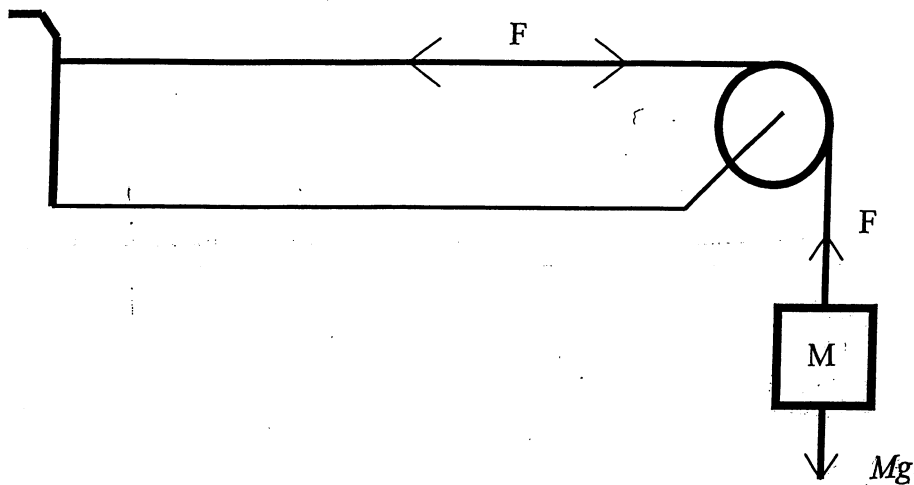


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## SPEED OF TRANSVERSE PULSE ON A STRETCHED STRING, PERIODIC WAVE, ENERGY TRANSPORT

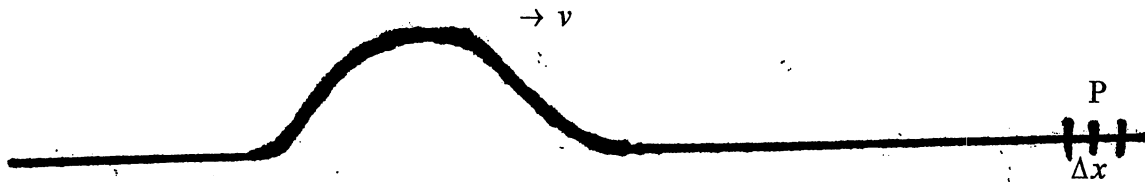
We now know that we can describe a transverse periodic wave of wavelength  $\lambda$  and a frequency  $f$  by the equation  $y = A \sin(kx - \omega t)$  with  $k = \frac{2\pi}{\lambda}$  and  $\omega = 2\pi f$  while  $A \perp \hat{x}$  with  $\omega = vk$  [same as  $v = \lambda f$ ]

To generate a "pulse" we need to sum up many, many periodic waves with different wavelengths, frequencies and amplitudes. Experimentally, all we need is to take a string of length  $L$  and mass  $m$  tie its one end, pass the other over a pulley and hang a mass  $M$ .

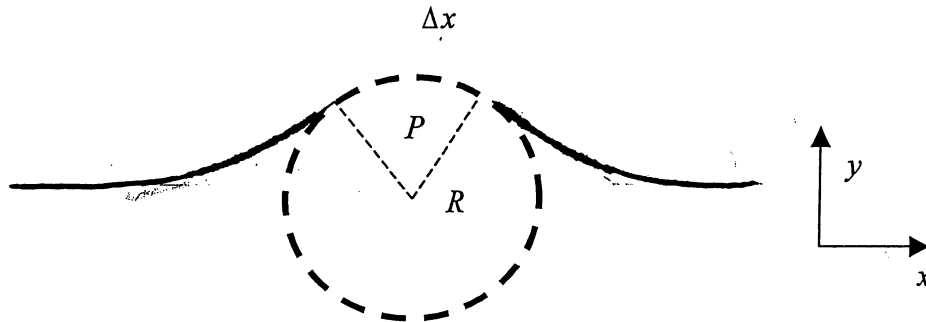


We define linear density  $\mu = \frac{m}{L}$

The string will develop tension  $F = (Mg)$  everywhere. We will make the string very long, so we do not need to worry about what happens at the ends as of yet. If we "tweak" it, we can observe a pulse such as shown below traveling along it.

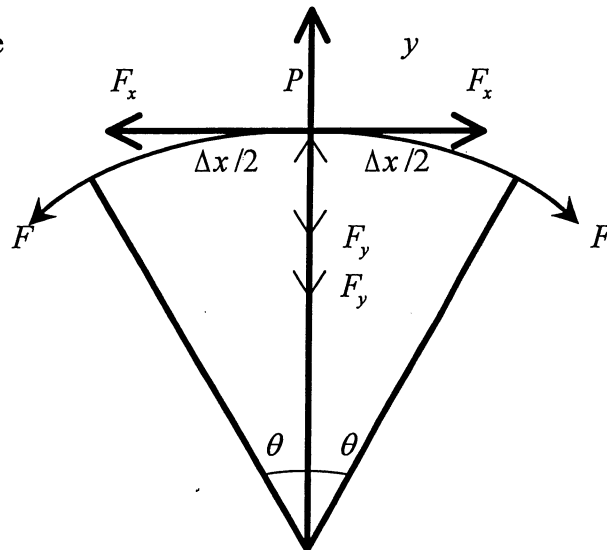


We will keep the amplitude small. Let us concentrate on a small piece of length  $\Delta x$  and ask what happens when the pulse comes along. As is clear  $\Delta x$  is lying there minding its own business when the pulse arrives and  $\Delta x$  must travel on a curved path to participate in the passage of the pulse. Indeed we can imagine that  $\Delta x$  is carried around a circle of radius  $R$  at speed  $v$ .



Since  $\Delta x$  has a mass of  $\mu\Delta x$  it needs a force  $F_c = -\frac{\Delta x v^2 \mu}{R}$  to go around the circle. The question is, what force is available to make this happen. Let us make  $\Delta x$  big and draw forces:

Immediately, we see that the net force along  $x$  (parallel to string) is zero. But the  $y$ -components due to the tension add



$$\begin{aligned}
 \text{So available force at } P \text{ is } \vec{F}_A &= -2F_y \hat{y} \\
 &= -2F \sin \theta \hat{y} \\
 &\cong -2F \theta \hat{y} \\
 &= -2F \frac{\Delta x}{2R} \hat{y} \\
 &= -F \frac{\Delta x}{R} \hat{y}
 \end{aligned}$$

Since  $\theta \ll 1$ .

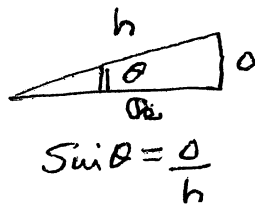
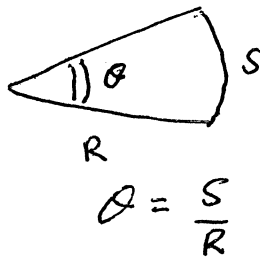
$$\theta = \frac{\Delta x}{2R}$$

While at  $P$  the needed  $\vec{F}_c$  is  $-\frac{\mu \Delta x v^2}{R} \hat{y}$ . If  $\vec{F}_A = \vec{F}_c$ ,  $\Delta x$  can happily participate in the pulse. That is, we must require  $\frac{F \Delta x}{R} = \frac{\mu \Delta x v^2}{R}$ . So  $v = \sqrt{\frac{F}{\mu}}$  is the speed of a small amplitude pulse in a string which has a tension  $F$  in it and a linear density (mass per unit length) of  $\mu \text{ kg/m}$ . It seems reasonable that for a periodic wave on our string we can write

$$\begin{aligned}
 y &= A \sin(kx - \omega t) \\
 \omega &= vk \\
 v &= \sqrt{\frac{F}{\mu}}
 \end{aligned}$$

Provided that we keep  $A \ll \lambda$  so all angles are small [we needed  $\theta \ll 1$  in our proof].

[Note that when  $\theta \ll 1$ ,  $\sin \theta \approx \theta$ ]



$$\left. \begin{aligned}
 \theta &\ll 1 \\
 a &\rightarrow h \\
 \frac{\theta}{a} &\approx \frac{\theta}{h} \approx \frac{S}{R}
 \end{aligned} \right]$$

ENERGY TRANSPORT BY SINE WAVE ON STRING

Every point on the string has a  $y$  coordinate which varies as  $y = A \sin \omega t$ . This is like linear harmonic motion so every point has a transverse velocity

$$v_y = A \omega \cos \omega t$$

A unit length of string will therefore have a kinetic energy  $K = \frac{1}{2} \mu A^2 \omega^2 \cos^2 \omega t$

Whose maximum value (which is equal to total energy, KIN + POTL) will be  $K_{\max} = \frac{1}{2} \mu A^2 \omega^2$

wave travels by  $v$  m/s so energy transport per second  $\eta = \frac{1}{2} \mu A^2 \omega^2 v$

Since  $F = \mu v^2$  we can also write  $\eta = \frac{1}{2} A^2 \omega^2 \frac{F}{v}$

\* TREAT A UNIT LENGTH OF STRING AS A "SPRING-MASS" OSCILLATOR with spring constant  $k_0$ .

$$\text{KIN. ENERGY } K = \frac{1}{2} \mu A^2 \omega^2 \cos^2 \omega t$$

$$\text{POT. ENERGY } U = \frac{1}{2} k_0 A^2 \sin^2 \omega t$$

$$\text{But } \omega = \sqrt{\frac{k_0}{\mu}}$$

$$\text{So } U = \frac{1}{2} \mu A^2 \omega^2 \sin^2 \omega t$$

Next, averaged over time  $\langle \sin^2 \omega t \rangle = \langle \cos^2 \omega t \rangle = \frac{1}{2}$

$$\text{So } \langle K \rangle + \langle U \rangle = \left( \frac{1}{4} + \frac{1}{4} \right) \mu A^2 \omega^2 = \frac{1}{2} \mu A^2 \omega^2$$

[Of course,  $\sin^2 \omega t + \cos^2 \omega t = 1$  so it is not surprising that  $\langle K \rangle + \langle U \rangle = K + U = K_{\max} = U_{\max}$

since our oscillator has no friction, total energy is CONSTANT!]