t must be greater than zero. Since the second term is larger, we must have

$$t = \frac{v_{oy}}{g} \left[\sqrt{1 + \frac{2hg}{v_{oy}^2}} - 1 \right]$$

S3). The point is that we don't want the bullet to bite the dust before it gets to the monkey i.e we want $x_{\text{max bullet}} \geq d$.

$$\Rightarrow \frac{2v_0^2 \sin \theta \cos \theta}{g} > d$$

$$v_0 \sin \theta > \frac{dg}{2 \cos \theta}$$



But since

$$x_b = v_0 \cos \theta t,$$

d is the time $t = t^*$ it would take the bullet to get to $x = d$. The time t^* on the other hand is less than or equal to the time it takes the monkey to drop to the ground i.e

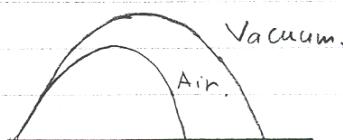
$$t^* \leq \sqrt{\frac{2h}{g}} \quad \text{so} \quad (\text{In the book, } h \leq d).$$

$$v_0 \sin \theta > \frac{t^* g}{2}$$

$$\geq \sqrt{\frac{2h}{g}} \cdot \frac{g}{2} \Rightarrow v_0 \sin \theta \geq \sqrt{\frac{hg}{2}}$$

S4).

since the horizontal component of velocity decreases gradually and the vertical component drops to zero, and since the drag force is $\propto v^2$, the effect of drag is greater on the range than on the maximum height.



Problem 55:

Note that the two things we're concerned about here is the initial velocity in the y-direction, as well as the acceleration in the y-direction.

Student 1:

Since s_1 is moving with constant speed w.r.t. Damian, we have

$$\vec{v}_{BD} = \vec{v}_{BS_1} + \vec{v}_{S_1 D}$$

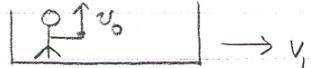
$$v_{BDy} = +v_0 + v_1$$

$$\vec{a}_{BD} = \vec{a}_{BS_1} + \vec{a}_{S_1 D}^{\uparrow} \Rightarrow \vec{a}_{BD} = g$$

$$\Rightarrow t_{max1} = \frac{v_0 y}{g} = \frac{v_0 + v_1}{g}, \quad \text{according to } S_1, \quad t_{max1} = \frac{v_0}{g}, \quad \text{according to Damian.}$$



Student 2: $\vec{v}_{BD} = \vec{v}_{BS_2} + \vec{v}_{S_2 D}$



$$v_{BDoy} = v_{BS_2 oy} + v_{S_2 D oy}$$

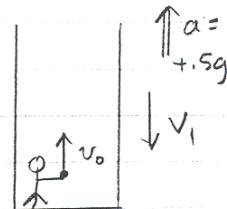
$$v_{BDoy} = v_0, \quad \text{and } \vec{a}_{BD} = g \text{ as above}$$

so $t_{max2} = \frac{v_0}{g}$; $t_{max2} = \frac{v_0}{g}$, according to S_2 , according to Damian.

Student 3: $\vec{v}_{BD} = \vec{v}_{BS_2} + \vec{v}_{S_2 D}$

$$v_{BDoy} = v_{BS_2 oy} + v_{S_2 D oy}$$

$$v_{BDoy} = v_0 - v_1$$



$$\vec{a}_{BD} = g \Rightarrow t_{\max}^{(3)} \text{ according to Damian} = \frac{v_0 - v_1}{g}$$

$$\vec{a}_{BS_3} = ? \quad \text{Use } \vec{a}_{BS_2} = \vec{a}_{BD} + \vec{a}_{DS_2} \\ = \vec{a}_{BD} - \vec{a}_{S_2 D} \\ \Rightarrow a_{BS_2 y} = -g - (+\tfrac{1}{2}g)$$

$$a_{BS_2 y} = -\tfrac{3}{2}g \Rightarrow t_{\max S_3} = \frac{v_0}{\tfrac{3}{2}g} = \frac{2}{3} \frac{v_0}{g}$$

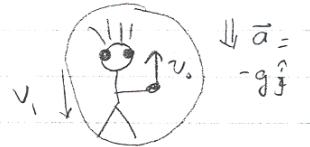
student 4. Now we've got the hang of it, so

$$t_{\max \text{ Damian}} = \frac{v_0 + v_1}{g}$$

$$t_{\max S_4} = \frac{v_0}{\tfrac{3}{2}g} \Rightarrow t_{\max S_4} = \frac{2}{3} \frac{v_0}{g}.$$

student 5. Freely falling Student

$$t_{\text{flight Damian}} = \frac{v_0 - v_1}{g}$$



$$\left. \begin{aligned} \vec{a}_{BS_5} &= \vec{a}_{BD} + \vec{a}_{DS_5} \\ &= \vec{a}_{BD} - \vec{a}_{S_5 D} \end{aligned} \right\} \text{with respect to the freely falling student. It keeps moving upwards with velocity } v_0. \text{ So there is no instant where the ball comes momentarily to rest with S}_5. \text{ So } t_{\max} \text{ is not applicable here!}$$

This implies that the ball does not accelerate

$$\text{so } t_{\max} s_1 = t_{\max} s_2 > t_{\max} s_3 = t_{\max} s_4.$$

$$t_{\max}^{(1)} \underset{\text{Dam.}}{=} t_{\max}^{(4)} \underset{\text{Dam.}}{=} \dots > t_{\max}^{(2)} \underset{\text{Dam.}}{=} \dots > t_{\max}^{(3)} \underset{\text{Dam.}}{=} t_{\max}^{(5)} \underset{\text{Dam.}}{=} \dots$$

Problem P1.

(a). The key difference is that in case I, the ball has's acceleration does not match the acceleration of the cart in the x-direction. However in Case II, even when the ball is air-borne, its acceleration i.e $\vec{a} = -g\hat{j}$ has a component parallel to the incline whose magnitude and direction is the same as the cart's acceleration along the incline. Since the ball has the same speed along the incline when the ball leaves the cart, and its accel. along the incline is also the same, it is always directly above the cart. So it falls back in the cart when its vertical position matches that of the cart.

(b). Case I:

$$\vec{r}_c = (x_0 + v_{0x_c} t + \frac{1}{2} a_{x_c} t^2) \hat{i} + (y_0 + v_{0y_c} t + \frac{1}{2} a_{y_c} t^2) \hat{j}$$

$$\Rightarrow \vec{r}_c = (v_{0x_c} t + \frac{1}{2} a_{x_c} t^2) \hat{i}$$

$$\vec{r}_{ball} = (x_0 + v_{0x_b} t + \frac{1}{2} a_{x_b} t^2) \hat{i} + (y_0 + v_{0y_b} t + \frac{1}{2} a_{y_b} t^2) \hat{j}$$

$$= (v_{0x_b} t) \hat{i} + (v_{0y_b} t - \frac{1}{2} g t^2) \hat{j}$$

For the ball to fall back in the cart at some time $t = t^*$, we must have $\vec{r}_{cart}(t^*) = \vec{r}_{ball}(t^*)$
 $\Rightarrow x_c(t^*) = x_b(t^*) \text{ & } y_c(t^*) = y_b(t^*)$.

$$y_c = y_b$$

$$\Rightarrow 0 = v_{0y_b} t - \frac{1}{2} g t^2 \Rightarrow \boxed{t^* = \frac{2v_{0y}}{g}}$$

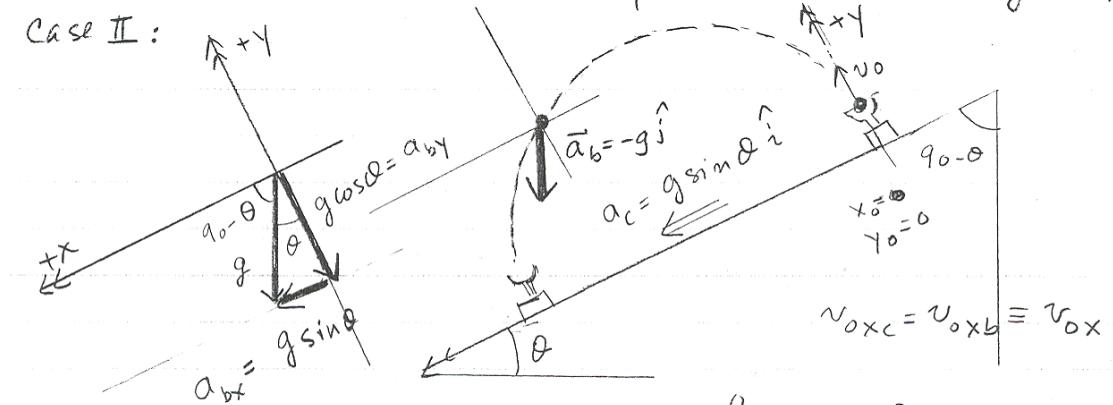
$$\text{But imposing } x_c(t^*) = x_b(t^*)$$

$$\Rightarrow v_{oxc}t + \frac{1}{2}a_{xc}t^2 = v_{oxb}t$$

$$\text{since } v_{oxc} = v_{oxb}, \text{ we find } \frac{1}{2}a_{xc}t^2 = 0.$$

i.e. the only solution for $a_{xc} \neq 0$ is $t=0$. So the ball does not ever meet up with the cart again.

Case II:



$$\vec{r}_{\text{cart}} = (v_{oxc}t + \frac{1}{2}a_{xc}t^2)\hat{i} + (v_{oyc}t + \frac{1}{2}a_{yc}t^2)\hat{j}$$

$$\vec{F}_{\text{cart}} = (v_{ox}t + \frac{1}{2}gsin\theta t^2)\hat{i}$$

$$\begin{aligned}\vec{r}_{\text{ball}} &= (v_{oxb}t + \frac{1}{2}a_{xb}t^2)\hat{i} + (v_{oyb}t + \frac{1}{2}a_{yb}t^2)\hat{j} \\ &= (v_{ox}t + \frac{1}{2}gsin\theta t^2)\hat{i} + (v_{oy}t - \frac{1}{2}gcos\theta t^2)\hat{j}\end{aligned}$$

To have $\vec{r}_{\text{cart}}(t) = \vec{r}_{\text{ball}}(t)$, we must have

$$\Rightarrow x_c(t=t^*) = x_b(t=t^*) \quad \& \quad y_c(t=t^*) = y_b(t=t^*)$$

$$\textcircled{1} \quad x_c = x_b$$

$$\Rightarrow v_{ox}t + \frac{1}{2}g sin^2\theta t^2 = v_{ox}t + \frac{1}{2}g sin^2\theta t^2$$

so $x_c = x_{\text{ball}}$ all the time! so when $y_b = y_c = 0$
the ball falls back in the cart at $t^* = \frac{2v_{oyb}}{g cos\theta}$

P21.(a). Need $x_{\max} = h_{\max}$.

Recall at $y = h_{\max}$, $v_y = 0$. So

$$v_{fy}^2 = v_{oy}^2 + 2ax$$

$$0 = v_{oy}^2 - 2g(y_f - y_0)$$

$$0 = v_{oy}^2 - 2g(h_{\max} - 0)$$

$$\Rightarrow h_{\max} = \frac{v_{oy}^2}{2g}$$

$$t_{\max} = \frac{v_{oy}}{g} \Rightarrow t_{\text{flight}} = \frac{2v_{oy}}{g} = \frac{2v_o \sin \theta}{g}$$

$$x_{\max} = (v_o \cos \theta) t_{\text{flight}} = v_o \cos \theta \cdot \frac{2v_o \sin \theta}{g}$$

$$= \frac{2v_o^2 \sin \theta \cos \theta}{g}$$

$$h_{\max} = x_{\max} \Rightarrow \frac{y_o \sin^2 \theta}{2g} = \frac{2v_o^2 \sin \theta \cos \theta}{g}$$

$$\Rightarrow \frac{\sin \theta}{\cos \theta} = 4 \Rightarrow \tan \theta = 4$$

$$\Rightarrow \theta = \tan^{-1} 4 \approx 76^\circ.$$

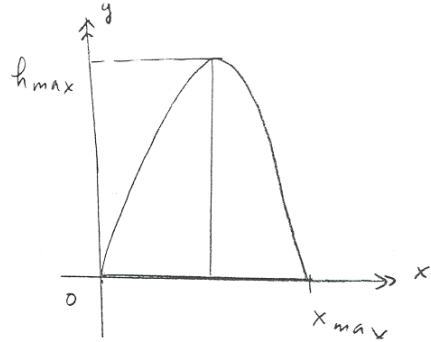
b) $x_{\max} = \frac{1}{2} h_{\max}$

$$\frac{2v_o^2 \sin \theta \cos \theta}{g} = \frac{x_o^2 \sin^2 \theta}{4g} \Rightarrow \tan \theta = 82.9^\circ$$

c) $x_{\max} (\theta = 45^\circ \pm \alpha) = \frac{2v_o^2 \sin 2\theta}{g} = \frac{2v_o^2 \sin(90^\circ \pm 2\alpha)}{g}$

$$= \frac{2v_o^2}{g} [\sin 90^\circ \cos 2\alpha \pm \cos 90^\circ \sin 2\alpha] = \frac{2v_o^2 \cos 2\alpha}{g}.$$

so for both $\theta = 45^\circ + \alpha$ & $45^\circ - \alpha$, $x_{\max} = \frac{2v_o^2 \cos 2\alpha}{g}$.



Problem P4.

a). The upper limit for H corresponds to the limit when $v_0 \rightarrow \infty$. In that case, the trajectory of the motorbike is essentially a straight line. (see Fig P4 b.)

$$\text{Therefore } \frac{H}{W} = \tan \theta$$

$$\text{or } H = W \tan \theta.$$

In real life, $v_0 \rightarrow \infty$ simply means v_0 is very very large. But what does it mean to say very very large? It must be very large compared to something. We will see that shortly,

$$b/. y = y_0 + v_{0y} t + \frac{1}{2} a_y t^2$$

$$H = 0 + v_0 \sin \theta t - \frac{1}{2} g t^2$$

At the time that the motorbike is at $y = H$, we want it to be at $x = W$.

$$x = x_0 + v_{0x} t + \frac{1}{2} a_x t^2$$

$$W = v_0 \cos \theta t \Rightarrow t = \frac{W}{v_0 \cos \theta}$$

$$\Rightarrow H = x_0 \sin \theta \cdot \frac{W}{v_0 \cos \theta} - \frac{1}{2} g \left(\frac{W}{v_0 \cos \theta} \right)^2$$

$$H = W \tan \theta - \frac{1}{2} g \frac{W^2}{v_0^2 \cos^2 \theta}$$

[Note that when $v_0^2 \gg g W$, $H \rightarrow W \tan \theta$.]

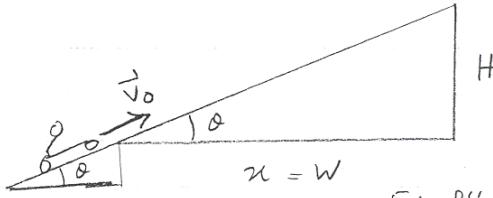
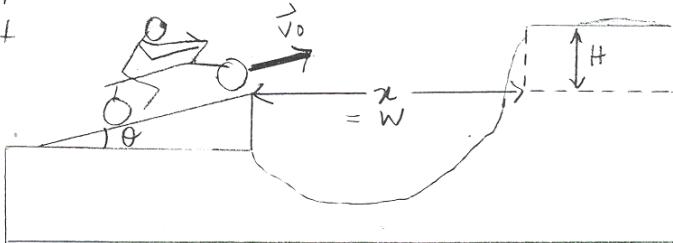
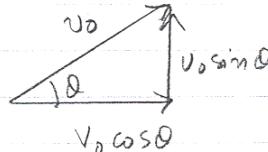


Fig P4 b.



That means that $v_0 \rightarrow \infty$ is physically the same thing as saying that v_0 is much much larger than \sqrt{gW} . Note that \sqrt{gW} has dimensions of speed. $\sqrt{\frac{L}{T^2} \cdot L} = \sqrt{\frac{L^2}{T^2}} = \frac{L}{T}$.

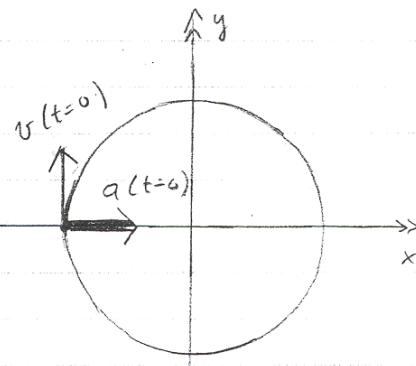
$$\Rightarrow \frac{w^2}{v_0^2 \cos^2 \theta} = \frac{2}{g} (w \tan \theta - H)$$

$$\Rightarrow v_0^2 = \frac{g w^2}{2(w \tan \theta - H) \cos^2 \theta}$$

when $\theta = 0$ $v_0^2 = \frac{g w^2}{-2H}$. This is non-sensical because v_0^2 cannot be negative. But this simply indicates that if $\theta = 0$, $v_{0y} = 0$, so the motorbike cannot gain any height whatsoever.

P5/. $\vec{r}(t) = -10m \cos \omega t \hat{i} + 10m \sin \omega t \hat{j}$

$$\begin{aligned} \text{(a)} |\vec{r}(t)| &= \sqrt{r_x^2 + r_y^2} \\ &= \sqrt{(-10m \cos \omega t)^2 + (10m \sin \omega t)^2} \\ &= \sqrt{(10m)^2 (\cos^2 \omega t + \sin^2 \omega t)} \end{aligned}$$



$|r(t)| = 10m$ since the magnitude of the position vector is independent of time, i.e., constant, it implies that the object moves in a circle.

(Note that at $t=0$, the object is at $r(t=0) = -10m \hat{i}$.)

$$b). \vec{v}(t) = \frac{d\vec{r}}{dt} = +(10m)\omega \sin \omega t \hat{i} + (10m)\omega \cos \omega t \hat{j}$$

$$|\vec{v}(t)| = \sqrt{v_x^2 + v_y^2} = \sqrt{(10m)^2 \omega^2 \sin^2 \omega t + (10m)^2 \omega^2 \cos^2 \omega t} \\ = \sqrt{100m \omega^2 (\sin^2 \omega t + \cos^2 \omega t)} = \sqrt{(100m^2) \omega^2}$$

$$|\vec{v}(t)| = (10\omega) m \quad [\text{note } [2r] = s^{-1}] \\ = 10 \cdot 2 s^{-1} \cdot m \\ = 20 m/s. \quad \text{Note, speed is constant. So } a_{\text{tangential}} = 0.$$

$$c). \vec{a}(t) = \frac{d\vec{v}}{dt} = (10m)\omega^2 \cos \omega t \hat{i} - (10m)\omega^2 \sin \omega t \hat{j}$$

$$|\vec{a}(t)| = \sqrt{(10m)^2 \omega^4 [\cos^2 \omega t + \sin^2 \omega t]} \\ = (10\omega^2)m \\ |\vec{a}(t)| = 40 m/s^2. \quad \text{The tangential acceleration is zero since } |\vec{v}| \text{ is constant.}$$

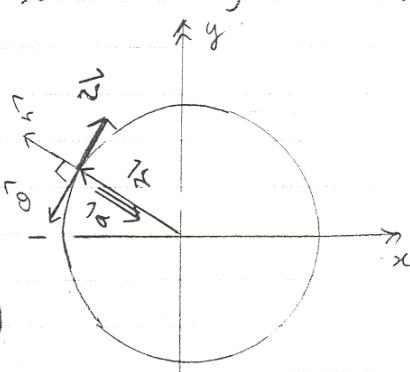
Note also: $v(t=0) = 10m\omega \hat{j}$ and
 $a(t=0) = (10m)\omega^2 \hat{i}$ so radially inwards.
 The object is moving clockwise.

$$d). \vec{r}(t) = r_0 \hat{r} = (10m) \hat{r}.$$

$$\vec{r}'(t) = -r_0 \omega \hat{\theta} \quad (\text{it's } -\hat{\theta} \text{ since}$$

The object is moving clockwise.

& $\hat{\theta}$ is normally defined + in the counterclockwise direction.)



$$\vec{r}'(t) = -r_0 \omega \hat{\theta}, \text{ negative sign because } \vec{a} \text{ points radially inwards} - \hat{r} \text{ is radially outwards.}$$

4.39 $v = (150^2 + 30.0^2)^{1/2} = \boxed{153 \text{ km/h}}$

$$\theta = \tan^{-1}\left(\frac{30.0}{150}\right) = \boxed{11.3^\circ \text{ north of west}}$$

- 4.40 For Alan, his speed downstream is $c + v$, while his speed upstream is $c - v$. Therefore, the total time for Alan is

$$t_1 = \frac{L}{c+v} + \frac{L}{c-v} = \boxed{\frac{2L/c}{1-v^2/c^2}}$$

For Beth, her cross-stream speed (both ways) is $\sqrt{c^2 - v^2}$

Thus, the total time for Beth is

$$t_2 = \frac{2L}{\sqrt{c^2 - v^2}} = \boxed{\frac{2L/c}{\sqrt{1-v^2/c^2}}}$$

Since $1 - \frac{v^2}{c^2} < 1$, $t_1 > t_2$, or Beth, who swims cross-stream, returns first.

- 4.41 α = Heading with respect to the shore

β = Angle of boat with respect to the shore

- (a) The boat should always steer for the child at heading

$$\alpha = \tan^{-1} \frac{0.600}{0.800} = \boxed{36.9^\circ}$$

(b) $v_x = 20.0 \cos \alpha - 2.50 = 13.5 \text{ km/h}$

$$v_y = 20.0 \sin \alpha = 12.0 \text{ km/h}$$

$$\beta = \tan^{-1} \left(\frac{12.0 \text{ km/h}}{13.5 \text{ km/h}} \right) = \boxed{41.6^\circ}$$

$$(c) t = \frac{d_y}{v_y} = \frac{0.600 \text{ km}}{12.0 \text{ km/h}} = 5.00 \times 10^{-2} \text{ h} = \boxed{3.00 \text{ min}}$$

