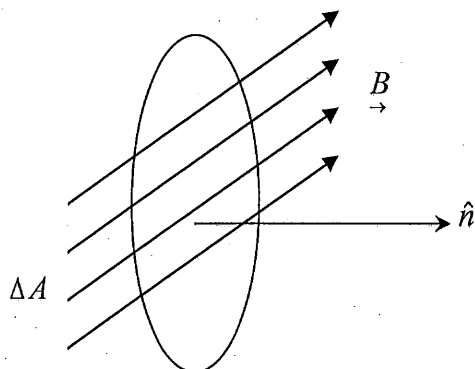


NON - COULOMB \underline{E} FIELD

(INDUCTION)

We begin by considering a uniform \underline{B} -field represented by a set of parallel lines. Next, imagine an area ΔA whose normal



is along \hat{n} . Then, as in the case of \underline{E} , we define flux of \underline{B} through ΔA to be

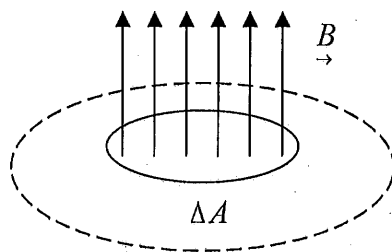
$$\Delta \Phi_B = \underline{B} \cdot \underline{\Delta A} = B \Delta A \cos(\underline{B}, \hat{n})$$

As before, flux is maximum when $\underline{B} \parallel \hat{n}$ (\perp to area) and zero if $\underline{B} \perp \hat{n}$ (\underline{B} lines lie in plane of ΔA).

Faraday's discovery was that if the flux of \underline{B} , $\Delta \Phi_B$, changes as a function of time, that is

$$\frac{\Delta \Phi_B}{\Delta t} \neq 0$$

There will be an \underline{E} -field induced in every loop surrounding the region where Φ_B is changing. For example, in the picture if the \underline{B} field varies with time an \underline{E} -field will appear at every point of the dashed loop.



This new \underline{E} -field we call a Non-Coulomb \underline{E} -field. The following points are noteworthy:

1. Flux of \underline{B} can vary with time in 3 ways

- (i) Magnitude of \vec{B} is a function of time
 - (ii) The area ΔA where \vec{B} is non-zero varies with time
 - (iii) Angle between \vec{B} and \hat{n} varies with time
2. If there is an \vec{E} -field at every point on a closed loop, there must be an $\mathcal{E}mf$ in the loop given by

$$\mathcal{E} = \sum_C \vec{E}_{NC} \cdot \underline{\Delta l}$$

since $\mathcal{E}mf$ is the work done on a unit charge by an \vec{E} -field.

3. Since the change of potential on a closed loop is NOT zero, \vec{E}_{NC} is NOT A CONSERVATIVE FIELD.
4. The \vec{E}_{NC} field lines close on themselves, there is no beginning and no end. This is totally different from the Coulomb \vec{E} which “started” at positive charges and “ended” at negative charges.
5. For \vec{E}_{NC} Gauss’ Law will always give

$$\sum_C \vec{E}_{NC} \cdot \underline{\Delta A} \equiv 0$$

Total flux of \vec{E}_{NC} through any closed surface is always ZERO.

6. However, a charge q placed in \vec{E}_{NC} experiences a force

$$\vec{F}_E = q \vec{E}_{NC}$$

exactly as for a Coulomb \vec{E} -field.

Although Faraday discovered the induced $\mathcal{E}mf$, it was left to Lenz to assert that the direction of \vec{E}_{NC} must be such as to oppose the change in the flux of \vec{B} that causes the \vec{E}_{NC} , namely

$$\mathcal{E} = \sum_C \vec{E}_{NC} \cdot \underline{\Delta l} = \frac{-\Delta \Phi_B}{\Delta t}$$

Eventually, Helmholtz showed that the minus sign on the right side is crucial because it arises from Conservation of Energy. Omitting the minus sign would lead to a “run away” situation causing an energy catastrophe.

To Summarize the above we say:

The circulation of non-coulomb \vec{E} around a closed loop is determined by the time rate of change of the flux of \vec{B} through the area inside the loop. The sense of \vec{E}_{NC} is such as to oppose the change in the flux of \vec{B}

Applications

1. Consider the picture.

There is a uniform $\underline{B} = B\hat{y}$ everywhere.

The radius of the loop is r . Let \underline{B}

increase with time $\frac{\Delta B}{\Delta t} > 0$. It will

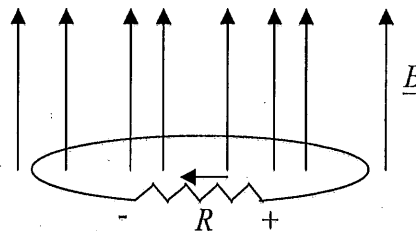
set up E_{NC} in the loop.

$$\mathcal{E} = E_{NC} 2\pi r = -\pi r^2 \frac{\Delta B}{\Delta t}$$

and the direction of E_{NC} in the loop must be clockwise. Current in R will flow from left to right.

$$I = \frac{\pi r^2}{R} \frac{\Delta B}{\Delta t}$$

[Notice: the loop acts like a "battery" and current flows from $-$ to $+$ inside the battery and from $+$ to $-$ in the load R .]



2. The coil of resistance R is in the xz -plane while $\underline{B} = B\hat{y}$

Flip the coil by 180°

Flux change will be

$$\begin{aligned} \Delta\Phi_B &= -B\pi r^2 - B\pi r^2 \\ &= -2B\pi r^2 \end{aligned}$$

Supposing this flip takes t sec. Then

$$\frac{\Delta\Phi_B}{\Delta t} = \frac{-2B\pi r^2}{tR}$$

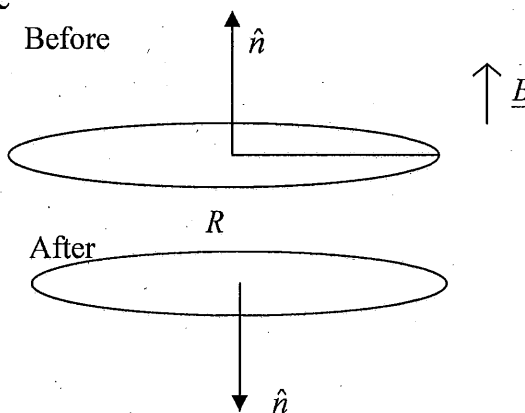
Since flux is reducing, \mathcal{E} in coil must be counter clockwise.

$$\text{current in coil} \quad I = \frac{2B\pi r^2}{tR}$$

$$\text{current is flow of charge} \quad I = \frac{\Delta Q}{t}$$

$$\text{so} \quad \Delta Q = \frac{2B\pi r^2}{R}$$

that is, when coil is flipped ΔQ coulombs of charge circulate counter clockwise in it. If you keep flipping it, current will continue to flow, *alternately CW and CCW.*



3. Calculate the E_{NC} for a solenoid at a distance r from its axis when the flux of B is varied

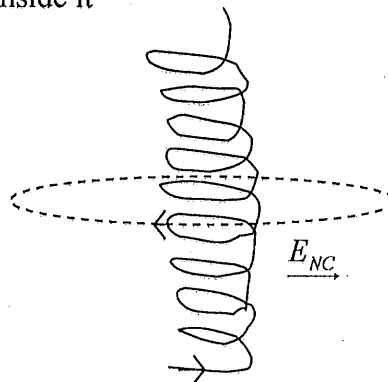
by time variation of the current in the solenoid. That is, $\frac{\Delta i}{\Delta t} \neq 0 \Rightarrow \frac{\Delta B}{\Delta t} \neq 0 \Rightarrow \frac{\Delta \Phi_B}{\Delta t} \neq 0$.

Consider a solenoid wound on a tube of radius a . If there are n turns per meter and the current flow is as shown, there is a uniform B inside it

$$\vec{B} = \mu_0 n i \hat{y}$$

and hence

$$\frac{\Delta \vec{B}}{\Delta t} = \mu_0 n \frac{\Delta i}{\Delta t} \hat{y}$$



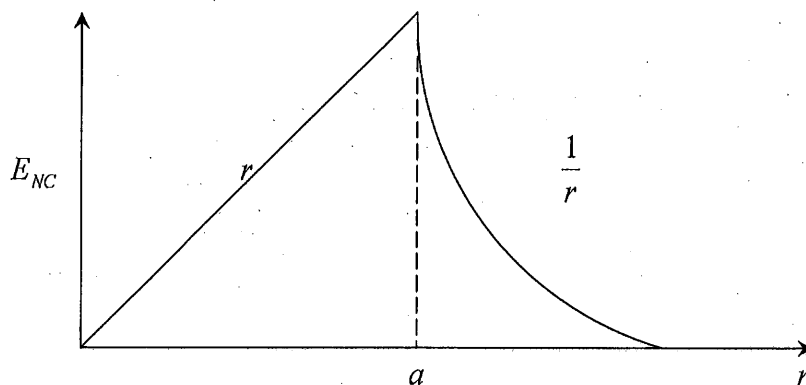
The problem has cylindrical symmetry about axis of solenoid, E_{NC} at r must be a function of r only and must be azimuthal. Let i increase with time. Then flux of B along $+\hat{y}$ is increasing with time. Take a circular loop. As shown direction of E_{NC} must be clockwise (as viewed from above) to oppose increase of Φ_B .

Next,

$$\text{If } r < a \quad E_{NC} 2\pi r = -\mu_0 n \pi r^2 \frac{\Delta i}{\Delta t} \quad [\text{LOOP INSIDE SOLENOID}]$$

$$\text{If } r > a \quad E_{NC} 2\pi r = -\mu_0 n \pi a^2 \frac{\Delta i}{\Delta t} \quad [\text{LOOP OUTSIDE SOLENOID}]$$

Magnitude



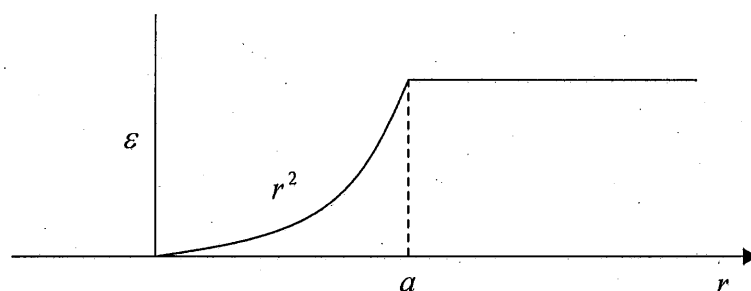
$$E_{NC} = \frac{-\mu_0 n r}{2} \frac{\Delta i}{\Delta t} \quad r < a$$

$$E_{NC} = \mu_0 n \frac{a^2}{2r} \frac{\Delta i}{\Delta t} \quad r > a$$

It is useful to ask what is the \mathcal{E} in these two loops $\mathcal{E} = \oint_C \underline{E}_{NC} \cdot \underline{\Delta l}$

$$r < a \quad \mathcal{E} = E_{NC} \cdot 2\pi r = \frac{-\mu_0 n \pi r^2 \Delta i}{\Delta t}$$

$$r > a \quad \mathcal{E} = -\mu_0 n \pi a^2 \frac{\Delta i}{\Delta t}$$



Inside \mathcal{E} increases as r^2
Outside it is constant for all r !

Non-Coulomb \underline{E} [Motional \mathcal{E}]

It is far from easy for us to understand the origin of the \underline{E}_{NC} and the consequent induced \mathcal{E} .

However, there is one analogous situation where we can use our present knowledge to see how motion of a conductor in a \underline{B} -field causes an \mathcal{E} to appear within the conductor.

Consider a copper bar of length l moving with velocity $v\hat{x}$ in the presence of a uniform \underline{B} field $\underline{B} = -B\hat{z}$

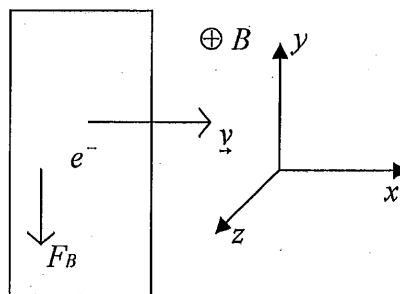
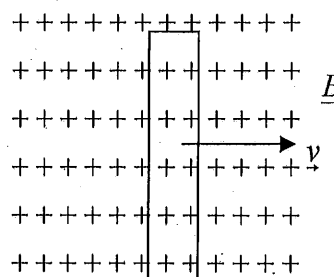
An electron inside the conductor feels the force

$$\underline{F}_B = q[\underline{v} \times \underline{B}]$$

which in the negative y-direction

$$\underline{F}_B = -evB\hat{y}$$

That will cause electrons to drift toward the bottom



leaving an excess positive charge at the top which will cause an \underline{E} field

$$\underline{E} = -E\hat{y}$$

and a force

$$\underline{F}_E = +eE\hat{y}$$

on the electron,

Further segregation of charge will stop when

$$\underline{F}_E + \underline{F}_B = 0$$

that is

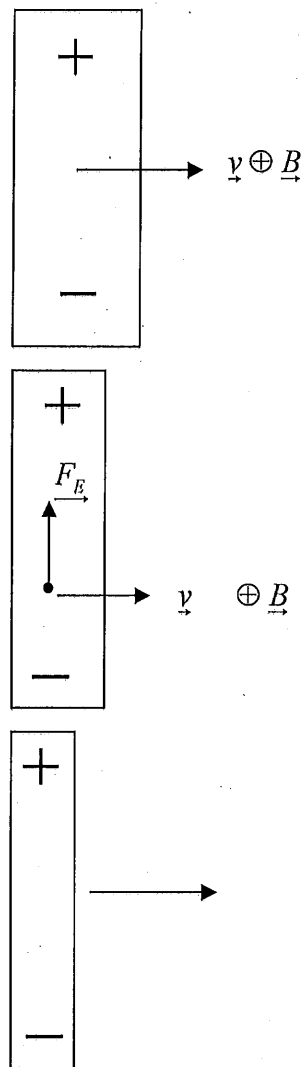
$$evB = eE$$

$$E = vB$$

and this will cause an \mathcal{E} to appear in the bar.

$$\mathcal{E} = VBl$$

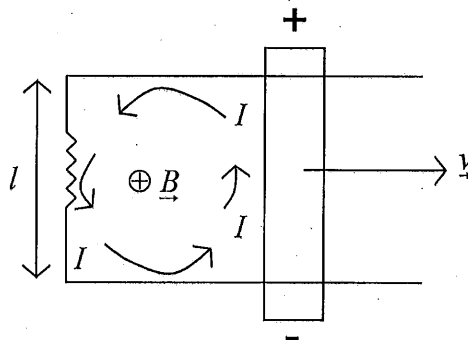
to appear in the bar.



Question \underline{B} field does no work on the electron, so where does the energy originate to set up E ?

The \mathcal{E} is exactly the answer you get if you consider the situation below and use $\mathcal{E} = \frac{-\Delta\Phi_B}{\Delta t}$

$$\begin{aligned} \frac{\Delta\Phi_B}{\Delta t} &= B \frac{\Delta A}{\Delta t} \\ &= Bvl \end{aligned}$$



and Lenz's Law tells us that the induced \mathcal{E} must oppose the increase in flux of \underline{B} into the page requiring negative at the bottom and plus at the top leading to a counterclockwise current.