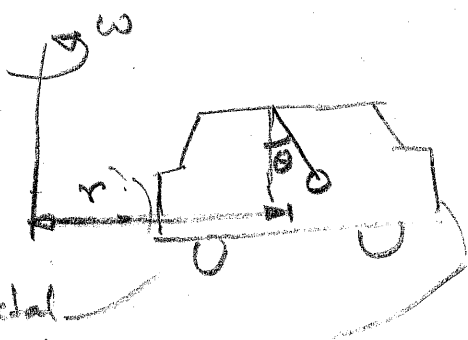


6-1

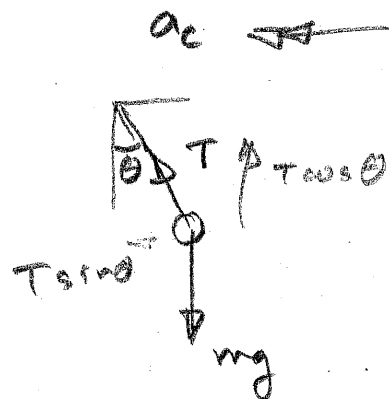
The pendulum cannot stay vertical. This is because the car along with the pendulum is describing a circular motion. The speed remains same, but due to constant change of the direction of velocity, we have a non-zero acceleration acting towards the centre of the circular path.

Now if the string of the pendulum remains vertical, the horizontal component of Tension would be zero. So there would be no force to provide the necessary centripetal force to the pendulum. Hence the string has to make an angle with the vertical to provide the necessary centripetal force to the pendulum bob.

6-2



F.B.I.D



$a_c =$  centripetal acceleration

$$= \frac{v^2}{r}$$

From F.B.I.D

$$T \cos \theta = mg$$

$$T \sin \theta = ma_c = \frac{mv^2}{r}$$

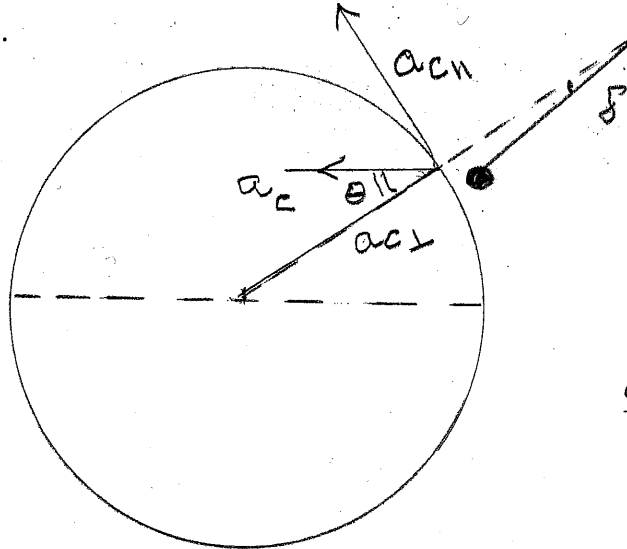
$$\therefore \tan \theta = \frac{v^2}{rg}$$

$$\theta = \tan^{-1} \left( \frac{v^2}{rg} \right) = \tan^{-1} \left( \frac{20 \times 20}{100 \times 9.8} \right)$$

$$= 0.387^\circ$$

### 6-3 This is another effect of the rotation of Earth.

This is the most subtle and happens for any  $\Theta$  other than 0 and  $\frac{\pi}{2}$  because  $\underline{a}_c$  is NOT parallel to  $\underline{g}$ .



$$\underline{a}_c = R_E \omega_E^2 \cos \Theta$$

Indeed now  $\underline{a}_c$  has a component parallel to surface of Earth

$$a_{c||} = R_E \omega_E^2 \sin \Theta \cos \Theta$$

and a component along radius of Earth

$$a_{c\perp} = -R_E \omega_E^2 \cos^2 \Theta$$

which is along  $\hat{r}$  so it modifies "g" slightly.

Since we have an  $a_{c||}$ , if you try to hang a pendulum, it cannot be vertical (parallel to  $\hat{r}$ ).

It must tilt to yield a force to produce  $a_{c||}$

$$\begin{aligned} \tan \delta &= \frac{a_{c||}}{g} \\ &\cong \frac{\sin \Theta \cos \Theta R_E \omega_E^2}{g} \end{aligned}$$

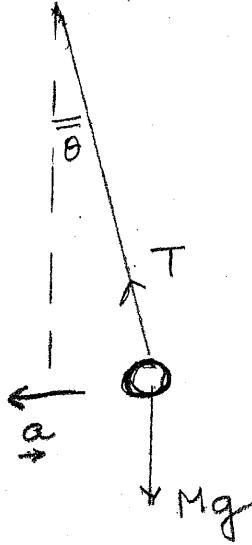
$$\delta \rightarrow 0 \quad \Theta = 0 \quad \text{and} \quad \Theta = \pi/2$$

The situation is exactly like the case of a pendulum hanging in a cart which has an acceleration  $a = -a\hat{x}$ . (Prob. 4-16)

$$-T \sin \Theta = -Ma$$

$$T \cos \Theta - Mg = 0$$

$$\tan \Theta = \frac{a}{g}$$



6-4 For a system to be inertial it cannot have an acceleration. However, the Earth goes around the Sun once a year so even at the pole there is a centripetal acceleration. In

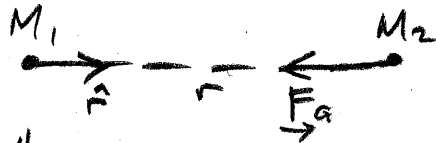
$$\text{this case } T = (365 \times 24 \times 3600) \text{ secs.}$$

$$R = 1.5 \times 10^{11} \text{ m}$$

$$\text{so } a_c = - \frac{1.5 \times 10^{11} \times \left( \frac{2\pi}{365 \times 24 \times 3600} \right)^2}{1} = -6 \times 10^{-3} \text{ m/s}^2 \hat{z}$$

Very small, but NOT zero.

(6-5) Because it is an attractive force, so the force on  $M_2$  should point to  $M_1$ .



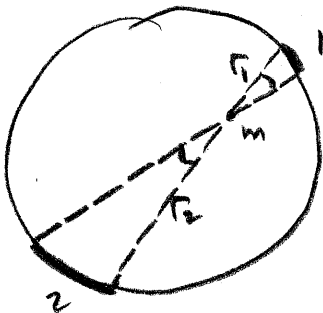
If we analyze the force on  $M_2$ , we see that  $\vec{F}_G$  will point against  $\hat{r}$ , the unit vector from  $M_1$  to  $M_2$ . So there should be a minus sign in the expression of  $\vec{F}_G$ .

Actually

$$\vec{F}_G = -\frac{GM_1 M_2}{r^2} \hat{r}$$

represents two forces ensuring that the force felt by  $M_2$  due to  $M_1$  is exactly equal but opposite to the force experienced by  $M_1$  due to  $M_2$ .

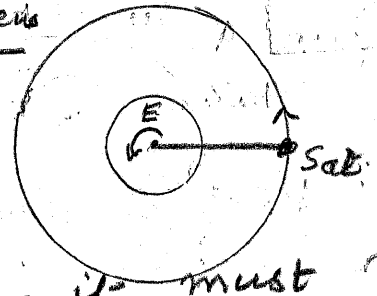
(6-6)



The force on  $m$  due to a piece of mass at position 1 is  $\frac{GmM_1}{r_1^2}$ . If the surface mass density is  $\sigma$ , then  $M_1 = \sigma A_1$ , where  $A_1$  is the surface area of  $M_1$ . So the force on  $m$  due to  $M_1$  is  $Gm\sigma A_1/r_1^2$ .

Similarly, for  $M_2$ , we have a force on  $m$  due to  $M_2$ , which is  $Gm\sigma A_2/r_2^2$ . Since the "solid angle" formed by  $M_1$  and  $M_2$  is the same, then we have  $\frac{A_1}{r_1^2} = \frac{A_2}{r_2^2}$ . Therefore, the force on  $m$  due to  $M_1$  and  $M_2$  will cancel each other. Since we choose  $M_1$  (hence  $M_2$ ) arbitrarily, we conclude that the net force on  $m$  is zero. A mass  $m$  does not feel gravitational force inside a hollow spherical surface.

T. 2nd View



The key concept is: the satellite has to be stationary with respect to a point on Earth.

Hence it must move with the same angular speed as that of earth.

Now

$$T_{\text{earth}} = 1 \text{ day} = 24 \times 60 \times 60 \text{ s} = 86400 \text{ s}$$

$$\therefore \omega = \frac{2\pi}{T} = 7.27 \times 10^{-5} \text{ rad/s}$$

Now to keep the satellite in a circular orbit an external agent has to provide the necessary centripetal force. In this case, this is the gravitational force due to Earth. So,

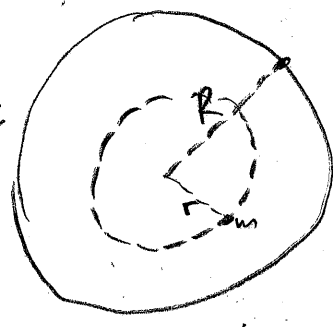
$$\frac{GMm}{r^2} = m\omega^2 r \quad \omega = \sqrt{\frac{GM}{r^3}} \quad \text{or} \quad r = \left(\frac{GM}{\omega^2}\right)^{1/3}$$

$$G = 6.7 \times 10^{-11} \frac{\text{N-m}^2}{(\text{kg})^2} \quad M_E = 6 \times 10^{24} \text{ kg}$$

$$r = \left(\frac{6.7 \times 10^{-11} \times 6 \times 10^{24}}{(7.27 \times 10^{-5})^2}\right)^{1/3} = 4.16 \times 10^7 \text{ m}$$

6-7

m is inside all the shells between r and R.



The gravitational force on a mass  $m$  at radius  $r$  is  $\vec{F}_g = -\frac{GM'm}{r^2} \hat{r}$ , where  $M'$  is the mass of the sphere inside the radius  $r$ . Why do we use  $M'$ , not  $M$  (the total mass)? Because the mass  $m$  that sits at radius  $r$  does not feel gravitational force from the mass of solid sphere that sits outside  $r$ .

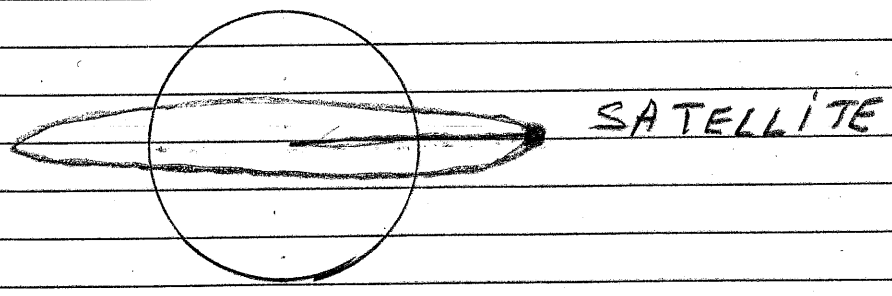
Now, since  $M_r = \rho \times \text{volume} = \rho \times \frac{4}{3} \pi r^3$ , we will have

$$\vec{F}_g = -\frac{G \rho \frac{4}{3} \pi r^3 m}{r^2} \hat{r} = -\frac{4\pi}{3} G \rho m r \hat{r} \quad (\text{for } r < R).$$

If, instead, the mass  $m$  sits on the radius  $r$  outside  $R$ , then this mass  $m$  will feel a force from the whole mass  $M$ , so the force on  $m$  is  $\vec{F}_g = -\frac{GMm}{r^2} \hat{r}$  (for  $r > R$ ).

6-8. IF THE SATELLITE IS TO REMAIN STATIONARY WITH RESPECT TO A POINT ON EARTH, THE PLANE OF ITS ORBIT MUST BE THE SAME AS THE PLANE OF THE CIRCLE TRAVELLED BY A POINT ON EARTH, SO ITS ORBIT MUST BE IN THE EQUATORIAL PLANE

ONE VIEW



6-9 Apply Newton's law on the moon of Earth,

$$\begin{aligned} \frac{GM_E m_m}{r^2} &= \frac{m_m v^2}{r} \\ &= \frac{m_m \omega^2 r^2}{r} \\ &= m_m \left(\frac{2\pi}{T}\right)^2 r \end{aligned}$$

[NOTE THAT FOR KEPLERIAN ORBITS  
 $T_{sat}^2 = \frac{4\pi^2}{GM_E} R_{sat}^3$  Earth

and  $T_{sat}^2 = \frac{4\pi^2}{GM_J} R_{sat}^3$  Jupiter]

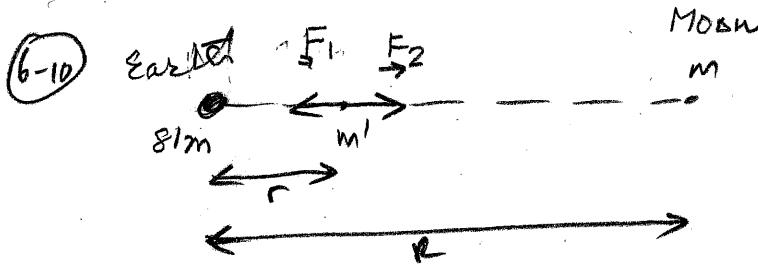
So,  $\frac{GM_E}{r^3} = \frac{4\pi^2}{T^2}$ , or  $r^3 = \frac{GM_E}{4\pi^2} T^2$

The equation for the moon of Jupiter is similar with the equation above. If we divide the two equations, we will have

$$\frac{r_1^3}{r_2^3} = \frac{\frac{GM_E}{4\pi^2} T_1^2}{\frac{GM_J}{4\pi^2} T_2^2}$$

$$\frac{(4 \times 10^5 \text{ km})^3}{(7 \times 10^5 \text{ km})^3} = \frac{M_E}{M_J} \frac{(27 \text{ days})^2}{(3.5 \text{ days})^2}$$

So  $\frac{M_E}{M_J} = 3.14 \times 10^{-3}$



There are two forces on  $m'$ . The total will be zero if they are equal and opposite.

Let's say I put a mass  $m'$  on the distance  $r$  from Earth.

Apply Newton's law,  $\frac{GMm'}{r^2} - \frac{Gmm'}{(R-r)^2} = 0$

$$\frac{GMm'}{r^2} = \frac{Gmm'}{(R-r)^2}$$

$$\frac{M}{r^2} = \frac{m}{(R-r)^2}$$

$$\left(\frac{R-r}{r}\right)^2 = \frac{m}{M}$$

$$\frac{R-r}{r} = \sqrt{\frac{m}{M}} = \sqrt{\frac{m}{81M}}$$

$$\frac{R-r}{r} = \frac{1}{9}$$

$$9(R-r) = r$$

$$9R = 10r$$

$$r = \frac{9}{10} R$$

(6-11)

$$\frac{GMm}{r^2} = \frac{mv^2}{r}$$

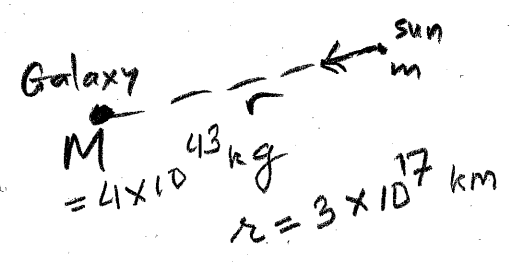
$$\text{So, } v^2 = \frac{GM}{r}$$

$$v = 9.42 \times 10^7 \text{ m/s}$$

and the period of the sun is

$$T = \frac{2\pi}{\omega} = \frac{2\pi}{v/r} = \frac{2\pi r}{v}$$

$$= \frac{2\pi \times 3 \times 10^{20}}{9.42 \times 10^7} = 2 \times 10^{13} \text{ Sec.}$$



(6-12)

The orbital speed is  $v_p = \frac{2\pi R_p}{T_p}$ . Since  $T_p^2 \propto R_p^3$ , then

$$T_p \propto R_p^{3/2}. \text{ So, } v_p = \frac{2\pi R_p}{T_p} \propto \frac{2\pi R_p}{R_p^{3/2}} = \frac{2\pi}{R_p^{1/2}}.$$

If  $R_p$  is increased, then  $v_p$  would decrease.

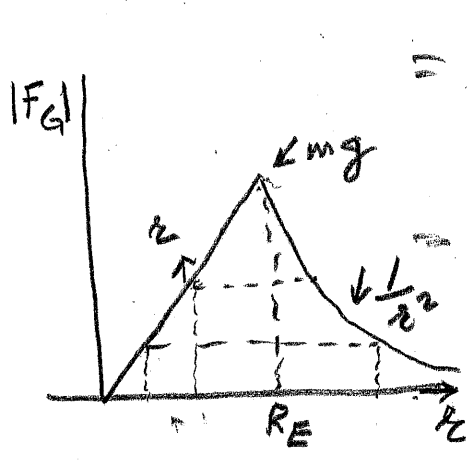


G-13

Referring to prob 6-7 we find that 9

INSIDE

$$\vec{F}_G = -\frac{4\pi}{3} G \rho m r \hat{r} \quad \text{for } r < R$$



$$= -\frac{4\pi}{3} G \frac{M}{\frac{4\pi}{3} R^3} m r \hat{r}$$

$$= -\left(\frac{GMm}{R^2}\right) \left(\frac{r}{R}\right) \hat{r}$$

At surface  $r = R_E$   $\vec{F}_G = -\frac{GMm}{R_E^2} \hat{r}$

$$= -mg \hat{r}$$

So  $\frac{GMm}{R_E^2} = mg$

or  $\vec{F}_G = -mg \left(\frac{r}{R_E}\right) \hat{r} \quad \text{for } r < R$

So if  $F_G = \frac{1}{4} mg$  then  $\frac{r}{R_E} = \frac{1}{4} \therefore r = \frac{1}{4} R_E$

Similarly  $r = \frac{1}{2} R_E$  for  $F_G = \frac{1}{2} mg$

Outside  $\vec{F}_G = -\frac{GMm}{r^2} \hat{r}$

So you get  $\frac{Mg}{2}$  at  $\sqrt{2} R_E$

$\frac{Mg}{4}$  at  $2 R_E$

6-19

$$\frac{GMm}{r^2} = m\omega^2 r$$

$$\omega = \frac{2\pi}{T}$$

$$\frac{GMm}{r^2} = \frac{4\pi^2 m r}{T^2}$$

m = mass of moon  
M = mass of earth  
r = radius of orbit

We have

$$T_1 = 27 \text{ days}$$

$$r_1 = 400,000 \text{ km}$$

$$T_2 = 1 \text{ day}$$

$$r_2 = ?$$

$$T^2 = \left(\frac{4\pi^2}{GM}\right) r^3$$

$$\left(\frac{T_1}{T_2}\right)^2 = \left(\frac{r_1}{r_2}\right)^3$$

$$r_2 = r_1 \left(\frac{T_2}{T_1}\right)^{2/3} = r_1 \left(\frac{1}{27}\right)^{2/3} = \frac{r_1}{9} \text{ km}$$

$$r_2 = \frac{400,000 \text{ km}}{9} = 4.44 \times 10^7 \text{ m}$$

(slight over estimate, compare 6-8)